

Impact of Viscous Dissipation on MHD Darcy-Forchheimer Nanoliquid Flow Comprising Gyrotactic Microorganisms past a Nonlinear Extending Surface

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Abstract

In the current article, MHD flow problem of Darcy-Forchheimer nanoliquid comprising motile microorganisms with viscous dissipation impact past a nonlinear elongated sheet is addressed. The inclusion of gyrotactic microorganisms in nanoliquid assists to boost up the thermal efficiency in numerous micro-biological systems. Iterative solution of the single-phase flow problem is attained by utilizing the procedure of SOR (Successive over Relaxation). The influences of leading parameters on flow velocity, temperature, density and concentration of motile microbes are considered and depicted through tables and graphs by employing software of MATLAB. Furthermore, comparison table is developed to check the accuracy of numerical results of the flow problem under consideration. An increment in the Forchheimer parameter causes reduction in the velocity distribution. It is portrayed that the Lewis number and Brownian motion parameter tend to enhance the rate of mass transport.

Keywords: Darcy-Forchheimer; Bioconvection; Viscous Dissipation; Nonlinear Stretching Sheet; Gyrotactic Microorganisms.

2010 Mathematics Subject Classification: 76A05; 76M20, 35Q35, 76D05, 80A20

Abbreviations:

u	x –Component of velocity	T_{bn}	Temperature at surface of sheet
v	y –Component of velocity	b_{ch}	Chemotaxis constant
μ	Dynamic viscosity	α	Thermal diffusivity
ν	Kinematic viscosity	Ec	Eckert number
ρ_{Fl}	Fluid's density	B_0	Strength of magnetic field
ρ_{Np}	Nanoparticles' density	n	Microorganisms' concentration
T	Temperature	C	Volume fraction of nanoparticles
c_p	Specific heat	Me	Magnetic parameter
b_o	Positive constant	Pr	Prandtl number
K_*	Permeability of medium	Le	Lewis number
c_B	Drag force coefficient	Sc	Schmidt number
D_{Tp}	Thermophoresis diffusion coefficient	Nb	Parameter of Brownian motion
D_{Bm}	Brownian coefficient	Nt	Parameter of thermophoresis
D_{Gm}	Microorganism diffusion coefficient	Pe	Peclet number
σ_{Ec}	Electric conductivity	λ	Porosity parameter
T_{inf}	Ambient temperature	η	Dimensionless variable
W_{cs}	Maximum swimming speed of cell	Fr	Forchheimer parameter

1. Introduction

The nanoliquid has well known contributions in the fields of medicine, chemical processes, devices of heat transfer, chillers, cooling systems and industry etc. Initially, Choi [1] introduced that combined mixture of micro metallic particles and conventional base fluids form a new kind of liquid, known as nanoliquid. This new type of engineering liquid has more capacity of heat transfer as compared to

conventional base fluids. When Reynolds number is greater than 1 then the flow is nonlinear. In such circumstances, it becomes impossible to ignore impacts of inertia. The non-Darcian model of permeable media is obtained from classical Darcian model that includes inertial drag, vorticity diffusion impacts, inertial drag and their combination as well. In order to demonstrate the impressions of boundary as well as inertia, Forchheimer [2] developed a square velocity term which is added in expression of velocity. Muskat [3] called this term as “Forchheimer term” which is valid for high Reynolds number. Darcy-Forchheimer theory was utilized by Seddeek [4] to investigate mixed convection Darcy-Forchheimer flow with thermophoresis and viscous dissipation impact through porous media. The solution of problem involving flow of Darcy-Forchheimer Maxwell nanoliquid was revealed by Sajid et al. [5] on a stretching linear shallow with support of shooting technique. Sheikholeslami et al. [6-8] numerically interpreted the nanofluid performance in a solar LFR system, heat storage with honeycomb configuration utilizing nanoparticles and solar system equipped with innovative hybrid nanofluid.

Darcy- Forchheimer viscous nanofluid flow past an elongating curved surface was analyzed by Saif et al. [9] through a penetrable region. Nasir et al. [10] investigated flow problem of thin film 2D Darcy-Forchheimer nanofluid past a flat unsteady stretchable surface in a porous media. Problem of micropolar 3D MHD (magnetohydrodynamics) Darcy-Forchheimer nanoliquid CNTs flowing among horizontal & parallel plates was inspected in a porous region by Khan et al. [11]. Ganesh et al. [12] scrutinized MHD Darcy-Forchheimer flow of nanofluid past a sheet in thermally stratified porous medium. They solved the flow problem by employing Runge-Kutta method of order four with shooting technique. Nanoliquid flow with thermophoresis impact and Brownian motion is numerically examined by Hayat et al. [13] under Darcy-Forchheimer relationship. Sadiq and Hayat [14] addressed the problem of MHD Darcy-Forchheimer two-dimensional flow due to a stretched surface.

The boundary-layer MHD flow of Maxwell nanoliquid generated by a stretching surface was investigated by Muhammad et al. [15] in a porous medium under influence of uniform magnetic field. Hayat et al. [16] scrutinized the boundary layer flow problem of Darcy-Forchheimer Williamson nanomaterial on a bidirectional nonlinearly stretching surface in a porous region. They adopted the Homotopy Analysis Method (HAM) to generate the numerical solution for governing nonlinear system of flow problem. Zakaullah et al. [17] computed the numerical solution of Darcy-Forchheimer nanofluid flow problem with partial slip impacts with assistance of shooting methodology. Moreover, they explored that temperature and concentration of nanoparticles have expanding behaviors for larger values of Forchheimer parameter. Turk and Tezer-Sezgin [18] addressed the flow problem of

micropolar nanoliquid inside a square under impact of magnetic field and computed its numerical solution by employing Finite Element Method (FEM). Shafiq et al. [19] inspected special impacts of thermal slip and Arrhenius activation energy in flow of Darcy-Forchheimer nanoliquid due to a stretching surface. Some of the investigations involving Darcy-Forchheimer flow have been published [20-22].

The swimming of gyrotactic microorganisms generates the phenomena of bioconvection. The spectacle of bioconvection boosts up microscale mixing and stability of nanoparticles. Effective distribution of heat is due to stabilized nanoparticles. Shahid et al. [23] numerically inspected the flow of blood comprising motile microbes passing through a stretching and penetrable region. They computed the numerical solution of blood flow problem by employing methodology of Successive Taylor Series Linearization (STSLM). Waqas et al. [24] envisaged Buongiorno's model to investigate flow characteristics of MHD Williamson nanoliquid due to a porous wedge containing gyrotactic microorganisms. The numerical solution of this time-dependent flow problem was achieved by shooting method. In the presence of motile microorganisms, the numerical solution of flow problem involving Maxwell nanofluid with thermal radiation impact was obtained by Sohail and Naz [25] with the assistance of HAM. Some recent relevant studies can be referred to the refs. [26-29].

