



# Innovative theory for the compliance computation in rotors

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## KEYWORDS

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**Abstract.** The conventional energy release rate approach to the computation of the crack compliance in rotors is critically reviewed and its shortcoming is highlighted. The state functions and the generic compliance are introduced, defined, and verified. The proposed functions are defined based on the end states' conditions and are exact. The crack compliance of the shaft in torsion is explicitly defined as the function of the crack depth ratio and recommended as a good alternative to the classical procedure of the energy release rate method. The accuracy and efficiency of the work are verified by concise mathematical formulation and comparison of the results of the work with the others in four examples.

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## 1. Introduction

The fracture stress for brittle materials is related to the crack size, first proposed by Griffith in 1921 [1]. He defined the energy release rate as a function of the stress, the crack depth, and the modulus of elasticity. The energy release rate is defined as the amount of energy per unit area of the crack surface. The Griffith theory is modified for ductile materials by Irwin in 1957 [2]. Rice [3-5] considered the plasticity of the crack tip and extended the method of energy release rate to nonlinear materials. He derived the energy release rate as a path-independent integral, then called the J-integral. Wells [6] proposed the Crack Tip Opening Displacement, CTOD. In 1981, Shih [7] showed a unique relationship between the J-integral and the CTOD. Thereafter, many experiments were conducted for verification of the fracture mechanics models. During the last four decades, a great deal of attention has been paid to the effect of cracks

on rotating machinery [8-17]. The most significant challenge in this area is modeling of the crack. The combination of the theory of energy release rate and the rotor dynamics was a base for the work in classical fracture mechanics in the context of the energy release approach [17], which has inherent difficulty. More information from these references is included in the next section. Qian and Fatemi [18] reviewed the criteria for crack growth problems under mixed mode loadings. The various parameters and criteria proposed in the literature for prediction of the rate and directions of crack growth and the limitations and physical basis for each criterion are reviewed. Several loading conditions are considered and the effects of important variables, such as load magnitudes, material strength, initial crack tip condition, mean stress, and crack closure, on mixed mode crack growth directions and/or rates are discussed. Lin and Smith proposed that a linear elastic finite-element analysis is used to estimate the stress intensity factor, and then the Paris law is employed to calculate crack advances at a few points along the crack front [19]. The results are then used to predict the shape of crack front via a numerical procedure. Carpinteri et al. [20] considered a circular arc circumferential notch in a round bar, and the stress concentration

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factor for both bending and tension is computed. An elliptical arc surface flaw is assumed at the notch root; for different values of stress concentration factor, the stress intensity factor is computed and the effect of the stress concentration on the stress intensity factor is discussed. Li et al. [21] proposed a simple method for calculating stress intensity factors of transverse cracked shaft subjected to shear, tension, and bending. The normalized energy density and the critical load are calculated. The initial crack growth angles are determined, and their relation with the mixed mode ratio is investigated. Fatigue behavior of lateral notched round bars made of high-strength steel under bending, torsion, and combined bending torsion was studied by Branco et al. [22]. The notch effect was analyzed using the equivalent strain energy density concept and the fatigue life was predicted via the Coffin-Manson model. They obtained a close correlation between experimental and predicted fatigue lives. The present paper is an effort to remove the difficulty. A new and exact theory for the determination of the compliance of cracked shaft as an explicit function of the crack depth ratio based on the original concepts of structural and fracture mechanics and concise mathematical logics is proposed in this paper. Only the crack compliance of round bars is included, and other issues of fracture mechanics and fatigue analysis are excluded.

## 2. Review of the energy release rate theory

The key concepts of the classical fracture mechanics are reviewed in this section. Crack compliance,  $c_S$ , is computed due to force component,  $P$ , by applying the Castigliano's theorem as follows:

$$c_S = \frac{d^2 E_C}{dP^2}, \quad (1)$$

where crack energy release,  $E_C$ , is defined as:

$$E_C = \int_0^\xi G A dx, \quad (2)$$

in which  $G$  is the energy release rate (energy per unit area of crack surface) and  $\xi = a/h$  is the crack depth ratio where  $a$  is the crack depth and  $h$  is the structure depth, see Figure 1. The energy release rate is defined in terms of stress intensity factor,  $K$ , and elastic modulus,  $E$ , as follows:

$$GE = K^2. \quad (3)$$

The stress intensity factor depends on the cracked section geometry and the applied load. For a typical specimen, the stress intensity factor is defined as:

$$K = \int_0^a \sigma(x) W(x, a) dx, \quad (4)$$

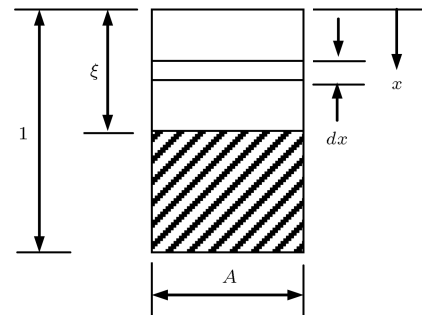


Figure 1. Crack surface area.

