Thermal study of variable conductivity and variable viscosity considering magnetic dipole on nanofluid flow with heat sink/source

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Abstract: The goal of the current study is to examine how the magnetic dipole effects on nanofluid flow over an extended surface. Based on a constant, non-porous material with a velocity slip condition, this investigation has been reported. The effects of the variable viscosity and thermal conductivity are explored for two different types of nanofluids. The experiment's findings involved dispersing in water and ethylene glycol base solutions. For both types of nanofluids, the fundamental governing equations are converted into nonlinear ordinary differential equations using the suitable transformation and solved using the bvp4c technique. The viscosity and porosity parameters decay the velocity field. Furthermore, the transport of heat is decaying function of viscous dissipation factor, growing function of Prandtl factor.

Keywords: Magnetic dipole; Variable viscosity, Nanofluid, Porous medium, Newtonian heating; Heat sink/source

1. Introduction

Currently, nanofluids have concerned a boundless deal of consideration because of their prospective in increasing the heat transport. The nano-sized metallic elements, that is, gold, iron, copper, aluminum, or their oxides are utilized as colloidal representatives with greatest collective base liquids for instance, glycol, water, ethylene, or lubricant so that value-added thermo-physical

aspects of heat transport of base fluid. Choi and Eastman [1] established a brand-new type of fluid named nanofluid. Nanofluids have uses in a number of industries, comprising transportation, energy production, electrical structures like microprocessors, and biology. Numerous scholars have looked into the topic of nanofluids both mathematically and experimentally for heat transport characteristics. The studies related to nanofluid are explored in Refs. [2-11]. Furthermore, Manigandan and Narayana [12] examined variable properties of thermal conductivity on hybrid nanofluid with slip mechanisms. Irfan [13] elaborated Joule heating influence and activation energy in Maxwell nanofluid.

Furthermore, an innovative kind of nanofluid that has newly been used in research is a hybrid nanofluid. It is formed by dissolving two different kinds of nanoparticles throughout the base fluid. The rapid growth of theoretical studies has led to the proposal of many new mathematical models. Scientists and researchers formed hybrid ferrofluids, which have nanoparticles distributed throughout, in an effort to increase fluid efficiency. However, a study of the works reveals that only a small number of studies [14–17] have concentrated on the flow and heat transmission of hybrid ferrofluids. Research on MHD hybrid ferrofluid and the impact of erratic heat sources and sinks on the flow of radiative thin films was suggested by Kumar et al. [18]. Tlili et al. [19] studied MHD hybrid ferrofluid under the impact of asymmetrical heat (rise/fall) was then continued. Manhet al. [20] investigated the role of radiation in the heat transmission of an MHD hybrid ferrofluid submerged in porous media. In a mixed convection flow on a hybrid ferrofluid ($Fe_3O_4 - CoFe_2Fe_2$), Zainodin et al. [21] examined exponentially deformable sheets at stagnation points with the effects of heat source/sink and velocity slip.

The study analysis of the works revealed that no study of hybrid nanoparticles that affect the thermal growth of heat in ferrofluid when it is exposed to a magnetic dipole. The mathematical model by Tiwari and Das [22] has been employed to address this issue. The current article explores the theory of nanofluid [23-25] in view of convective conditions on ferro-magnetic hybrid nanofluid [26-27] with variable viscosity [28-32], thermal conductivity in porous surface considering magnetic dipole. The byp4c method utilized for solutions.

2. Description of the problem

Here the flow of ferromagnetic nanofluid $(\gamma Al_2O_3 \ H_2O, \gamma Al_2O_3 \ C_2H_6O_3)$ along stretched surface is taken into consideration. This flow is stable, two-dimensional, and incompressible.

2.1.Model assumptions and conditions

The following assumptions are taken into account when analyzing the mathematical model:

- Nanoparticles
- Variable thermal conductivity
- Magnetic dipole
- Varying viscosity
- Porous substance
- Heat sink/source

The sheet is stretched with velocity $U_{\omega} = cx$ denotes the velocity of the extended sheet. The temperature *T* of a fluid temperature of the surface and T_{ω} and T_c denote the initial and ambient temperatures, respectively, i.e., $T_{\omega} < T_c < T_c$. The magnetic dipole study are also considered shown in Fig.[1].