The present study of Darcy-Forchheimer nanofluid flow has different valuable applications in biomedical sciences and biotechnology. According to our knowledge, no work on the numerical investigation of viscous dissipation in Darcy-Forchheimer nanofluid flow comprising gyrotactic microbes employing SOR (Successive over Relaxation) technique has been done yet. This fact encouraged us to conduct the current study. This paper is concerned with behavior of numerous parameters on the MHD flow of Darcy-Forchheimer nanoliquid containing gyrotactic microorganisms under influence of viscous dissipation. The numerical solutions are clarified with the help of figures and tables of various pertinent parameters involved in the concerned problem through MATLAB.

2. Mathematical Formulation

Let us assume steady viscous 2D incompressible Darcy-Forchheimer flow of MHD nanofluid due to a stretching sheet which behaves nonlinearly. The sheet has motion along x -axis with stretching velocity of form $U_{bn} = b_o x^m$, where m and $b_o > 0$ are constant terms.

The geometry of the modeled flow problem is given below in Fig. 1:

The system of differential equations representing the current flow problem has form:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = v \left(\frac{\partial^2 u}{\partial y^2} \right) - \left(\frac{\sigma_{Ec} B_o^2 x^{m-1}}{\rho_{Fl}} + \frac{v}{K_*} \right) u - \frac{c_B}{\sqrt{K_*}} u^2 \quad (2)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha \frac{\partial^2 T}{\partial y^2} + \frac{(\rho c)_{Np}}{(\rho c)_{Fl}} \times \left\{ D_{Bm} \left(\frac{\partial C}{\partial y} \frac{\partial T}{\partial y} \right) + \frac{D_{Tp}}{T_{Inf}} \left(\frac{\partial T}{\partial y} \right)^2 \right\} + \frac{v}{c_p} \left(\frac{\partial u}{\partial y} \right)^2 \quad (3)$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_{Bm} \frac{\partial^2 C}{\partial y^2} + \frac{D_{Tp}}{T_{Inf}} \frac{\partial^2 T}{\partial y^2} \quad (4)$$

$$u \frac{\partial n}{\partial x} + v \frac{\partial n}{\partial y} + \frac{b_{ch} W_{cs}}{(C_{bn} - C_{Inf})} \left\{ \frac{\partial}{\partial y} \left(n \frac{\partial C}{\partial y} \right) \right\} = D_{Gm} \frac{\partial^2 n}{\partial y^2} \quad (5)$$

The boundary conditions (BCs) related to above equations are:

$$u = U_{bn} = b_o x^m, \quad v = 0, \quad T = T_{bn}, \quad C = C_{bn}, \quad n = n_{bn} \quad \text{at } y=0 \quad (6)$$

$$u \rightarrow 0, \quad T \rightarrow T_{Inf}, \quad C \rightarrow C_{Inf}, \quad n \rightarrow n_{Inf} \quad \text{as } y \rightarrow \infty \quad (7)$$

Let's define dimensionless quantities

$$\eta = \frac{1}{2} \sqrt{\frac{2b_o(m+1)\rho_{Fl}}{\mu}} x^{\frac{m-1}{2}} y, \quad \theta(\eta) = \frac{T - T_{Inf}}{T_{bn} - T_{Inf}}, \quad \chi(\eta) = \frac{n - n_{Inf}}{n_{bn} - n_{Inf}}, \quad (8)$$

$$\phi(\eta) = \frac{C - C_{Inf}}{C_{bn} - C_{Inf}}, \quad u = b_o x^m f'(\eta), \quad v = -\frac{1}{2} \sqrt{2b_o(m+1)\nu} x^{\frac{m-1}{2}}.$$

The velocity components $u = \frac{\partial \psi}{\partial y}$ & $v = -\frac{\partial \psi}{\partial x}$ provide assistance to develop stream function ψ and

η denotes the dimensionless similarity variable. The components u and v help to satisfy Eq. (1) in form of ψ . The transformation specified in (8) can be exploited to convert Eqs. (2) – (5) into ODEs given below:

$$f''' + ff'' - \left[\frac{2m}{m+1} \right] f'^2 - Me^2 f' - \lambda f' - Fr(f')^2 = 0, \quad (9)$$

$$\theta'' + Pr \left[Nb\theta'\phi' + Nt\theta'^2 + \theta'f \right] + Pr Ec f''^2 = 0, \quad (10)$$

$$\phi'' + \frac{Nt}{Nb} \theta'' + Le Pr (f'\phi') = 0, \quad (11)$$

$$\chi'' + Sc(f\chi') - Pe \left[\chi'\phi' + \phi''(\chi + \sigma) \right] = 0, \quad (12)$$

Also, BCs written in (6) as well as (7) have dimensionless form:

$$f(0)=0, \quad f'(0)=1, \quad \theta(0)=1, \quad \phi(0)=1, \quad \chi(0)=1, \quad \text{at } \eta = 0, \quad (13)$$

$$f' = 0, \quad \theta = 0, \quad \phi = 0, \quad \chi = 0, \quad \text{as } \eta \rightarrow \infty \quad (14)$$

Where primes symbolize derivative w.r.t η . Further, Me is magnetic parameter, Le is Lewis number, Nt is thermophoresis parameter, λ is porosity parameter, Ec shows Eckert number, Pr is Prandtl number, Sc is Schmidt number, Nb is parameter of Brownian motion, Pe signifies Peclet number, σ_{Ec} is a constant term, Fr is the Forchheimer parameter, which have mathematical forms as:

$$\begin{aligned} Nt &= \frac{(\rho c)_{Np} D_{Th} (T_{bn} - T_{Inf})}{(\rho c)_{Fl} \nu T_{Inf}}, & Nb &= \frac{(\rho c)_{Np} D_{Bm} (C_{bn} - C_{Inf})}{(\rho c)_{Fl} \nu}, & Pr &= \frac{\nu}{\alpha}, \\ Ec &= \frac{(u_{bn})^2}{c_p (T_{bn} - T_{Inf})}, & Sc &= \frac{\nu}{D_{Gm}}, & Le &= \frac{\nu}{D_{Bm}}, & Me &= \sqrt{\frac{2\sigma_{Ec} B_o^2}{b_o \rho_{Fl} (m+1)}}, \\ Pe &= \frac{b_{ch} W_{cs}}{D_{Gm}}, & \sigma_{Ec} &= \frac{n_{Inf}}{(n_{bn} - n_{Inf})}, & \lambda &= \frac{2\nu}{K_* b_o (m+1) x^{m-1}}, & Fr &= \frac{2c_B x}{K_*^{\frac{1}{2}} (m+1)}. \end{aligned} \quad (15)$$

3. Numerical Computation

Many researchers are contributing in various fields to find numerical solutions of nonlinear modeled problems. In our problem, the numerical elucidation of resulted nonlinear ODEs (9)-(12) involving BCs described in (13)-(14) is accomplished with the assistance of a methodology based on SOR. We have adopted this procedure because it provides the facility of fast convergence. Since, it is complicated to compute analytical solution of nonlinear differential equations of higher order. A detail discussion on our numerical technique can be seen in our prior papers [30-35].

