



Bioconvection phenomenon for the boundary layer flow of magnetohydrodynamic Carreau liquid over a heated disk

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Received 8 July 2019; received in revised form 25 October 2020; accepted 26 April 2021

KEYWORDS

Brownian diffusion;
 Heat and mass transport;
 Carreau fluid;
 Bioconvective process;
 Rotating heated disk;
 Boundary layer analysis.

Abstract. This study carried out a numerical examination on the effect of magnetohydrodynamic steady flow of Carreau fluid on the transfer of thermal energy and mass species comprising nanoparticles with gyrotactic microorganisms through a heated disk. The roles of thermophoresis and Brownian motion were also considered in resolving this flow problem. Governing equations were solved using boundary layer theory emphasizing the coupled system of Partial Differential Equations (PDEs) involving boundary conditions. The highly non-linear system of Ordinary Differential Equations (ODEs) was generated using transformation approach. Due to the highly nonlinear form of the system of transformed equations, an approximate solution was presented which was evaluated using optimal homotopy method. Moreover, the effects of prominent parameters on velocity, thermal energy, mass species, and motile density microorganisms were graphically examined. In addition, graphical observations regarding the mass species, thermal energy, and velocities were briefly discussed. It was estimated that in an intensified magnetic field, the motion of fluid particles and temperature of fluid would decrease and increase, respectively.

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1. Introduction

Due to their extensive applicability, non-Newtonian fluids have drawn much academic attention. This research is motivated by the abundance of these fluids in nature, as well. Such fluids are widely used in

different fields such as foodstuffs, extrusion of molten polymers and plastics, fiber synthesis, drilling gas and oil wells, etc. Many studies have been conducted on non-Newtonian fluids and their phenomenal roles. For instance, a number of researchers have investigated how to improve the thermal conductivity of non-Newtonian liquids. In addition, some others have shifted their focus to the difference in modeling the power law for non-Newtonian fluids. However, the power law model suffers some drawbacks including low and high shear rates and then, authors have examined the viscosity

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model called Carreau rheology. This is a special type of Newtonian liquid and the shear rate is a function of viscosity; this phenomenon is useful for high shear-rate liquids. Carreau [1,2] employed an important theory regarding rheological equations based on network molecular models. Griffiths [3] studied the flow behavior of generalized Newtonian liquid through disk by applying the Carreau fluid model. Machireddy and Naramgari [4] discussed the role of transfer of thermal energy and mass species with cross-diffusion involving Magnetohydrodynamic (MHD) Carreau fluid on the stretch surface. Another notable study on Carreau fluid can be found in [5–8] with various explorations therein. Nanoparticles are tiny particles made of dense nanoparticles or nanofibers ranging between (1–100 nm) which are normally equated through the conventional heat carrying fluids, and they enjoy an advanced thermal conductivity. As can be found in the literature, Choi and Eastman [9] initially put forward the idea of utilizing fluid consisting of nano-sized particles and base fluid, called nanofluid. There are many applications containing insulation of energy, astronomical, cooling processes, solar amusement and defense, magnetic sticking, mass/heat transport strengthening and medical instruments, etc. Such applications require substantial point altering from conventional fluids. Therefore, the current authors of this study have become drawn to nanofluid and discussed some useful findings which were previously extracted from Refs. [10–15].

MHD is an added working zone of building sciences nowadays that includes the effect of magnetic fields. Applications of such types of fluid flow are pumps, power generators, magnetic drug treatment, accelerators, plasma studies, and flow meters. Bhatti et al. [16] discussed mathematical modeling of mass and heat effects on the flow of electrically conducting fluid for two-phase peristaltic propulsion through a porous medium with Darcy-Brinkman-Forchheimer. Many researchers [17–20] interpreted different aspects of non-Newtonian fluid such as Williamson, Micropolar and Carreau fluids, etc.