in which  $\sigma(x)$  is the internal stress distribution,  $W(x, a)$  is the weight function [23], and  $x$  is position in the crack surface. A typical weight function used for one-dimensional mode I crack was proposed by Shen and Glinka [24] Glinka and Shen [25] as follows:

$$W(x, a) = \frac{2F}{\sqrt{2\pi(a-x)}} \left( 1 + A_1(1-x/a)^{1/2} + A_2(1-x/a)^1 + A_3(1-x/a)^{3/2} \right), \quad (5)$$

in which splitting force  $F$  and parameters  $A_1$ ,  $A_2$ , and  $A_3$  are to be computed as in [24]. The difficulty in using the weight function method is its definition and the accurate computation of the integral in Eq. (4). Stress field,  $\sigma(x)$ , can be very complex and nonlinear and the analytical integration is seldom easy in practice. Also, the standard numerical integration methods are inaccurate due to the singularity of the weight function at the crack tip. It has been well known since 1983 [8] that the local flexibility of a cracked shaft can be calculated in terms of the strain energy release rate function over the crack surface area as in Eq. (1). The process commences with the calculation of the stress intensity factor in Eq. (4) followed by calculation of the energy release rate in Eq. (3) and ends in determination of the compliance in Eq. (1). With reference to Papadopoulos [16,17], since the stress intensity factors for circular cross-sections are not available, the analysis of the computation of compliance of the cracked shaft is done by the energy release rate method as follows. The non-dimensional form for torsional compliance,  $c_T$ , is computed (see Figure 2) as follows:

$$c_T = \int_{-b}^b dx \int_0^a \frac{32}{\pi} \left( 1 - \left( \frac{x}{R} \right)^2 \right) \frac{y}{R} F^2(x, y) dy, \quad (6)$$

in which  $F$  is defined as:

$$F(x, y) = \sqrt{\frac{\tan(\pi y/2h_x)}{(\pi y/2h_x)}} \times \frac{[0.923 + 0.199(1 - \sin(\pi y/2h_x))^4]}{\cos(\pi y/2h_x)}, \quad (7)$$

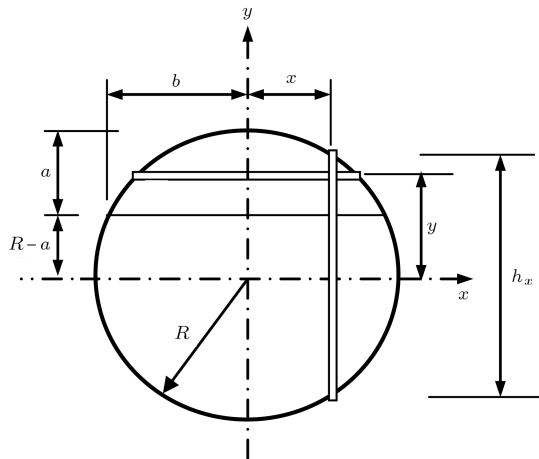


Figure 2. Geometry of cracked shaft.

and:

$$b = \sqrt{R^2 - (R - a)^2}, \quad (8)$$

and:

$$h_x = 2\sqrt{R^2 - x^2}. \quad (9)$$

The shaft is considered to be the sum of independent strips and the compliance is obtained by integrating along the crack tip. In this way, there is difficulty in using the double integral in Eq. (6), because the functions present a singularity when  $x$  is equal to  $b$ . If the crack depth exceeds the radius of the shaft, the compliance presents a divergence. Abraham et al. [26] as well as Dimarogonas [27] suggested us to select integration length ratio,  $\eta$ , to make the integration near but not on the boundaries  $-b_m$  and  $b_m$  defined as follows:

$$b_m = \eta b, \quad 0 < \eta < 1. \quad (10)$$

A  $\eta$  value in range  $0.9 \leq \eta \leq 0.95$  should be used in order to avoid the singularity and estimate a realistic value of the compliance for deep cracks. The main aim of this paper is to remedy the problem by proposing a new theory as follows.