To find numerical solution of coupled system of nonlinear ODEs, we applied Finite Difference Method. As the ODE given in (9) is of order 3, its order is reduced by putting $f' = r$. Finally, the resultant system of equations with their boundary conditions has form:

$$r'' + fr' - \left[\frac{2m}{m+1} \right] r^2 - Me^2 r - \lambda r - Fr(r)^2 = 0, \quad (16)$$

$$\theta'' + Pr \left[Nb\theta'\phi' + Nt\theta'^2 + \theta'f \right] + Pr Ecr'^2 = 0, \quad (17)$$

$$\phi'' + \frac{Nt}{Nb} \theta'' + Le Pr(f\phi') = 0, \quad (18)$$

$$\chi'' + Sc(f\chi') - Pe \left[\chi'\phi' + \phi''(\chi + \sigma_{Ec}) \right] = 0, \quad (19)$$

The boundary conditions get the form:

$$\begin{aligned} f(0) = 0, \quad r(0) = 1, \quad \theta(0) = 1, \quad \phi(0) = 1, \quad \chi(0) = 1, \quad \text{at } \eta = 0, \\ r = 0, \quad \theta = 0, \quad \phi = 0, \quad \chi = 0, \quad \text{as } \eta \rightarrow \infty \end{aligned} \quad (20)$$

Firstly, the domain $[0, \infty)$ of flow problem under observation is discretized for the numerical manipulation of above-mentioned system of equations. In order to integrate the equation $f' = r$, Simpson's numerical iterative process is employed. Additionally, well-known central difference approximation is employed to discretize the involved derivatives in (16)-(20) as well as an iterative solution is found by applying the methodology of SOR. This iterative scheme is terminated when

$$\max \left(\|r_{j+1} - r_j\|_2, \|f_{j+1} - f_j\|_2, \|\theta_{j+1} - \theta_j\|_2, \|\phi_{j+1} - \phi_j\|_2, \|\chi_{j+1} - \chi_j\|_2 \right) < Tol_{irr}$$

is satisfied up to four consecutive iterations.

4. Results and Discussion

In this section, we are committed to present our outcomes through physical elucidations. The relevant flow model dimensionless equations (9)-(12) are numerically treated with the help of SOR method. The velocity $f'(\eta)$, motile microorganisms (microbes) density distribution $\chi(\eta)$, concentration $\phi(\eta)$ and temperature $\theta(\eta)$ display asymptotic appearance against the well-adjusted step size η . The movements of microorganisms and nanoparticles are independent of each other. The motion of nanoparticles is because of the Brownian and thermophoresis effects while bioconvection phenomenon

within the flow causes the motion of the mobile microorganisms. The microorganisms own self-propelled movement unlike the nanoparticles. The impact of the physical parameters on the shear stress, mass transport rate, motile microbes' density number and heat transport rate are also examined and tabularized.

The motile microbes' density, velocity, concentration and temperature profiles are drawn in Figs. 2-5 respectively for multiple values of magnetic field parameter Me . It can be evidently seen from these figures that the effect of magnetic parameter enhances microorganisms' density, temperature and concentration while its effect on velocity shows a decreasing trend. A retarding force is induced in the direction of motion because of the applied magnetic field. This force strictly opposes the motion which causes a reduction in the velocity of fluid. The same profiles are examined for porosity parameter, as appeared in Figs. 6-9. It is obvious from the analysis of these figures that the porosity and magnetic parameters insert the same effects on the aforementioned profiles.

The impacts of local inertia or Forchheimer parameter on the profiles such as $f'(\eta)$, $\chi(\eta)$, $\theta(\eta)$ and $\phi(\eta)$ are portrayed in Figs. 10-13. The velocity profile gets decrease with the influence of Forchheimer parameter while the remaining profiles (microbes' density, concentration and temperature) are increased with an increment in this parameter. Both parameters (porosity and Forchheimer) act parallel to each other in terms of their effects on the dimensionless profiles. The porosity of the medium and the surface drag are proportional to the Forchheimer parameter or inertial coefficient. Due to this reason the resistive force increases which causes an enhancement in surface friction. Consequently, the raising values of Fr cause a decrement in the velocity and increment in the motile microbe's density. A numerical data comparison, to validate our code, is presented in Table 1. The outcomes are correlated with those presented by Rasool et al. [36]. The results are matched in limiting cases for distinct values of Forchheimer parameter. A better correlation of the results in Table 1 assures the accuracy of our numerical solutions.

Tables 2-5 are provided to check out the impacts of the preminent parameters on local motile density number $\chi'(0)$, local Sherwood number $\phi'(0)$, skin friction coefficient $f''(0)$ and local Nusselt number $\theta'(0)$. The variation of skin friction for several parameters is depicted in Table 2. It is noticed here that various values of porosity, magnetic and Forchheimer parameters cause an enhancement in the skin friction. The impacts of parameters like Nt , Nb , λ , Fr , Ec and Me on the Nusselt number can be examined in Table 3. The values of $\theta'(0)$ are diminishing

with the influence of these parameters. The magnetic, porosity and thermophoresis parameters tend to devaluate the mass transport rate on the sheet surface, on the other hand; the Lewis number, Brownian motion parameter and the Forchheimer parameter cause a substantial increase in the Sherwood number (mass transfer rate) as portrayed in Table 4.

The desired consequences can be attained by providing the suitable values to the preeminent parameters. The pores sizes can affect the local density number of microorganisms. If the sizes of pores are so small such that the tiny organisms do not cross them easily then the bioconvection may not occur. In order to avoid this situation, it will be better to take medium having high porosity or much smaller permeability. Table 5 reveals the influences of Me , λ , Pe , Sc and Fr over the motile microorganisms density number $\chi'(0)$. The density of microbes is enhanced with the escalating values of Pe , Sc and Fr . Contrarily, the magnetic and the porosity parameter tend to decline the density of the microbes. In the tabular analysis, it is worth mentioning that the surface drags as well as $\chi'(0)$ is increased on the surface of sheet with the influence of the Forchheimer parameter.