Heat forms out of energy and is transmitted from one place to another, the difference of which may turn into thermal energy. Physics can help read the investigations to produce the energy difference called heat transfer. The mathematical form of heat transfer is derived from Fourier's law of conduction. For heat transfer measurement, thermal conductivity is a very important factor which is defined as the ability to measure heat conduction. The enhancement of thermal conductivity for a material is suggestive of being a good conductor while a poor insulator results from low thermal conductivity. Similarly, transport of mass species involves the movement or diffusion of fluid particles from one place to another. Transport

of thermal energy and mass species are the kinetic processes that may be investigated, either separately or jointly. The transport of thermal energy and solubility of nanofluid can be investigated based on Fick's and Fourier's laws. These movements are modeled by similar mathematical equations in the form of convection and diffusion and both transfer types must be considered jointly in some cases, i.e., ablation and evaporative cooling. Applications of mass species and thermal energy transport in different fields including oil transport wonder, dispersion of specific medications in blood, food preparation, cooling of electronic equipment, manufacturing/materials processing, absorption, drying, precipitation, membrane filtration, and evaporation can be seen in [21]. For further information on mass and heat transfer applications, readers are referred to [22–24] and the studies cited therein for further insights.

Rotation is the most powerful and useful tool for such as applications medical equipment, gas turbine, food processing, and computer operating, while numerous other applications have been found in food of rotating geometries (disk, cylinder, and surface). It is evident that rotating disk has an important role from the research viewpoint. Concept of rotating disk was developed by Karman first [25]. He employed transformation (Von-Karman) to evaluate solution flow problems over heated rotating disk. Several other applications for rotations can be found in [26–28].

The particles of this impact (bio-convection) which are not self-boosted microorganisms have been investigated. Another terminology for the bio-convection approach is called boosted microorganism. So, this type of terminology was initiated by Platt [29] who concluded that drag force could be generated from the movement of microorganisms while gravitation torque was produced due to the equilibrium position of particles in the cells of swimming microorganism. Chakraborty et al. [30] explored the remaining unknown aspects of the magnetic field and nanoparticles with emphasis on gyrotactic microorganisms. Impact between gyrotactic microorganisms and nanoparticles and the resulting radiation were measured by Khan et al. [31]. For deeper insights, readers are referred to the works mentioned in [32–34] and other references therein.

Satisfactory performance of a solution should be taken into consideration from two aspects: the case of approximate solution and generation of an accurate approximate solution with different parameters. Numerous approaches are employed to find a solution to linear flow problems. In the case of analytical technique, an analysis approach (Optimal Homotopy Analysis Method (OHAM)) captures a solution to non-linear flow problems involving Boundary Conditions (BCs). Recently, the OHAM approach has been

adopted by Marinca et al. [35], with few relevant studies on algorithm being accessible in [36–39]. Makinde and Animasaun [40] proposed a new bouncy induced procedure for nanoparticles and considered volume fraction by causing variations in thermal conductivity for this flow problem while solving the flow problem by the (RK4SM) approach. Makinde and Animasaun [41] investigated the phenomenon of flow in terms of Brownian motion, bouncy force, and bio-convection through parabolic surface with microorganisms. Mutuku and Makinde [42] discussed the characteristics of bio-convection subject to hydro-magnetics considering nanoparticles. Khan et al. [43] explored lesser known aspects of gyrotactic microorganisms with nanoparticles and their impact on the transport of mass species and heat energy in the magnetic field. Makinde et al. [44] developed a flow model upon taking into account the influence of radiations, chemical reaction, Brownian motion, and magnetic force over vertical plate. The latest important contributions dealing with the flow problems were made and reported in [45–47].

This paper aims to delve into the phenomenon of flow emphasizing electrical conducting of Carreau rheological fluid with nanoparticles and gyrostatic microorganisms through heated cones. The system of Ordinary Differential Equations (ODEs) was derived from the system of Partial Differential Equations (PDEs) using Von-Karman transformations and the approach of OHAM analysis. This study is structured as follows: after presented a literature review, the mathematical formulation is developed in Section 2. The formulation of flow problem and numerical solution are captured in Sections 3 and 4, respectively. The key points regarding flow problem are added in Section 5. In the end, references are listed.