### 3. Basic principles

The governing equation for the free vibration of the shaft in torsion is defined as follows:

$$\theta^{(2)} + \lambda_\theta^2 \theta = 0, \quad (11)$$

where  $\theta$  is the displacement, the number in round brackets denotes the order of derivative with respect to  $x$  in the super script, the number without round brackets denotes the power, and torsion parameter,  $\lambda_\theta$ , is defined as follows:

$$J\mu\lambda_\theta^2 = I_m\omega^2, \quad (12)$$

where  $\mu$  is the shear modulus,  $J$  is the section polar moment of inertia,  $I_m$  is the mass moment of inertia, and  $\omega$  is the natural frequency. In fracture mechanics, the effect of crack on shaft is defined as:

$$\Delta\theta = c_T\theta^{(1)}, \quad (13)$$

in which  $\Delta$  denotes the change, and  $c_T \in [0, \infty]$  is the crack compliance of the shaft. Eq. (13) may be defined [28] in terms of the new crack displacement  $\theta_c$  as follows:

$$\theta_c = c_T\theta^{(1)}H(x - x_i), \quad (14)$$

where  $H(x - x_i)$  is the Heaviside unit step function with center at crack position  $x_i$ . The derivative of Eq. (13) with respect to  $x$ , the golden derivative, is written as follows:

$$\theta_c^{(1)} = c_T\theta^{(1)}\delta(x - x_i), \quad (15)$$

where  $\delta(x - x_i)$  is the Dirac delta. Eqs. (11) and (15) are combined to obtain the governing equation for the free vibration of the cracked shaft as:

$$\left(\theta^{(1)} - c_T\theta^{(1)}\delta(x - x_i)\right)^{(1)} + \lambda_\theta^2\theta = 0. \quad (16)$$

The validity of Eq. (16) is verified via its similarity with the governing equation for the free vibration of axial bar [28]. To derive the finite-element equation for Eq. (16), the conventional method of weighted residual is used. The weighted error is written as follows:

$$\int_0^L \psi \left[ \left(\theta^{(1)} - c_T\theta^{(1)}\delta(x - x_i)\right)^{(1)} + \lambda_\theta^2\theta \right] dx = 0, \quad (17)$$

where  $\psi$  is the weight function. The chain rule of mathematics is used to write the weak form equation as follows:

$$\int_0^L \psi^{(1)}\theta^{(1)}dx - \int_0^L \psi^{(1)}c_T\delta(x - x_i)\theta^{(1)}dx - \lambda_\theta^2 \int_0^L \psi\theta dx = 0. \quad (18)$$

There was no proper chain rule for the second term; hence, the rule of the others is used. The weight function is defined in terms of nodal value,  $\psi_\alpha$ , and shape function,  $N_\alpha(x)$ , as follows:

$$\psi = \psi_\alpha N_\alpha(x). \quad (19)$$

The displacement is defined in terms of nodal value,  $\theta_\beta$ , and shape function,  $N_\beta(N)$ , as follows:

$$\theta = \theta_\beta N_\beta(x), \quad (20)$$

where the repeated indices denote summation and are dummy. Substitution of Eqs. (19) and (20) into Eq. (18) led to the finite-element equation as:

$$(k_{\alpha\beta} - k_{\alpha\beta}^{co} - \lambda_\theta^2 m_{\alpha\beta}) \theta_\beta = 0, \quad (21)$$

where stiffness matrix,  $k_{\alpha\beta}$ , is defined as:

$$k_{\alpha\beta} = \int_0^L N_\alpha^{(1)} N_\beta^{(1)} dx, \quad (22)$$

and mass matrix,  $m_{\alpha\beta}$ , is defined as:

$$m_{\alpha\beta} = \int_0^L N_\alpha N_\beta dx. \quad (23)$$

Original crack equivalent stiffness matrix,  $k_{\alpha\beta}^{co}$ , is defined as:

$$k_{\alpha\beta}^{co} = \int_0^L N_\alpha^{(1)} c_T \delta(x - x_i) N_\beta^{(1)} dx. \quad (24)$$

Eq. (21) is implemented by a personal software (The software determines the Eigen values, frequencies, of Eq. (21).) and is used for the free vibration analysis of a cracked shaft with a given closed-form solution. The finite-element analysis has failed. Close investigation revealed that the equivalent crack stiffness in Eq. (24) and its corresponding term in the weak form of Eq. (18) are not correct and should be modified. The modification is done [29,30] via the first stiffness reduction function,  $F_R \in [0,1]$ , as:

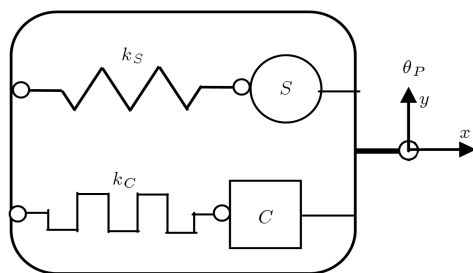
$$k_{\alpha\beta}^c = F_R \int_0^L N_\alpha^{(1)} c_T \delta(x - x_i) N_\beta^{(1)} dx. \quad (25)$$

For the definition of  $F_R$ , the shaft is modeled by a one-degree-of-freedom element (core model), see Figure 3. For this model, Eq. (20) is applied, but now without summation convention. The derivative of Eq. (20) with respect to  $x$  is then defined as:

$$\theta^{(1)} = \theta_\beta N_\beta^{(1)}(x). \quad (26)$$

Divide both sides of Eq. (26) by  $\theta_\beta$  as:

$$\frac{\theta^{(1)}}{\theta_\beta} = N_\beta^{(1)}(x). \quad (27)$$



**Figure 3.** Single degree of freedom core model.

The left side of Eq. (27) is the force per unit displacement, i.e. the stiffness, and then stiffness  $k_S$  is equal to the right-side term as:

$$k_S = N_\beta^{(1)}(x). \quad (28)$$

By substitution of Eq. (28) into Eq. (25) and making use of the property of the Dirac delta, the crack equivalent stiffness in Eq. (25) is written as:

$$k_{\alpha\beta}^c = S_R k_S, \quad (29)$$

where the second stiffness reduction coefficient,  $S_R \in [0,1]$ , is defined as:

$$S_R = c_T F_R k_S. \quad (30)$$

The stiffness of the cracked shaft should be non-negative. This is written as:

$$k_S - S_R k_S \geq 0, \quad (31)$$

or:

$$S_R \leq 1. \quad (32)$$

In view of Eq. (30),  $S_R$  should satisfy:

$$S_R = 0 \quad \text{at} \quad c_T = 0, \quad (33)$$

and:

$$S_R = 1 \quad \text{at} \quad c_T = \infty, \quad (34)$$

$S_R$  with the previous properties is defined as:

$$S_R = \frac{c_T k_S}{1 + c_T k_S}. \quad (35)$$

$F_R$  is defined as follows:

$$F_R = \frac{1}{1 + c_T k_S}. \quad (36)$$

It is interesting to note that the values of these functions are bounded in  $[0,1]$  interval. Reduction of the range of computation from infinity for the compliance to one for the proposed functions is a positive sign for further investigation. The stiffness reduction functions are customized for the unit stiffness and are called the state functions. The first state function,  $F_g$ , is defined as:

$$F_g = \frac{1}{1 + c_g}, \quad (37)$$

and the second state function,  $S_g$ , is defined as:

$$S_g = \frac{c_g}{1 + c_g}. \quad (38)$$

The generic compliance,  $c_g$ , is defined as:

$$c_g = \frac{S_g}{F_g}. \quad (39)$$

The state functions are defined for the unit stiffness; consequently, they are free functions, i.e. they are not defined for a specific structure. This property recommended the authors to look for a free definition of the state functions. The result is as described in the next section.

#### 4. State functions

Experience in the analysis of cracked members led to the need of the definition of the state functions. To that end, it is assumed that there exists an entity in its first state that should be transited smoothly to its second state. This transition happens through the change of state variable,  $\xi$ , which is equal to 0 at the first state and is equal to 1 at the second state,  $\xi \in [0, 1]$ .  $F_g$  and  $S_g$  are proposed for defining the smooth transition between the two states. Based on this definition, the functions should satisfy the following essential end conditions:

$$F_g(\xi) = \begin{cases} 1 & \text{at } \xi = 0 \\ 0 & \text{at } \xi = 1 \end{cases} \quad (40)$$

and:

$$S_g(\xi) = \begin{cases} 0 & \text{at } \xi = 0 \\ 1 & \text{at } \xi = 1 \end{cases} \quad (41)$$