Conclusions:

We have numerically investigated the MHD Darcy-Forchheimer nanofluid flow due to nonlinear elongating sheet in a porous medium. Magnetic effects not only maintain the temperature but also regulate flow. The numerical outcomes are attained via the Successive over Relaxation method. Major findings are specified below:

- It has been elucidated that the magnetic parameter tends to enhance microorganisms' density, temperature and concentration while its impact on flow velocity demonstrates a declining trend.
- An increase in Forchheimer parameter causes a decrement in the velocity and increment in the motile microbe's density.
- It is portrayed that the porosity and thermophoresis parameters tend to devaluate the mass transport rate on the sheet surface, on the other hand; the Lewis number and Brownian motion parameter cause a substantial increase in the Sherwood number (mass transfer rate).
- The Forchheimer parameter tends to escalate the $\chi'(0)$ on the surface of sheet.

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Fig. 1: The geometry of the flow problem

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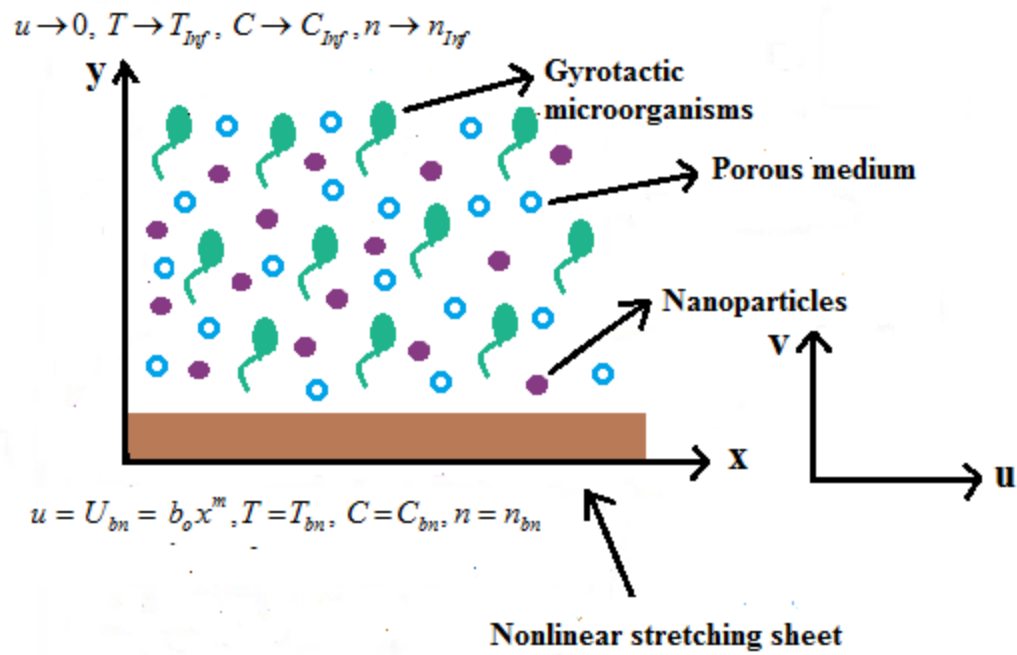


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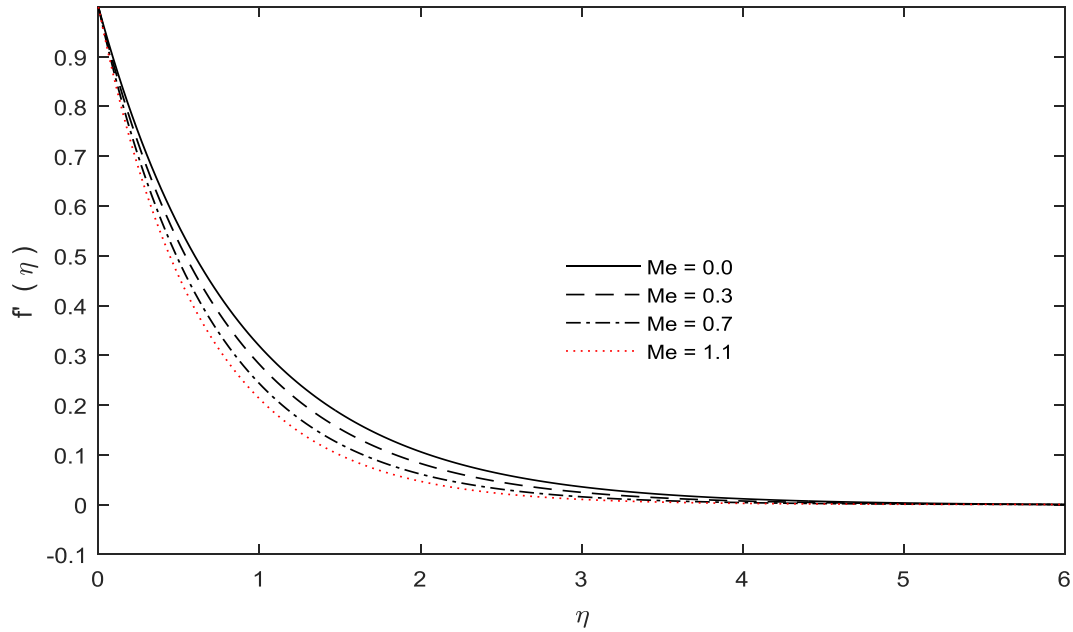


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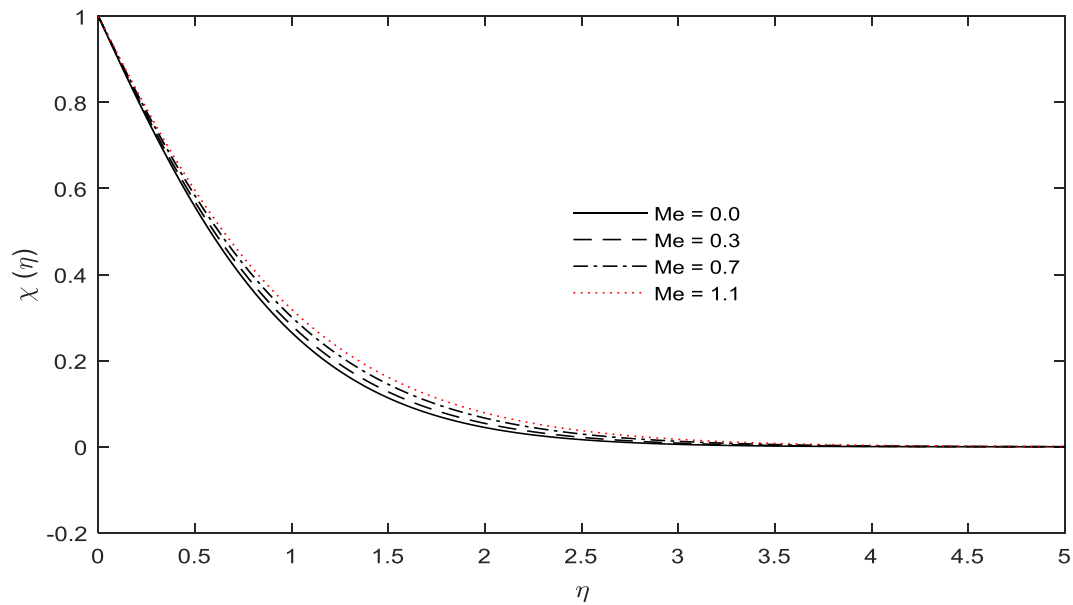