2. Mathematical formulation and fluid rheology

In this analysis, we have considered the flow of electrically conducting fluids including two-dimensional time-independent incompressible Carreau fluid and nanoparticles with motile gyrotactic microorganisms, as induced by a rotating disk. The magnetic field strength (B_0) on a boundary layer acts along the z -direction, while the motion of fluid particles is generated by the movement of wall velocity $u_s (= r l_0)$ where (l_0) is a constant. Geometrical flow under the current assumption is captured and given in Figure 1. The effect of induced magnetic field is not taken into account, but the features of thermophoresis and ambient motion are observed. The angular velocity (Ω_1) represents the rotational velocity of a rotating disk involving viscous dissipation. The velocity components are based on directions of (r, θ, z) . Initially, a disk is heated at (T_0) temperature, after which it adapts to

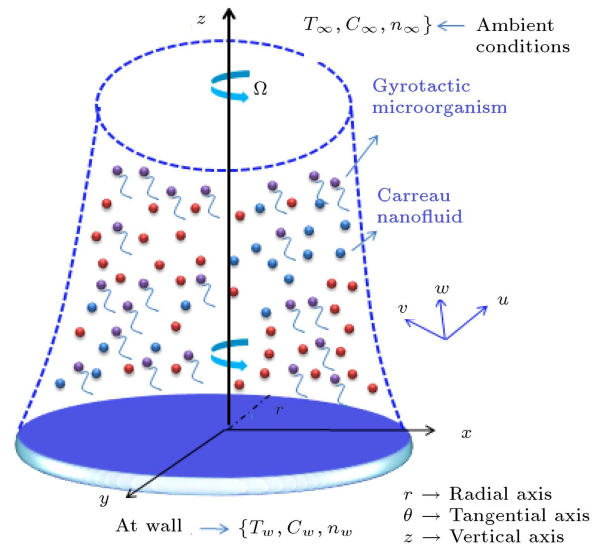


Figure 1. Fluid flow geometry.

and takes ambient temperature (T_∞). Bio-convective patterns are detected based on the movement of motile microorganisms from higher areas to lower regions. The concentrations of reference and ambient microorganisms are taken by n_0 and n_∞ , respectively. The stress tensor [1] is expressed as follows:

$$\tau_1 = \left[\eta_0 \left(1 + \lambda^2 \dot{\gamma}^2 \right)^{\frac{n-1}{2}} \right] \dot{\gamma}, \quad (1)$$

$$\dot{\gamma} = \left(\frac{1}{2} \sum_i \sum_j (\gamma_{ji} \gamma_{ij}) \right)^{1/2}, \quad (2)$$

where n is the power law index. It was estimated that Carreau rheology would become ($0 < n < 1$) shear thinning and ($n > 1$) shear thickening. Governing laws under motile microorganism and nanomaterial are handled in the following equations:

$$\nabla \cdot V = 0,$$

$$\rho_f [V \cdot \nabla] V = -\nabla P + \nabla \cdot \tau_1 + J_1 \times B,$$

$$J_1 = \sigma (V \times B),$$

$$B = [0, 0, B_0],$$

$$V \cdot \nabla T - \alpha^* \nabla^2 T = \tau \left[D_b \nabla T \cdot \nabla C + \left(\frac{D_t}{T_\infty} \right) \nabla T \cdot \nabla T \right],$$

$$(V \cdot \nabla) C = D_b (\nabla^2 C) + D_b (\nabla^2 C) + \left(\frac{D_t}{T_\infty} \right) \nabla^2 T,$$

$$\nabla \cdot J^* = 0,$$

where V is velocity. u_1, v_1, w_1 are flow components; P is pressure; ρ_f is fluid density; T is temperature; α^* is thermal diffusivity, C is concentration, D_t, D_b are

thermophoretic diffusion and Brownian numbers; and J^* is microorganisms flux.