The transition should be isolated from outside of the state variable domain. This requirement is defined by the fresh (natural) end conditions as:

$$F'_g(\xi) = \begin{cases} 0 & \text{at } \xi = 0 \\ 0 & \text{at } \xi = 1 \end{cases} \quad (42)$$

and:

$$S'_g(\xi) = \begin{cases} 0 & \text{at } \xi = 0 \\ 0 & \text{at } \xi = 1 \end{cases} \quad (43)$$

where the prime sign denotes the derivative with respect to  $\xi$ .  $F_g$  with the previous end conditions is defined as:

$$F_g(\xi) = (1 - \xi)^2(1 + 2\xi), \quad (44)$$

and  $S_g$  is defined as follows:

$$S_g(\xi) = \xi^2(1 + 2(1 - \xi)). \quad (45)$$

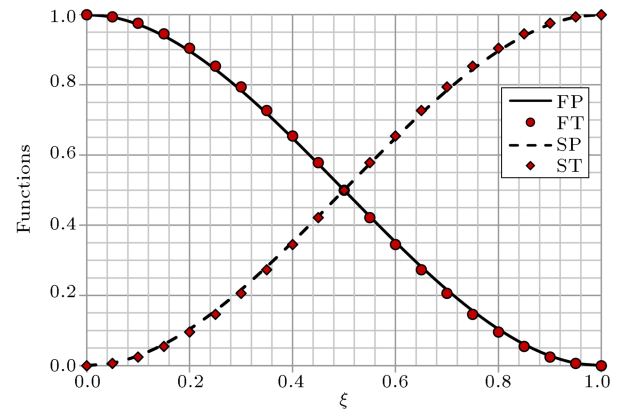
Alternative forms of the state functions are defined in terms of the trigonometric functions as:

$$2F_g(\xi) = 1 + \cos(\pi\xi), \quad (46)$$

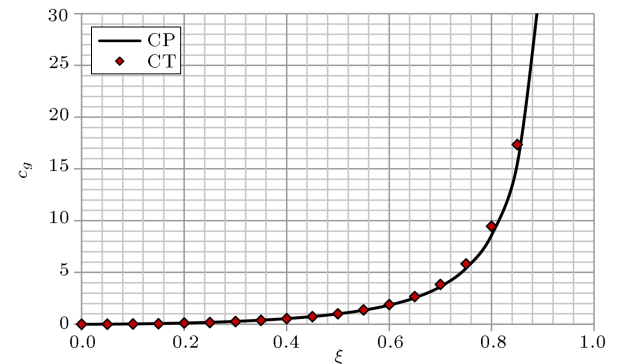
and:

$$2S_g(\xi) = 1 - \cos(\pi\xi). \quad (47)$$

Polynomial state functions FP and SP as well as trigonometric state functions FT and ST are shown in Figure 4. The state functions are quite general, but they are applied to fracture problems where the crack depth ratio is selected as the state variable



**Figure 4.** Polynomial and trigonometric state functions compared.



**Figure 5.** Generic crack compliance.

here. Finally, via Eq. (39), the generic compliance is explicitly defined by the ratio of the two state functions. The generic compliance based on polynomial functions and trigonometric functions, CP, is shown in Figure 5.

The left side of Eq. (27) is written as:

$$k_S = k_m \theta^{(1)}(x_a). \quad (48)$$

In structural mechanics, shaft stiffness,  $k_S$ , is defined as:

$$k_S = k_m / L. \quad (49)$$

Set Eq. (49) equal to the right side of Eq. (27) to obtain stiffness coefficient,  $k_m$ , the non-dimension stiffness, in terms of element length,  $L$ , as:

$$k_m = LN_{\beta}^{(1)}(x_{\beta}). \quad (50)$$

For linear shape function  $k_m = 1$ , and hence the stiffness for the shaft is defined as:

$$k_S = k_m \theta^{(1)}(x). \quad (51)$$

#### 5. Crack compliance

The state functions derived from the previous sections are innovatively used to derive the crack compliance of the rotor shaft in this section.

In fracture mechanics, the crack is defined by a spring with compliance  $c_T$  which produces a jump in the displacement as in Eq. (13). The right sides of Eqs. (13) and (51) are similar, and then crack stiffness (the reduction in the structure stiffness),  $k_C$ , is defined as:

$$k_C = c_m \theta^{(1)}, \quad (52)$$

in which  $c_m$  is the crack compliance. Divide Eq. (52) by Eq. (51) to obtain:

$$\frac{k_C}{k_S} = \frac{c_m}{k_m}. \quad (53)$$

The ratio in Eq. (53) is independent of the geometrical properties. It is valid for all stiffness coefficients. In view of the definition of the generic compliance, Eq. (53) is written as follows:

$$c_m = k_m c_g. \quad (54)$$

Eq. (54) applies to cracked members under any load component. For the torsional member (shaft),  $k_m = 1$ ; consequently, the crack compliance of the cracked shaft in torsion,  $c_T$ , is defined as:

$$c_T = c_g. \quad (55)$$

In spite of the conventional fracture mechanics, the crack compliance of the cracked shaft is an explicit function of the crack depth ratio. Eq. (55) introduced the great saving in the analysis of cracked shafts, because it is a replacement for the energy release rate method. That is a great achievement which increases the accuracy and removes the singularity and other systematic problems. The state functions are defined for gradual and smooth change between the intact and cracked states without any special assumption of hindering their accuracy and generality.