The governing equations are

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \tag{1}$$

$$\rho_{nf}\left(u\frac{\partial u}{\partial x}+v\frac{\partial v}{\partial y}\right) = \frac{\partial}{\partial y}\left(\mu_{nf}\left(T\right)\frac{\partial u}{\partial y}\right)+\mu_{0}M\frac{\partial H}{\partial x}-\frac{\varepsilon v_{nf}}{k}u$$

$$\left(u\frac{\partial T}{\partial x}+v\frac{\partial T}{\partial y}\right)+\frac{\mu_{0}}{\left(\rho c_{p}\right)_{nf}}T\frac{\partial M}{\partial T}\left(u\frac{\partial H}{\partial x}+v\frac{\partial H}{\partial y}\right)=\frac{1}{\left(\rho c_{p}\right)_{nf}}\frac{\partial}{\partial y}\left(K_{nf}\left(T\right)\frac{\partial u}{\partial y}\right)+\frac{Q_{0}\left(T-T_{c}\right)}{\left(\rho c_{p}\right)_{nf}}$$

$$(3)$$

Here μ_{nf} , ρ_{nf} , v_{nf} , ε , K, $(\rho c_p)_{nf}$, K_{nf} , M, H, c_p , Q_0 denotes the dynamic viscosity, density, kinematic viscosity, Porosity and permeability of porous medium, specific heat, thermal conductivity of the nanofluid, magnetization, magnetic field, specific heat and heat sink/source coefficient, respectively.

The boundary conditions are:

$$u = U_{\omega}(x) = cx, v = 0, \frac{\partial T}{\partial y} = -h_s T \bigg\} at \ y = 0$$
(4)

$$u \to 0, v \to 0, T \to T_{\infty} = T_c \}$$
 when $y \to \infty$ (5)

Here h_s the heat transfer coefficient.

The viscosity and thermal conductivity as an exponential function of temperature

$$\mu_{f}(T) = \mu_{f} \exp\left[-a(T - T_{c})\right]$$

$$k(T) = k \exp \varepsilon \left(\frac{T_{c} - T}{T_{c} - T_{\omega}}\right)$$
(6)
(7)

2.2.Thermo-physical properties of $\gamma A l_2 O_3 - H_2 O$ and $\gamma A l_2 O_3 - C_2 H_6 O_3$:

A few features associated with nanofluids, their dynamic viscosity, effective density, specific heat capacity and thermal conductivity.

$$\mu_{nf}\left(T\right) = \frac{\mu_{f} \exp\left(-m_{l}\theta\right)}{\left(1-\varphi\right)^{2.5}}, \ \rho_{nf} = \left(1-\varphi\right)\rho_{f} + \varphi\rho_{s}$$
(8)

$$\left(\rho c_{p}\right)_{nf} = (1-\varphi)\left(\rho c_{p}\right)_{f} + \varphi\left(\rho c_{p}\right)_{s}$$

$$k_{nf}\left(T\right) = \left(\frac{\left(k_{s}+2k_{f}\right)-2\varphi\left(k_{f}-k_{s}\right)}{\left(k_{s}+2k_{f}\right)+\varphi\left(k_{f}-k_{s}\right)}\right) * k\left(\exp\varepsilon\left(\frac{T_{c}-T}{T_{c}-T_{\omega}}\right)\right)$$
(10)

The relations determine the values of density, dynamic viscosity, specific heat and thermal conductivity of nanofluids. Base fluid and nanoparticles thermal conductivities are characterize as k_s and k_f , while φ denotes the volume of solid in the nanofluid , ρ_s and ρ_f represent the densities of the base fluid and nanoparticles, respectively.

2.3. Magnetic dipole

The magnetic field created by the magnetic dipole affects the flow of ferrofluid caused by stretching the sheet.

$$\phi = \left(\frac{\gamma}{2\pi}\right) \left(\frac{x}{x^2 + (y+d)^2}\right)$$
 is the magnetic scalar potential of a ferrofluid. (11)

Here, γ denotes the strength of the magnetic field at the source, and *d* denotes the separation between the magnetic field's centre and the *x*-axis. H_x and H_y , which are defined as

$$H_{x} = -\frac{\partial \phi}{\partial x} = \left(\frac{\gamma}{2\pi}\right) \left(\frac{x^{2} - (y+d)^{2}}{\left[x^{2} + (y+d)^{2}\right]^{2}}\right)$$
(12)

$$H_{y} = -\frac{\partial\phi}{\partial y} = \left(\frac{\gamma}{2\pi}\right) \left(\frac{2x(y+d)}{\left[x^{2} + \left(y+d\right)^{2}\right]^{2}}\right)$$
(13)

$$\left|H\right| = \left[\left(\frac{\partial\phi}{\partial x}\right)^2 + \left(\frac{\partial\phi}{\partial y}\right)^2\right]^{\frac{1}{2}}$$
(14)

$$\frac{\partial H}{\partial x} = -\left(\frac{\gamma}{2\pi}\right) \left(\frac{2x}{\left(y+d\right)^4}\right) \tag{15}$$

$$\frac{\partial H}{\partial y} = \left(\frac{\gamma}{2\pi}\right) \left(\frac{-2}{\left(y+d\right)^3} + \frac{2x^2}{\left(y+d\right)^5}\right)$$
(16)