$$J^* = nV + n \cdot \hat{V} - D_m \nabla n,$$

$$\hat{V} = \left(\frac{bW_c}{\Delta C} \right) \nabla C.$$

The above equations after the boundary layer approximations are expressed as follows:

$$\frac{\partial u_1}{\partial r} + \frac{u_1}{r} + \frac{\partial w_1}{\partial z} = 0, \quad (3)$$

$$\begin{aligned} \rho_f \left(u_1 \frac{\partial u_1}{\partial r} - \frac{v_1^2}{r} + w_1 \frac{\partial u_1}{\partial z} \right) = -\frac{\partial p}{\partial r} + \frac{\partial^2 u_1}{\partial z^2} \eta_0 \\ \left[1 + \frac{(n-1)}{2} \lambda^2 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^2 + \left(\frac{\partial u_1}{\partial z} \right)^2 \right\} + \frac{(n-1) \lambda^4 (n-3)}{4(2)} \right. \\ \left. \left\{ \left(\frac{\partial v_1}{\partial z} \right)^4 + \left(\frac{\partial u_1}{\partial z} \right)^4 + 2 \left(\frac{\partial v_1}{\partial z} \frac{\partial u_1}{\partial z} \right)^2 \right\} \right] \\ + \eta \frac{\partial u_1}{\partial z} \frac{\partial}{\partial z} \\ \left[\frac{(n-1)}{2} \lambda^2 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^2 + \left(\frac{\partial u_1}{\partial z} \right)^2 \right\} + \frac{(n-1) \lambda^4 (n-3)}{4(2)} \right. \\ \left. \left\{ \left(\frac{\partial v_1}{\partial z} \right)^4 + \left(\frac{\partial u_1}{\partial z} \right)^4 + 2 \left(\frac{\partial v_1}{\partial z} \frac{\partial u_1}{\partial z} \right)^2 \right\} \right] \\ - \sigma B_0^2 u_1, \end{aligned} \quad (4)$$

$$\begin{aligned} \rho_f \left(u_1 \frac{\partial v_1}{\partial r} - \frac{u_1 v_1}{r} + w_1 \frac{\partial v_1}{\partial z} \right) = \frac{\partial^2 v_1}{\partial z^2} \eta_0 \\ \left[1 + \frac{(n-1)}{2} \lambda^2 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^2 + \left(\frac{\partial u_1}{\partial z} \right)^2 \right\} + \frac{(n-1) (n-3)}{8} \lambda^4 \right. \\ \left. \left\{ \left(\frac{\partial v_1}{\partial z} \right)^4 + \left(\frac{\partial u_1}{\partial z} \right)^4 + 2 \left(\frac{\partial v_1}{\partial z} \frac{\partial u_1}{\partial z} \right)^2 \right\} \right] \\ + \eta_0 \frac{\partial v_1}{\partial z} \frac{\partial}{\partial z} \\ \left[\frac{(n-1)}{2} \lambda^2 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^2 + \left(\frac{\partial u_1}{\partial z} \right)^2 \right\} + \frac{(n-1) (n-3)}{8} \lambda^4 \right. \\ \left. \left\{ \left(\frac{\partial v_1}{\partial z} \right)^4 + \left(\frac{\partial u_1}{\partial z} \right)^4 + 2 \left(\frac{\partial v_1}{\partial z} \frac{\partial u_1}{\partial z} \right)^2 \right\} \right] \\ - \sigma B_0^2 v_1, \end{aligned} \quad (5)$$

$$\begin{aligned} \left(u_1 \frac{\partial T}{\partial r} + w_1 \frac{\partial T}{\partial z} \right) = \alpha^* \frac{\partial^2 T}{\partial z^2} \\ + \tau \left[D_b \left(\frac{\partial T}{\partial z} \frac{\partial C}{\partial z} \right) + \frac{D_t}{T_\infty} \left(\frac{\partial T}{\partial z} \right)^2 \right], \end{aligned} \quad (6)$$