Consequently, the compliance in Eq. (55) is exact. The crack compliance of the cracked shaft is determined independent of the energy release rate. Eq. (55) shows that the compliance is a geometrical property of the cracked shaft, and hence independent of stress, load, energy release or any other external effects. As will be shown in the next section, the energy release rate and the stress intensity factor are defined as functions of the generic compliance. To see the significance of the achievement of the work, the reader is referred to reference [17] for comparison.

## 6. Parameters of fracture mechanics

Structure energy,  $E_S$ , of the core model in Figure 3 is defined as:

$$E_S = 0.5 k_m \theta^{(1)2} \theta_P^2, \quad (56)$$

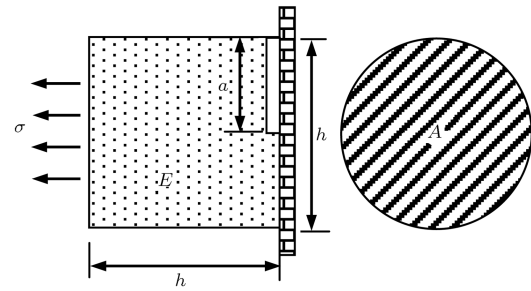


Figure 6. Affected region.

and crack energy release,  $E_C$ , is defined as:

$$E_C = 0.5 c_m \theta^{(1)2} \theta_P^2, \quad (57)$$

where  $\theta_P$  is the core displacement. The ratio of Eqs. (57) and (56) is written as follows:

$$\frac{E_C}{E_S} = \frac{c_m}{k_m}, \quad (58)$$

in view of Eqs. (54) and (58),  $E_C$  is defined as:

$$E_C = E_S c_g. \quad (59)$$

The structure energy for the affected region of the structure, Figure 6, is defined as:

$$E_S = \frac{\sigma^2}{2E} \times Ah, \quad (60)$$

in which  $\sigma$  is the stress. Substitution of Eq. (60) into the derivative of Eq. (59) with respect to  $\xi$  leads to:

$$E'_C = \frac{\sigma^2 Ah c'_g}{2E}. \quad (61)$$

In classical fracture mechanics, the derivative of  $E_C$  is defined as:

$$E'_C = GA. \quad (62)$$

The derivatives in Eqs. (61) and (62) are set equal to each other to obtain energy release rate,  $G$ , for the core model as follows:

$$G = \frac{\sigma^2 h c'_g}{2E}, \quad (63)$$

where  $c'_g$  is explicitly defined as follows:

$$c'_g = \frac{6\xi}{(1-\xi)^3(1+2\xi)^2}. \quad (64)$$

Stress intensity factor,  $K$ , is defined in terms of  $G$  as follows:

$$K = \sqrt{GE/h}. \quad (65)$$

Substitution of Eq. (63) into Eq. (65) leads to:

$$K = \sigma \sqrt{c'_g/2}, \quad (66)$$

or:

$$K = \sigma \sqrt{3\xi} f_R, \quad (67)$$

where:

$$f_R^{-1} = (1 - \xi)^{3/2} (1 + 2\xi). \quad (68)$$

Thanks to the present formulation, for the first time in the history of the fracture mechanics, computation of the stress intensity factor for the circular shaft has become possible. As will be verified in the next section, the proposed compliance and other parameters are exact and in close agreement with the test results reported in the literature.

## 7. Verifications

The proposed theory is verified in this section through four examples. The verification is done for both of the stress intensity factor and the compliance. In spite of the classical fracture mechanics in which the parameters are empirical, the proposed theory is exact. The proposed stress intensity factor is compared with the results in the literature in Examples 1 and 2.

**Example 1:** The stress intensity factor is defined by the ASTM special committee E-399 for bend specimen, BSTP (see Figure 7) [31] as:

$$K = \frac{3PL}{2Bh^{3/2}} \times \text{BSTP},$$

$$\text{BSTP} = \frac{\xi^{1/2}}{(1 + 2\xi)(1 - \xi)^{3/2}}$$

$$\left[ 1.99 - \xi \left( 1 - \xi \right) \left( 3.15 - 3.93\xi + 2.7\xi^2 \right) \right], \quad (69)$$

in which  $L$  is length,  $B$  is section width,  $h$  is the section depth, and  $P$  is the force applied to the center.