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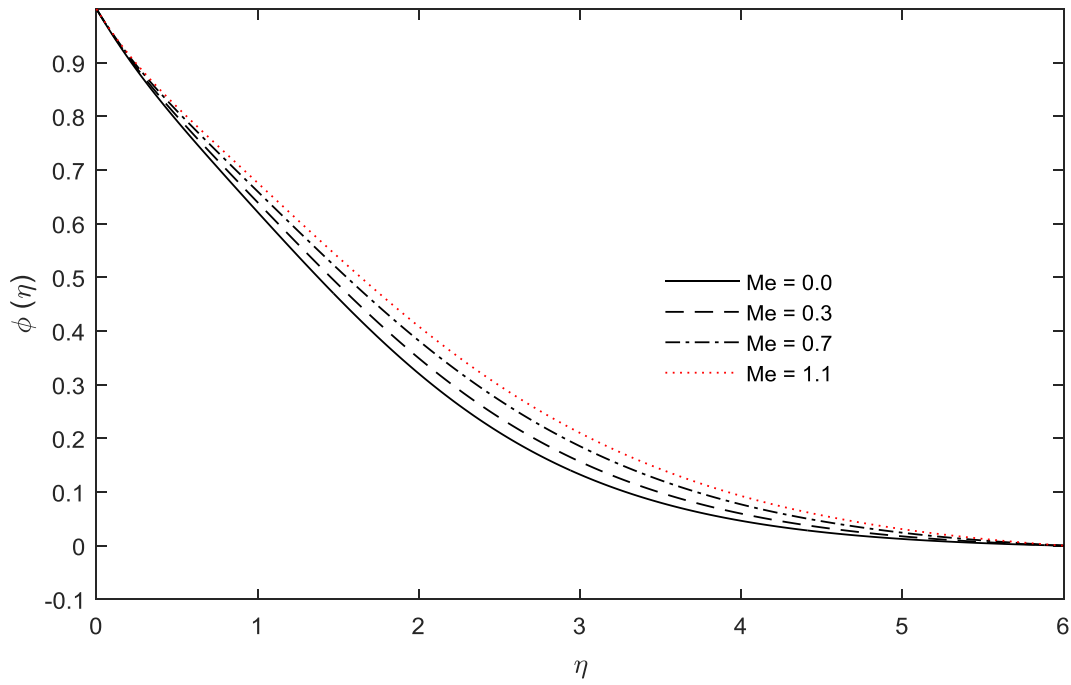


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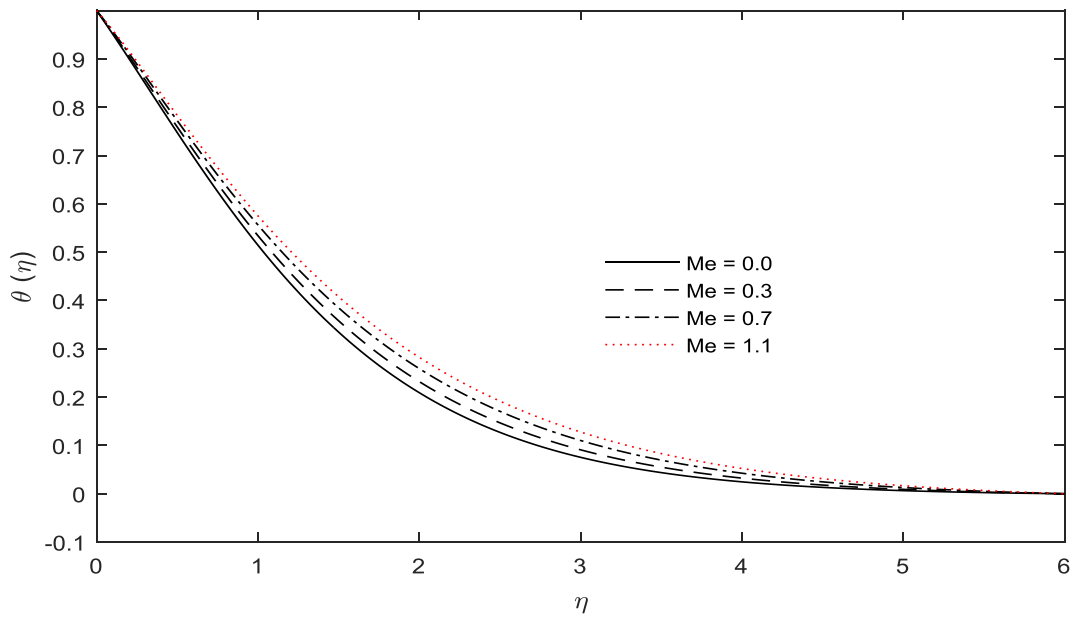


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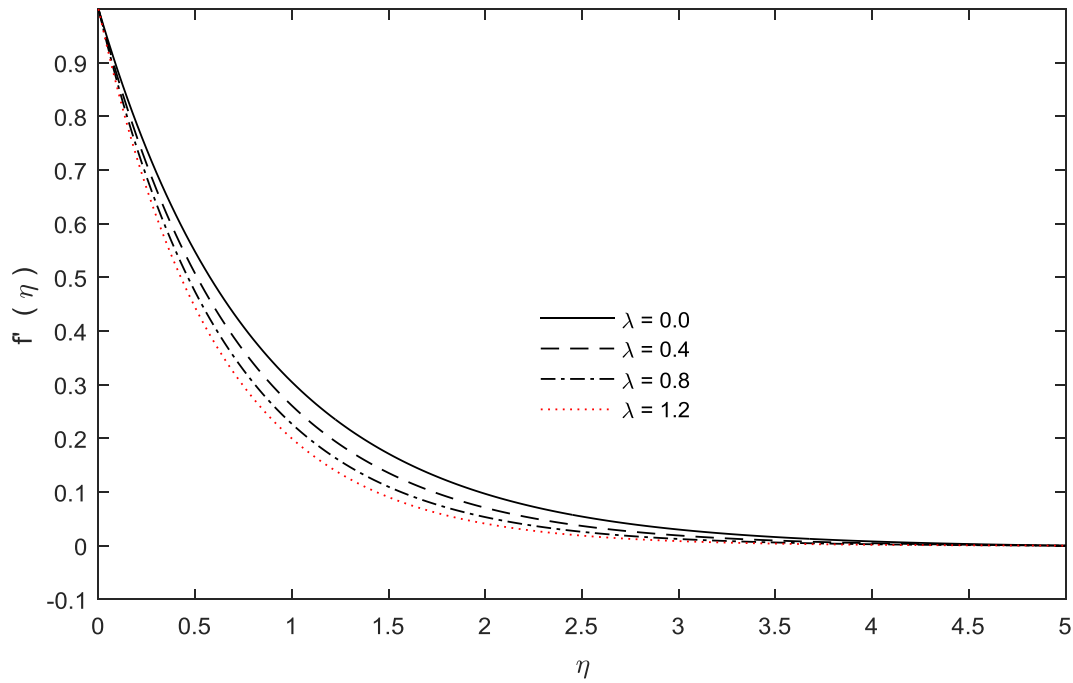