$$\left(u_1 \frac{\partial C}{\partial r} + w_1 \frac{\partial C}{\partial z} \right) = D_b \left(\frac{\partial^2 C}{\partial z^2} \right) + \frac{D_t}{T_\infty} \left(\frac{\partial^2 T}{\partial z^2} \right), \quad (7)$$

$$\begin{aligned} \left(u_1 \frac{\partial n}{\partial r} + w_1 \frac{\partial n}{\partial z} \right) = D_m \left(\frac{\partial^2 n}{\partial z^2} \right) \\ - \frac{bW_c}{\Delta C} \left(\frac{\partial}{\partial z} \left(n \frac{\partial C}{\partial z} \right) \right). \end{aligned} \quad (8)$$

2.1. Boundary conditions

The BCs subjected to fluid flow are captured as follows:

$$u_1 = u_s = l_0 r, \quad v_1 = r \Omega_1, \quad w_1 = 0,$$

$$T = T_0, \quad n = n_0, \quad C = C_0, \quad \text{at } z = 0,$$

$$T \rightarrow T_\infty, \quad v_1 = u_1 \rightarrow 0, \quad n \rightarrow n_\infty,$$

$$C \rightarrow C_\infty, \quad \text{at } z \rightarrow \infty. \quad (9)$$

2.2. Similarity analysis

The transformations called selection of Von-Karman are expressed as follows:

$$u_1 = (\Omega_1 r) f(\eta), \quad v_1 = (\Omega_1 r) g(\eta),$$

$$w_1 = \left(\sqrt{\nu_f \Omega_1} \right) H(\eta), \quad \eta = \left(\sqrt{\frac{\Omega_1}{\nu_f}} \right) z,$$

$$\theta(\eta) = \frac{T - T_0}{T_0 - T_\infty}, \quad \phi(\eta) = \frac{C - C_0}{C_0 - C_\infty},$$

$$\xi(\eta) = \frac{n - n_\infty}{n_0 - n_\infty}. \quad (10)$$

The set of dimensionless ODEs is generated as follows:

$$-(f + f) + H' = 0, \quad (11)$$

$$\begin{aligned} H''' - HH'' - \frac{1}{2} H'(H') - \frac{g^2}{2} + \frac{n-1}{2} \lambda_1 \text{Re} \\ \left[-H''(g')g'' - H''' g'^2 - H''' \frac{3}{4} H'' H'' \right] \\ + \frac{(n-3)(n-1)}{4(2)} (\lambda_1 \text{Re}) \lambda_1 \text{Re} \\ \left[-(g')^4 H''' - \frac{5}{16} (H'')^4 H''' - (H'')^2 \frac{3}{2} g'^2 H''' \right. \\ \left. - g'' 2H''(g')^3 - g'' (H'')^3 g' \right] + M H' = 0, \end{aligned} \quad (12)$$

$$\begin{aligned} g'' - g'H - H'g + \lambda_1 \frac{n-1}{2} \text{Re} \\ \left[\frac{g'' H''^2}{4} + 3g'^2 g'' + g' \frac{H''}{2} H''' \right] \\ + \frac{(n-3)(n-1)}{4(2)} (\lambda_1 \text{Re})^2 \left[5g'' g'^4 + \frac{1}{16} g'' H''^4 \right. \\ \left. + \frac{3}{2} H''^2 g'^2 g'' + \frac{g' H^{(iv)}}{4} H'''^3 + \frac{(g')^3 H'' H''' }{4} \right] - M g = 0, \end{aligned} \quad (13)$$

$$\theta'' + N_b \theta' \phi' + N_t \theta'^2 - \text{Pr} H \theta' = 0, \quad (14)$$