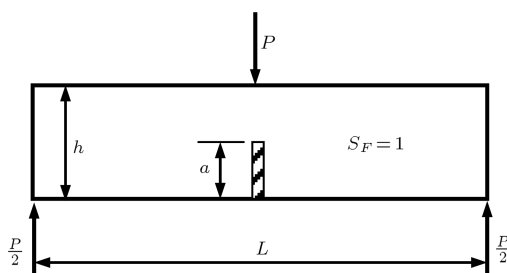


Figure 7. Bend specimen.

**Solution:** The dimensionless part of the proposed stress intensity factor,  $\text{PROPK} = \sqrt{c'_g/2}$ , is compared with BSTP dimensionless stress intensity factor of Eq. (69) in Figure 8. The results are in close agreement. The function defined in Eq. (69) has considerable similarity with the proposed function. That is the reason for close agreement of the results used as a verification for the proposed formulation. The strain energy of the bend specimen is equal to that of the core model for  $L = 2h/3$ .

**Example 2:** The stress intensity factor is defined by the ASTM special committee E-399 for Compact Tension Specimen (CTS), in Figure 9 [31], as:

$$K = \frac{P}{Bh^{1/2}} \times \text{CSTP}$$

$$\text{CSTP} = \frac{(2 + \xi)}{(1 - \xi)^{3/2}}$$

$$\left[ \left( 0.866 + 4.64\xi - 13.32\xi^2 + 14.72\xi^3 - 5.6\xi^4 \right) \right]_{(70)}$$

**Solution:** The dimensionless part of the proposed stress intensity factor,  $\text{PROPK} = \sqrt{c'_g/2}$ , is compared with CSTP, dimensionless stress intensity factor of Eq. (70), in Figure 10. Close agreement of the results

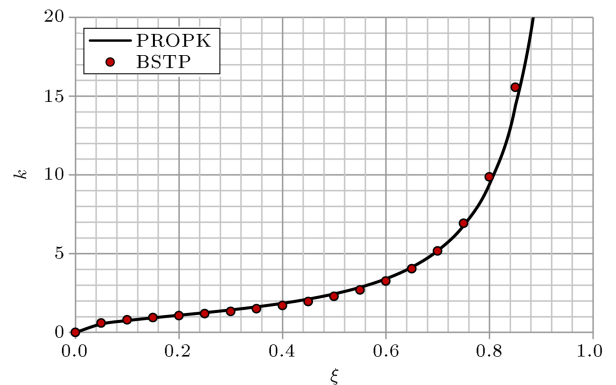


Figure 8. Comparison of PROPK with BSTP.

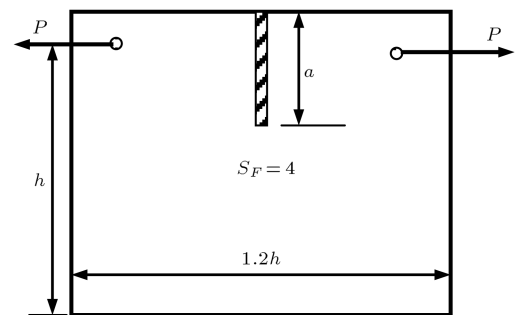


Figure 9. Compact tension specimen.

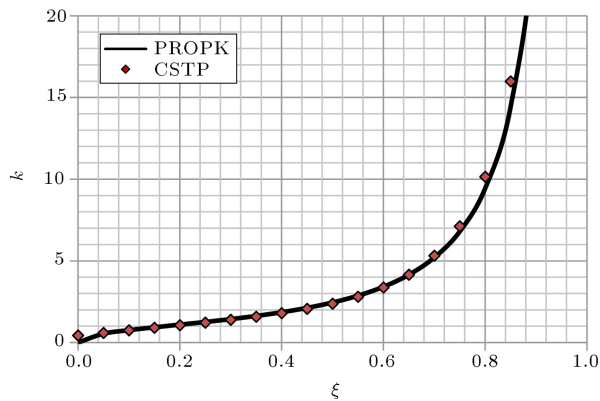


Figure 10. Comparison of PROPK with CSTP.

verified the work. The function in Eq. (70) has some similarities with the proposed formulation. Eq. (70) does not converge to zero for  $\xi = 0$ . This is the reason for some differences between this and the proposed function for small  $\xi$ s.