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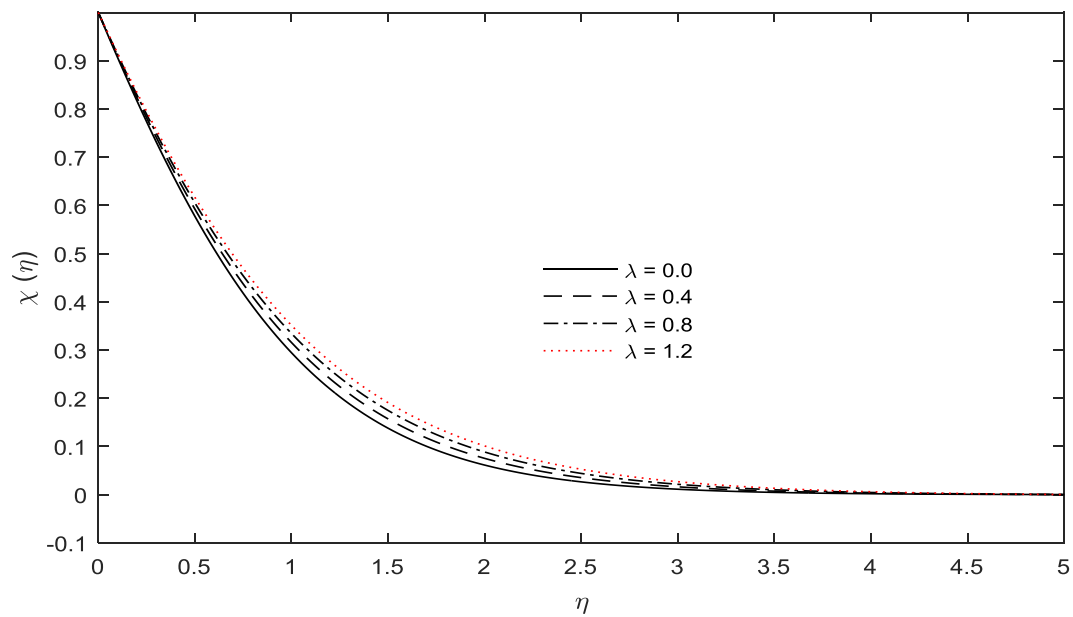


Fig. 7: Microorganisms' density distribution for various values of λ by setting $Me = 0.3, Fr = 0.2, Pr = 2.2, Nb = 0.1, Nt = 0.2, m = 1.3, Ec = 0.1, \sigma_{Ec} = 0.2, Le = 1.9, Pe = 0.1$ & $Sc = 2$.

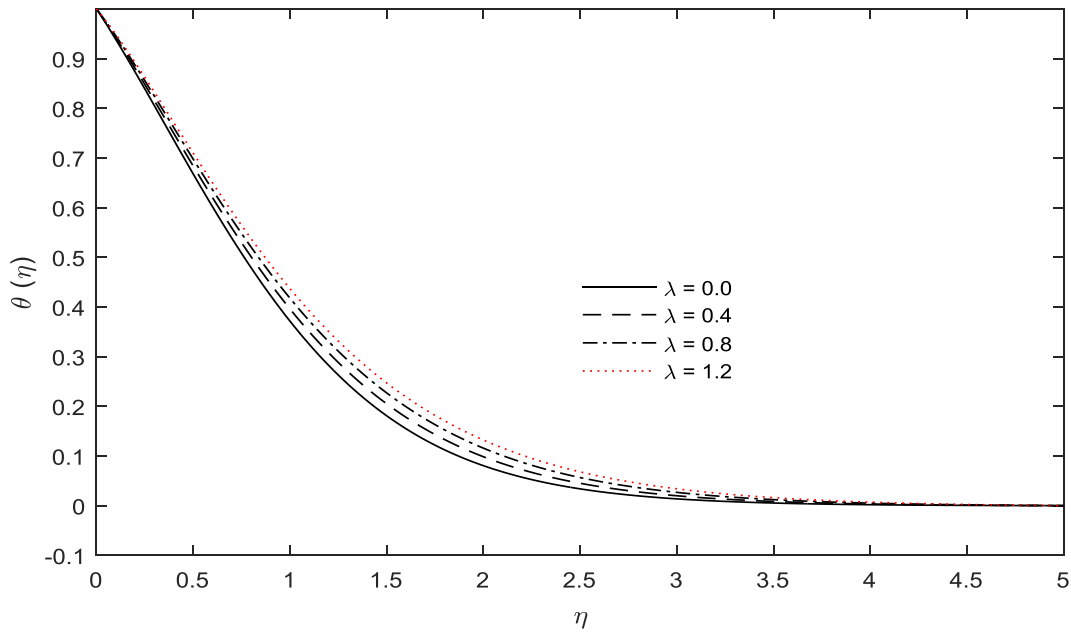


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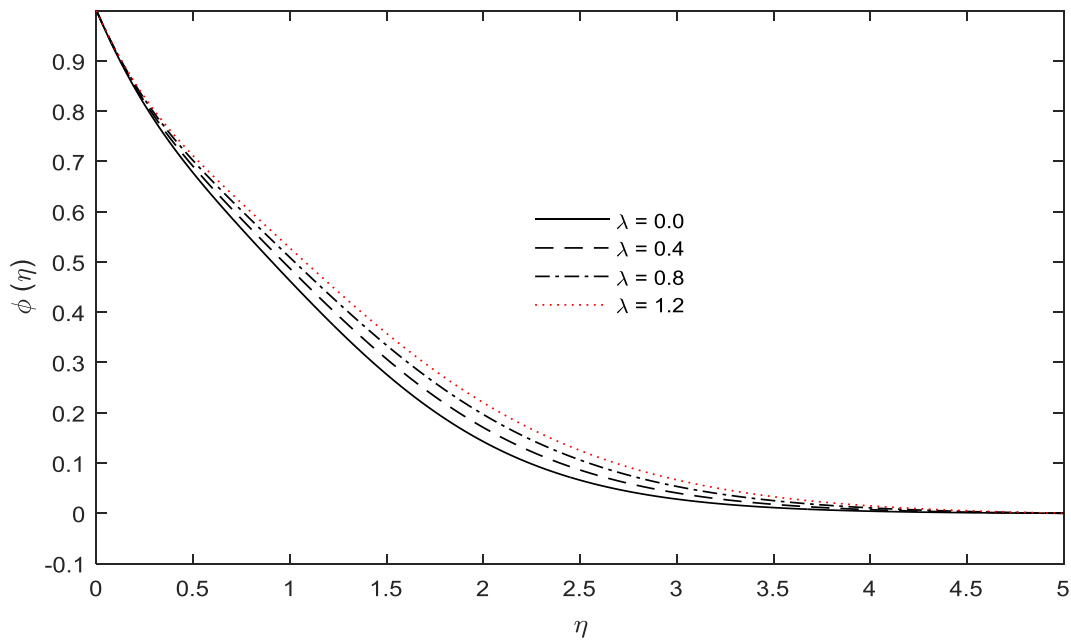