$$\phi'' + \frac{N_t}{N_b} \theta'' - \text{Sc} \phi' = 0, \quad (15)$$

$$\xi'' - Sc\xi'H + Pe[\phi''(\xi - \Omega) + \xi'\phi'] = 0. \quad (16)$$

BCs regarding the flow problem are developed as follows:

$$\begin{aligned} \text{At } \eta = 0, \quad H = 0, \quad -2St = H', \quad g = 1, \\ \theta = 1, \quad \phi = 1, \quad \xi = 1. \\ \text{At } \eta \rightarrow \infty, \quad H' = 0, \quad g = 0, \quad \theta = 0, \\ \phi = 0, \quad \xi = 0. \end{aligned} \quad (17)$$

The influential dimensionless parameters are listed as follows:

$$\begin{aligned} \lambda_1 = r^2\Omega_1^2, \quad \text{Re} = \frac{r^2\Omega_1}{\nu_f}, \quad M = \frac{\sigma B_0^2}{\rho_f\Omega_1}, \\ \text{Pr} = \frac{\nu_f}{\alpha^*}, \quad St = \frac{l_0}{\Omega_1}, \quad N_b = \frac{D_b\tau}{\alpha^*}(C_0 - C_\infty), \\ N_t = \frac{\tau}{\alpha^*} \frac{D_t}{T_\infty}(T_0 - T_\infty), \quad \text{Pe} = \frac{bW_c}{D_m}, \\ \text{Sc} = \frac{\nu_f}{D_m} = \frac{\nu_f}{D_b}, \quad \Omega = \frac{n_\infty}{(n_\infty - n_0)}. \end{aligned} \quad (18)$$

where λ_1 is Carreau fluid, M Hartmann, N_t thermophoresis motion, Sc Schmidt number, St stretching rate, Pr Prandtl, Re local Reynolds, Pe bio-convection Peclet, and N_b Brownian motion numbers.

2.3. Gradient velocity and flux numbers

The gradient velocity (C_f), Nusselt (Nu_r), Sherwood and motile microorganisms (Nn_r) numbers are formulated as follows:

$$C_f = \frac{\sqrt{\tau_{rz} + \tau_{r\theta}}}{\rho_f(r\Omega_1)^2}, \quad Nu_r = \frac{rq_1}{k(T_0 - T_\infty)}, \quad (19)$$

$$Sh_r = \frac{rq_2}{D_b(C_\infty - C_0)}, \quad Nn_r = \frac{rq_3}{Dm(n_\infty - n_0)}, \quad (20)$$

$$\tau_{rz}|_{z=0} = \frac{\partial u_1}{\partial z}\eta_0$$

$$\left[\begin{aligned} &1 + \frac{(n-1)}{2}\lambda^2 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^2 + \left(\frac{\partial u_1}{\partial z} \right)^2 \right\} + \\ &\frac{(n-1)(n-3)}{8}\lambda^4 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^4 + \left(\frac{\partial u_1}{\partial z} \right)^4 + 2\left(\frac{\partial v_1}{\partial z} \frac{\partial u_1}{\partial z} \right)^2 \right\} \end{aligned} \right], \quad (21)$$

$$\tau_{\theta z}|_{z=0} = \frac{\partial v_1}{\partial z}\eta_0$$

$$\left[\begin{aligned} &1 + \frac{(n-1)}{2}\lambda^2 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^2 + \left(\frac{\partial u_1}{\partial z} \right)^2 \right\} + \\ &\frac{(n-1)(n-3)}{8}\lambda^4 \left\{ \left(\frac{\partial v_1}{\partial z} \right)^4 + \left(\frac{\partial u_1}{\partial z} \right)^4 + 2\left(\frac{\partial v_1}{\partial z} \frac{\partial u_1}{\partial z} \right)^2 \right\} \end{aligned} \right], \quad (22)$$