The proposed formulation of the crack compliance is compared with the compliances in the literatures in Examples 3 and 4.

**Example 3:** The compliance of a cracked shaft with circular section is computed by Papadopoulos [16] as:

$$\begin{aligned} \text{PAPA} = 1.2 \times & \left( +31.985\xi^6 - 18.846\xi^5 + 4.9219\xi^4 \right. \\ & \left. + 1.5285\xi^3 + 1.5785\xi^2 - 0.0301\xi \right). \end{aligned} \quad (71)$$

**Solution:** The PAPA compliance computed in Eq. (71) is compared with the proposed compliance,  $\text{PROP} = c_T$ , in Figure 11. The data for the compliance in reference [16] are given in tabular form. The function in Eq. (71) is obtained by fitting curve on these data by the present authors. The results are in close agreement at least for  $\xi < 0.5$ . The agreement was anticipated

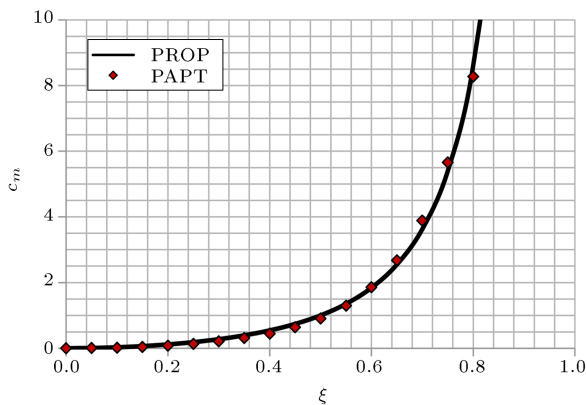


Figure 11. Comparison of PROP with PAPT.

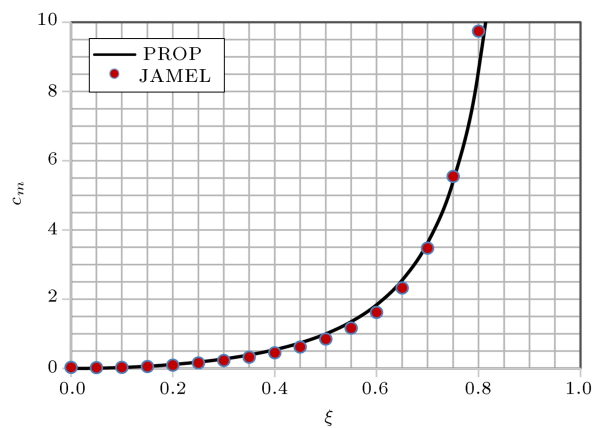


Figure 12. Comparison of PROP with JAMEL.

because the result in Eq. (71) was previously criticized and discussed in several papers [26,27].

**Example 4:** The proposed compliance is compared with that of a cracked shaft proposed by Jamel et al. [32] as:

$$\begin{aligned} \text{JAMEL} = 5\mu^{1.5} & \left( +0.035\mu^{-4} + 0.01\mu + 0.029\mu^2 \right. \\ & + 0.0086\mu^3 + 0.0044\mu^4 + 0.0025\mu^6 \\ & \left. + 0.0017\mu^7 + 0.008\mu^9 - 0.092 \right) \end{aligned} \quad (72)$$

where  $\mu = (1 - a/R)$ .

**Solution:** The proposed compliance,  $\text{PROP} = c_T$ , is compared with JAMEL of Eq. (72) in Figure 12. There were some differences in the results which are partly removed by suitable modification of the compliance function in Eq. (72) by adding the term  $5\eta^{1.5}$  to the original equation. The difference observed is due to problems in computation and extrapolation by polynomial function.

## 8. Conclusions

The following conclusions are obtained from this study:

- Based on the prior experience in analyses of cracked structures, the state functions and the generic compliance are introduced, defined as the explicit functions of the crack depth ratio, and verified;
- The proposed state functions are defined for the gradual and smooth changes between the two end states. The functions are derived based on the end conditions without any assumption in between. As a result, the compliance and other functions based on the state functions are exact;



- The crack compliance of the shaft in torsion is explicitly defined as the function of the crack depth ratio. The proposed explicit function is a good alternative for the classical procedure of the energy release rate method;
- The singularity problems, integration difficulty, and other systematic problems available in the classical energy release rate theory for the circular sections are not present in the proposed theory;
- For the first time in the history of the classical fracture mechanics, the computation of the stress intensity factor for circular shaft has become possible;
- The accuracy and efficiency of the work are verified by concise mathematical formulation and comparison of the results of the work with the others.

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