Making use of the following similarity variables

$$\Psi(\zeta,\eta) = \left(\frac{\mu_f}{\rho_f}\right) \zeta f(\eta), \theta(\zeta,\eta) = \frac{T_c - T}{T_c - T_\omega} = \theta_1(\eta) + \zeta^2 \theta_2(\eta), \eta = \left(\frac{\rho_f}{\mu_f}c\right)^{\frac{1}{2}} y, \qquad (17)$$

$$\zeta = \left(\frac{\rho_f}{\mu_f}c\right)^{\frac{1}{2}} xu_{,} = \frac{\partial\Psi}{\partial y} = cxf'(\eta), \quad v = -\frac{\partial\Psi}{\partial y} = -\sqrt{cv_f}f(\eta)$$
(18)

Equations [6-7] and [15-16] with [18-19] gives Eqs. [2-4]

$$\frac{(1-m_{l}\theta)[f'''-m_{l}\theta'f'']}{(1-\varphi)^{2.5}\left(1-\varphi+\varphi\frac{\rho_{s}}{\rho_{f}}\right)} - (f')^{2} + ff'' = \frac{2\beta\theta_{l}}{\left(1-\varphi+\varphi\frac{\rho_{s}}{\rho_{f}}\right)(\eta+\alpha)^{4}} + \frac{\wedge_{m}f'}{\left(1-\varphi+\varphi\frac{\rho_{s}}{\rho_{f}}\right)}$$
(19)

$$\frac{\frac{k_{nf}}{k_{f}}(1+\varphi\theta_{1})\theta_{1}^{"}}{\left(1-\varphi+\varphi\frac{\left(\rho_{s}c_{p}\right)_{s}}{\left(\rho_{f}c_{p}\right)_{f}}\right)}+\frac{2\lambda\beta f}{\left(\eta+\alpha\right)^{3}}\left(\theta_{1}-\varphi\right)+CPr\left(f\theta_{1}^{'}-2f\theta_{1}^{'}\right)-4\lambda\left(f^{'}\right)^{2}+\frac{Q_{h}}{\left(1-\varphi+\varphi\frac{\left(\rho c_{p}\right)_{s}}{\left(\rho c_{p}\right)_{f}}\right)}\theta_{1}+PrQ_{h}\theta_{1}=0 \quad (20)$$

$$\frac{\frac{k_{nf}}{k_{f}}(1+\zeta\theta_{2})\theta_{2}''}{\left(1-\varphi+\varphi\frac{\left(\rho c_{p}\right)_{s}}{\left(\rho c_{p}\right)_{f}}\right)} + \frac{2\lambda\beta f}{\left(\eta+\alpha\right)^{3}}\theta_{2} - \lambda\beta\left(\theta_{1}-\zeta\right)\left(\frac{2f'}{\left(\eta+\alpha\right)^{4}} + \frac{4f}{\left(\eta+\alpha\right)^{5}}\right) - \Pr\left(4f'\theta_{2} - f\theta_{2}'\right) - \lambda\left(f''\right)^{2}\frac{Q_{h}}{\left(1-\varphi+\varphi\frac{\left(\rho c_{p}\right)_{s}}{\left(\rho c_{p}\right)_{f}}\right)}\theta_{2} + \Pr Q_{h}\theta_{2} = 0$$

$$(21)$$

subject to boundary conditions

$$f(\eta) = 0, f'(\eta) = 1, \theta'_1(\eta) = -\gamma_1(1 + \theta_1(0)), \theta_2(\eta) = 0 \text{ at } \eta = 0,$$

$$f'(\eta) \to 0, \theta_1(\eta) \to 0, \theta_2(\eta) \to 0, \text{ when } \eta \to \infty$$
(22)
(23)

Here (Prandtl number) Pr, (porosity parameter) \wedge_m , (Curie temperature ratio), (Newtonian heating parameter) γ_1 , (Viscosity parameter) m, (Viscous dissipation) λ , (ferro hydrodynamic interaction parameter) β , (dimensionless distance) α and (heat sink/source) Q_h .