$$q_1 = \left| -k \frac{\partial T}{\partial z} \right|_{z=0},$$

$$q_2 = \left| -D_b \frac{\partial C}{\partial z} \right|_{z=0},$$

$$q_3 = \left| -D_m \frac{\partial n}{\partial z} \right|_{z=0}. \quad (23)$$

Through relations, gradient velocity (C_f), Nusselt (Nu_r), Sherwood and motile microorganisms (Nn_r) numbers resulting from the dimensionless form are as follows:

$$\begin{aligned} \text{Re}^{\frac{1}{2}}(C_f) &= (g'^2 + f'^2)^{1/2} [1 + \lambda_1 \text{Re}(g'g' + f'^2)]^{\frac{n-1}{2}}, \\ Nu_{1r} \text{Re}^{-\frac{1}{2}} &= -\theta'(0), \\ Sh_r \text{Re}^{-\frac{1}{2}} &= -\phi'(0), \quad Nn_r \text{Re}^{-\frac{1}{2}} = -\xi'(0). \end{aligned} \quad (24)$$

3. Solution analysis and physical description

The present section captures the role of fluid flow, heat energy, and mass species curves versus different parameters such as Carreau (λ_1), Hartmann (M), Prandtl (Pr), thermophoresis motion (N_t), Brownian motion (N_b), bioconvective Peclet (Pe), and power-law index (n) numbers. Figures 2–19 are sketched by applying OHAM using Mathematica 10.0. Bar charts (Figures 18–21) and numerical values in tabular forms (Tables 1–4) of different parameters are analyzed for gradient velocity (C_f), rate of heat energy transfer (Nu_r), Sherwood number (Sh_r), and motile microorganisms (Nn_r).

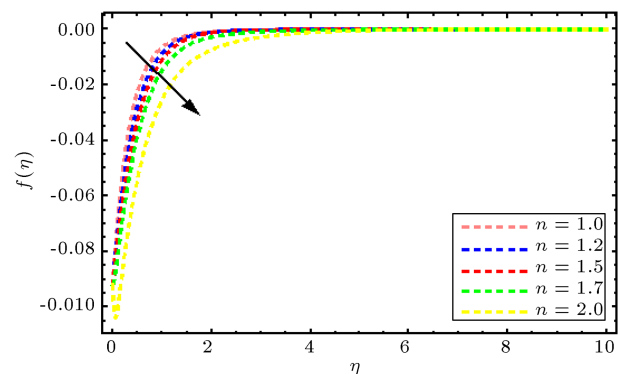


Figure 2. Character of n regarding $f(\eta)$.