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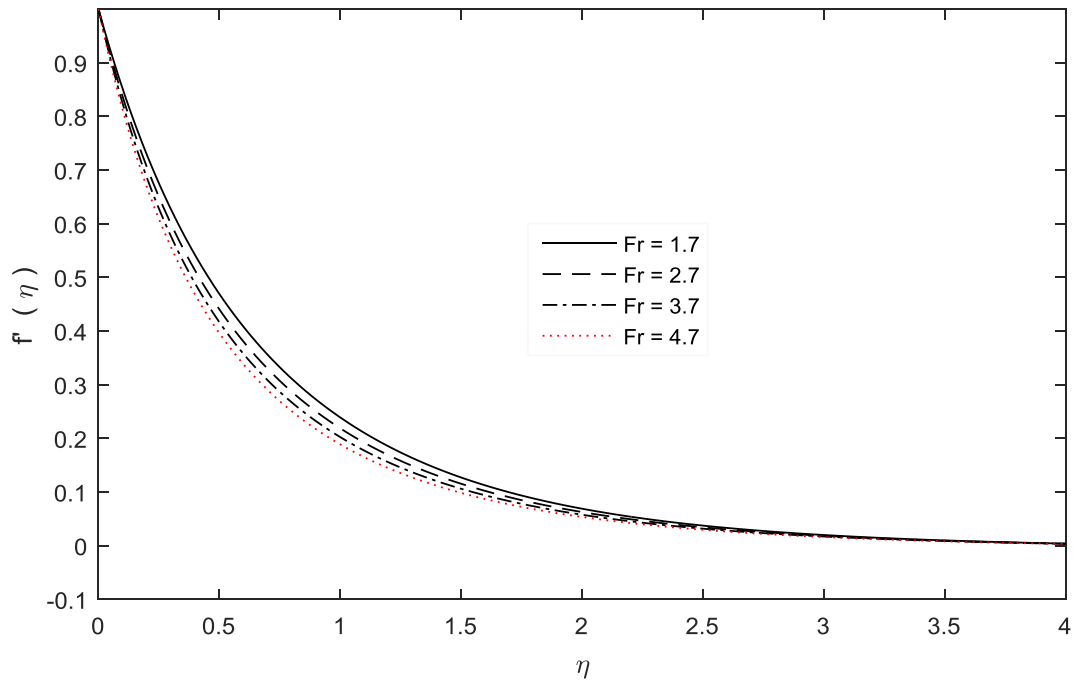


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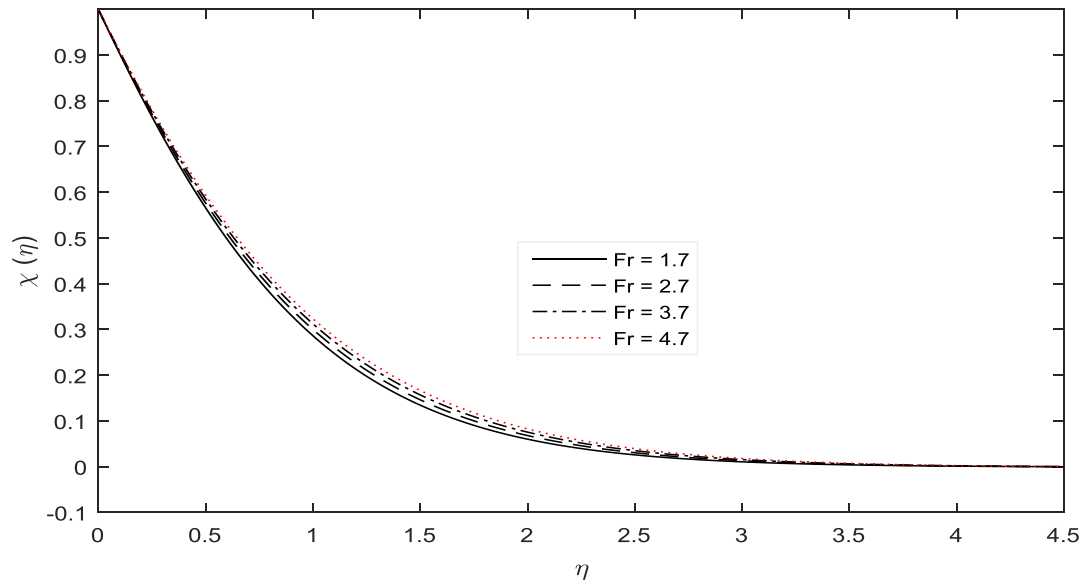


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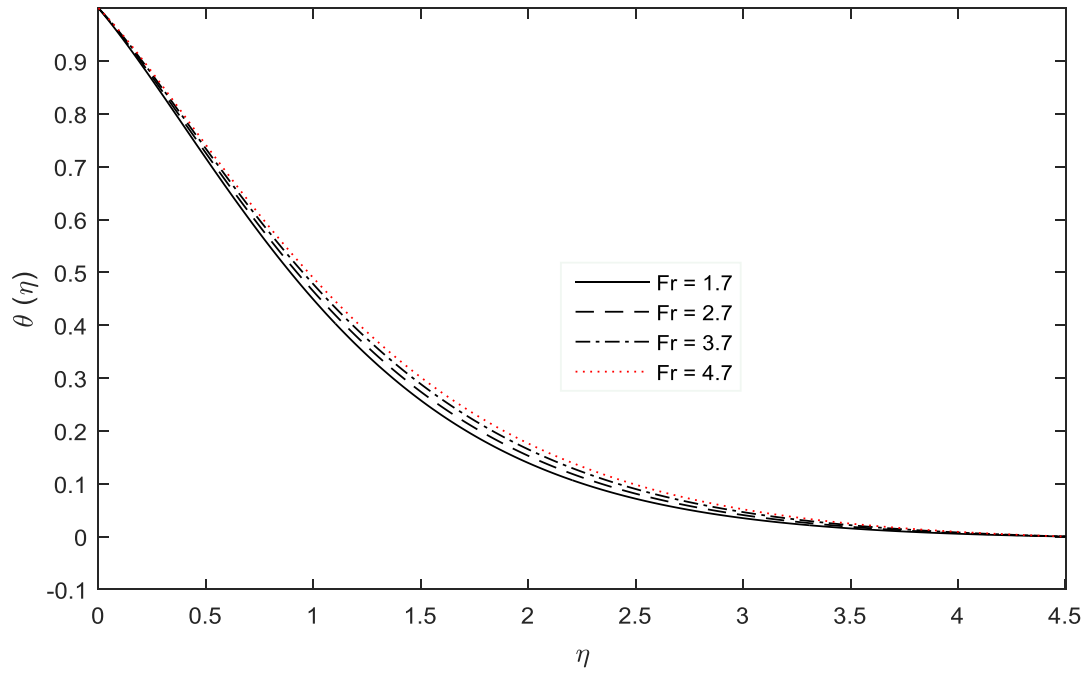