$$Pr = \frac{k}{\alpha_f}, \ \wedge_m = \frac{\varepsilon v_f}{kc}, \ \ = \frac{T_{\infty}}{(T_c - T_{\omega})}, \ \gamma_1 = h_s \sqrt{\frac{v_f}{c^2}}, \ \ m = -a(T_c - T_{\omega}), \ \lambda = \frac{c\mu_f^2}{\rho K_1(T_c - T_{\omega})}$$

$$\beta = \frac{\gamma \mu_0 K_1(T_c - T_{\omega})\rho_f}{2\pi \mu_f^2}, \ \ \alpha = \sqrt{\frac{c\rho_f d^2}{\mu_f}}, \ Q_h = \frac{Q_o}{c(\rho c_p)_{nf}}$$
(24)

3. Physical interest

The Skin fraction and Local Nusselt mathematically as:

$$C_{fx} = \frac{-2\tau_{\omega}}{\rho_{nf}U_{\omega}^{2}}, \tau_{\omega} = \mu_{nf}\left(T\right)\frac{\partial u}{\partial y}|_{y=0}$$
(25)

$$Nu_{x} = \frac{xk_{nf}}{k_{f}\left(T_{c} - T_{\omega}\right)} \frac{\partial T}{\partial y}\Big|_{y=0}$$
(26)

Dimensionless form

$$\frac{1}{2} \left(Re_x \right)^{\frac{1}{2}} C_{f_x} = \frac{\left(1 - m_1 \theta(0) \right)}{\left(1 - \varphi \right)^{2.5}} (0), \left(Re_x \right)^{\frac{1}{2}} Nu_x = -\gamma_1 \frac{k_{nf}}{k_f} \left(1 + \frac{1}{\theta_1(0) + \zeta^2 \theta_1(0)} \right)$$
(28)

where $Re_x = U_{\omega}(x) x / v_f$ is local Reynold number.

4. Discussion

Here numerical approach has been worked for solutions and present graphically and in tables form.

The values of fixed parameters taken are $\beta = 0.8$, $m_1 = Q_h = 0.1$, $\lambda = 0.01$, $\varphi = 0.15$, $\gamma_1 = 0.5$, $\wedge_m = 0.8$, $\gamma = 0.5$, z = 2, $\Pr = 6.2$. Figures (2–5) elaborate on the implications m_1, β, \wedge_m and γ_1 and on velocity profile. Figure 2 depicts the influence of viscosity change m_1 for the two nanofluids. This has dealt with a fall in Ferro-fluid velocity brought about by rise in exponential viscosity. The momentum boundary layer thickens as a result of the increase in viscous forces relative to inertial forces brought on by an upsurge in viscosity. Additionally, increasing the fluid's concentration through the inclusion of nanoparticles slows down the nanoliquid. The graph illustrates that, in the scenario of $\gamma A l_2 O_3 - H_2 O$ the fluid's velocity falls as the concentration of nanoparticles rises. The axial velocity of the nano fluid decreases when the density of the base fluid is increased with solid particles. The consequences of β are apparent in Figure 3. In the existence of $\gamma Al_2O_3 - H_2O$ and $\gamma Al_2O_3 - C_2H_6O_3$, as revealed in figure 3, the fluid viscosity increases and causes the axial velocity to decrease for increasing values of β . Figure 4 displays declining performance of Λ_m in flow of $\gamma A l_2 O_3 - H_2 O$ and $\gamma Al_2O_3 - C_2H_6O_3$ for the velocity distribution. Biological applications use porous media that contain pore structures filled with fluid. The γ_1 graph is explored in fig 5. The escalation in γ_1 alter the nanofluid axial velocity, which decreases velocity for γ_1 . The impact of the viscosity parameter and the ferromagnetic parameter on the fluctuation in the temperature profile is explored by employing figures 6 and 7. The change in viscosity the resistive forces produced and results display the temperature rises. The temperature of the nanofluid is revealed to be elevated by deeper values of β as in Figure 7. This arises from the interaction of a magnetic field's action with nanoparticle motions, i.e., rise of temperature field caused decay in nanoparticle activities. Figures 8 and 9 illustrate a boost in Curie temperature relation, leading to an enhancement in the fluid's temperature profile. The influence of λ (viscous dissipation parameter) on the temperature contour is perceived in Figure 10. Regardless of the growing λ , it appears that the temperature contour improves. γ_1 impact on temperature profile reports in figure 11. Increase in γ_1 increase heat transport and intensifies temperature. Figure 12 illustrates a pattern towards cooling in the temperature profile. For higher quantities of Pr, the temperature field drops. As thermal diffusivity falls, heat is isolated, missing from the heated sheet and enhancing the surface temperature gradient.

Thermo physical properties of nanomaterial are shown in Table 1. Numerical worth of skin friction

for two nanofluid for Pr, β , α parameter is provided in **Table 2**. The skin friction is increased by the higher levels of these limitations. Skin friction is exacerbated by a rise in resistance force attributed to an elevation in the volume fraction of nanoparticles and a greater range of magnetic restrictions. The consequence is comparatively prominent when $\gamma Al_2O_3 - H_2O$ nanofluid is employed. As shown in **Table 3**, larger values of Pr raise the cooling influence and, as a consequence, raise the Nusselt number, leading to greater skin friction and slow the heat transmission rate. Because of the $\gamma Al_2O_3 - C_2H_6O_3$ Nano fluid's strong thermo physical characteristics, cooling occurs quickly while Prandtl number is increased. The Nusselt number decreases as the hydrodynamic interaction constraint β and viscous dissipation term λ have rising values. The Nusselt number decreases as the amount of these parameters rises given that they physically enhance the Nano fluids' capacity to boost heat transfer. Also, **tables 2** and **3** characterize the comparison exploration in a limiting manner with the results of [33]. The discrepancy between numerical numbers and values from the literature demonstrates the applicability of the chosen technique whose step size has been set to h = 0.001 with an error tolerance of 10^{-6} and has been meticulously designed for convergence.

5. Conclusion

Here the magnetic dipole in ferromagnetic flow considering nanofluid has been examined numerically. The aluminum oxide nanoparticles utilizing water and ethylene glycol as the base fluid with variable viscosity and thermal conductivity have also studied. The key insight of the mentioned paper is

- The velocity profile was decreased by altering the viscosity parameter, while temperature profile enhanced.
- Temperature raised with higher values of the Curie temperature ratio, and temperature of the fluid decayed with increasing values of the Prandtl number.
- The temperature field enhanced for the rising value of ferromagnetic interaction factor.
- Up surging values of Curie temperature parameter raised the temperature field.
- The increasing λ and λ_1 caused the temperature field to become stronger, and that the nanofluid $\gamma Al_2O_3 H_2O$ has a relatively better influence on this consequence.

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Table 1: List of the current study's (thermo-physical characteristics).

Contents	Density	Specific heat	Thermal conductivity
$Water(H_2O)$	998.3	4182	0.6
Ethylene glycol $(C_2H_6O_3)$	1116.6	2382	0.249
Gamma-alumina (γAl_2O_3)	3970	765	40

Table 2: shows numerical quantities of skin friction coefficient $\frac{1}{2} (Re_x)^{\frac{1}{2}} C_{f_x}$ for two nanofluids when $m_1 = 1, \lambda = 0.1, z = 0.5$. Obtained results are compared with the results of Abbas Khan [33].

Khan et al. [33]			Present			
φ	β	α	$\gamma Al_2O_3 - H_2O$	$\gamma A l_2 O_3 - C_2 H_6 O_3$	$\gamma Al_2O_3 - H_2O$	$\gamma Al_2O_3 - C_2H_6O_3$
0.01	1	0.2	0.529705	0.52471	0.529700	0.52461
0.03			0.897501	0.879547	0.897491	0.879536
0.05			0.54147	0.50264	0.54137	0.50252

2		2.00877	1.96994	2.00863	1.96982
3		2.47608	2.43724	2.47600	2.43711
	0.3	3.68432	3.1646	3.68421	3.1632
	0.4	4.68448	4.14641	4.68436	4.14653

Table 3: Comparison of numerical results reported by of local Nusselt number for two nanofluids when $\varphi = 0.1$, $\alpha = 0.2$, z = 0.5.

Khan et al. [33]					Present		
Pr	β	λ	$\gamma A l_2 O_3 - H_2 O$	$\gamma A l_2 O_3 - C_2 H_6 O_3$	$\gamma Al_2O_3 - H_2O$	$\gamma A l_2 O_3 - C_2 H_6 O_3$	
5.6	1	0.3	1.8123	1.89077	1.8112	1.89068	
6.6			2.29341	2.39237	2.29331	2.39227	
7.6			2.77451	2.89397	2.77462	2.89297	
	2		1.85614	1.93647	1.85602	1.93637	
	3		0.937765	0.978975	0.937763	0.978964	
		0.35	0.47953	0.501218	0.479532	0.501207	
		0.4	0.0212945	0.0234611	0.0212941	0.0234601	

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