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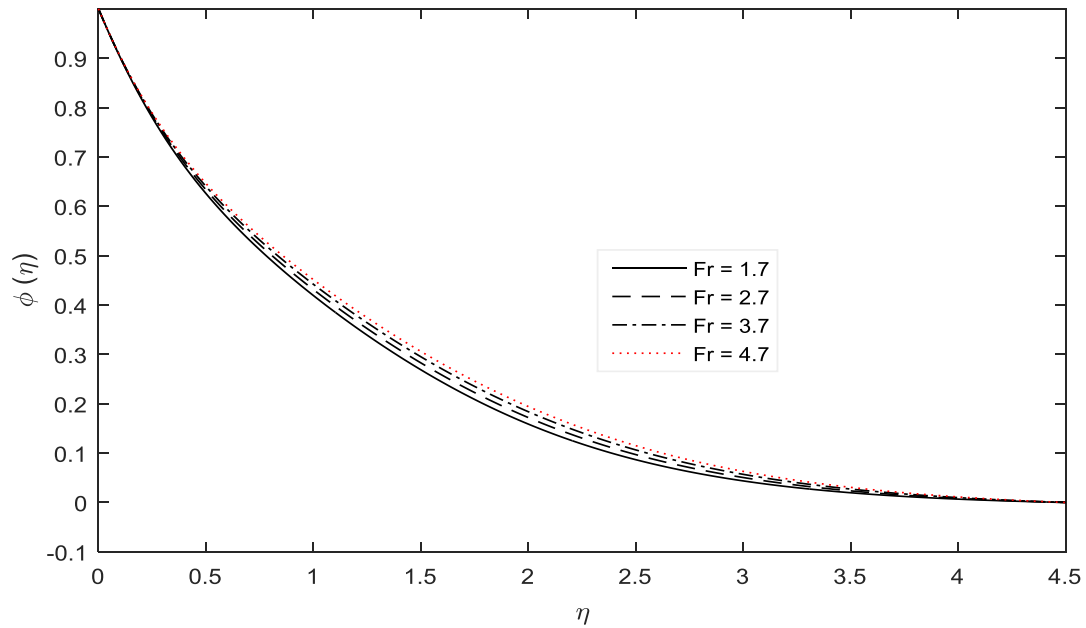


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Table 1. Comparison with Rasool et al. [36] by setting $Ec = 0, \sigma_{Ec} = 0, Pe = 0$ & $Sc = 0$.

Fr	$-\text{Re}_x^{-1/2} Nu_x$		$-\text{Re}_x^{-1/2} Sh_x$	
	Present	Rasool et al. [36]	Present	Rasool et al. [36]
0.0	0.5133	0.5115	0.4088	0.4085
0.6	0.4936	0.4953	0.3767	0.3925
1.2	0.4853	0.4814	0.3728	0.3789

Table 2. Variation in $f''(0)$, for diverse values of λ , Fr and Me .

λ	Me	Fr	$f''(0)$
0.0	0.2	0.9	-1.423938
0.3			-1.53581
0.5			-1.606089
0.2	0.0	0.9	-1.423938
	0.2		-1.499455
	0.4		-1.571352
0.2	0.2	0.0	-1.267705
		0.3	-1.349067
		0.6	-1.426118

Table 3. Variation in $-\theta'(0)$, against various values of Nt , Nb , λ , Fr , Ec and Me .

Nt	Nb	λ	Fr	Ec	Me	$-\theta'(0)$
0.1	0.2	0.2	0.9	0.1	0.2	0.445417
0.2						0.415899
0.4						0.385252
0.1	0.1	0.2	0.9	0.1	0.2	0.467118
	0.2					0.432433
	0.4					0.391482
0.1	0.2	0.0	0.9	0.1	0.2	0.456385
		0.3				0.426793
		0.5				0.414351
0.1	0.2	0.2	0.0	0.1	0.2	0.466133
			0.3			0.446389
			0.6			0.437961
0.1	0.2	0.2	0.9	0.0	0.2	0.499539
				0.1		0.431806
				0.2		0.376081
0.1	0.2	0.2	0.9	0.1	0.0	0.456385
					0.2	0.432169
					0.4	0.419783

Table 4. Variation in $-\phi'(0)$, for diverse values of Me , λ , Nb , Nt , Le and Fr .

Me	λ	Le	Nb	Nt	Fr	$-\phi'(0)$
0.0	0.2	0.8	0.2	0.1	0.9	0.430925
0.2						0.401441
0.4						0.397490
0.2	0.0	0.8	0.2	0.1	0.9	0.430925
	0.3					0.400779
	0.5					0.396747
0.2	0.2	0.3	0.2	0.1	0.9	0.307160
		0.5				0.311506
		0.7				0.363977
0.2	0.2	0.8	0.1	0.1	0.9	0.328577
			0.2			0.392125
			0.4			0.447677
0.1	0.2	0.8	0.2	0.1	0.9	0.429622
				0.2		0.335406
				0.4		0.237252
0.2	0.2	0.8	0.7	0.1	0.0	0.431477
					0.3	0.402987
					0.6	0.399828

Table 5. Variation in $-\chi'(0)$, for diverse values of Me , λ , Pe , Sc and Fr .

Me	λ	Pe	Sc	Fr	$-\chi'(0)$
0.0	0.2	0.2	0.3	0.9	0.455078
0.2					0.416871
0.4					0.410072
0.2	0.0	0.2	0.3	0.9	0.455078
	0.3				0.415472
	0.5				0.408550
0.2	0.2	0.1	0.3	0.9	0.422488
		0.3			0.442529
		0.5			0.498732
0.2	0.2	0.2	0.1	0.9	0.398371
			0.3		0.411313
			0.5		0.470771
0.2	0.2	0.2	0.3	0.0	0.457298
				0.3	0.420555
				0.6	0.415024

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