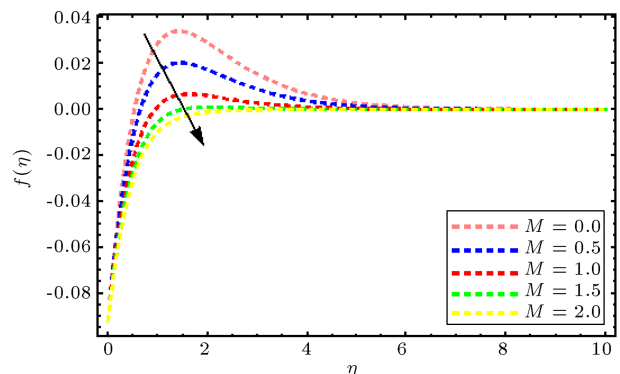
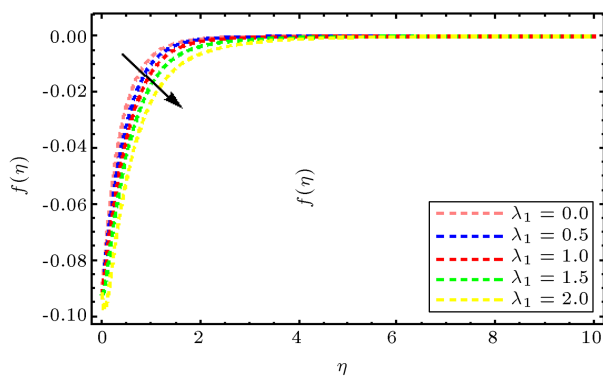
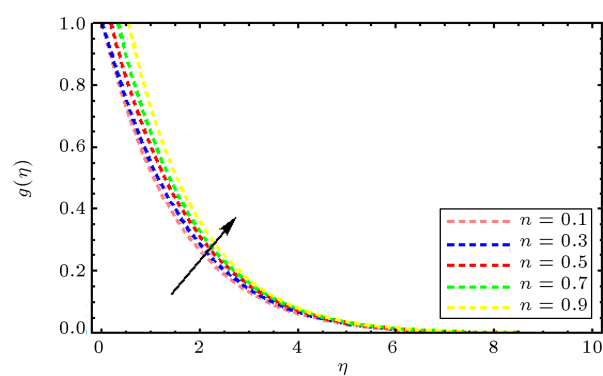
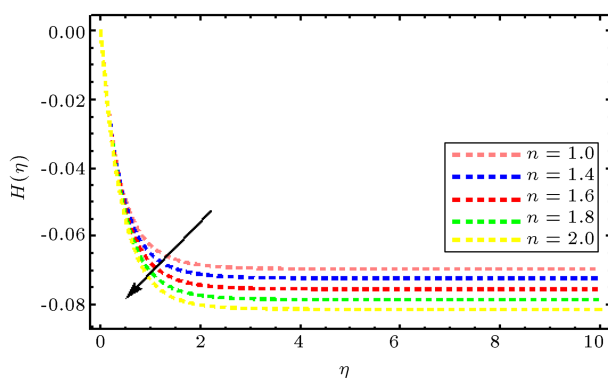
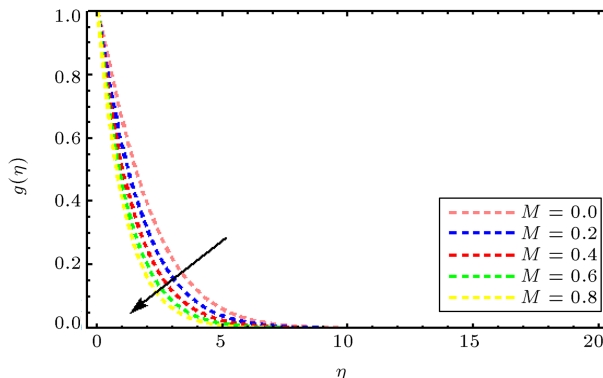
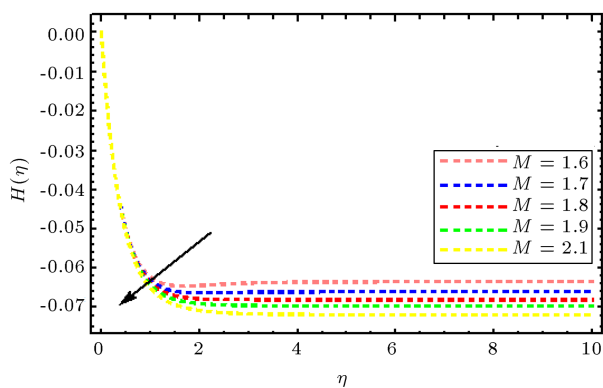
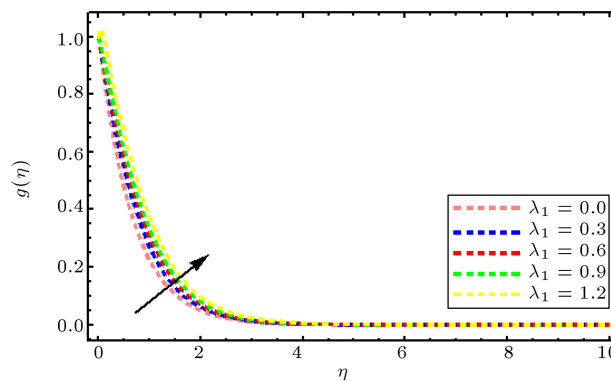
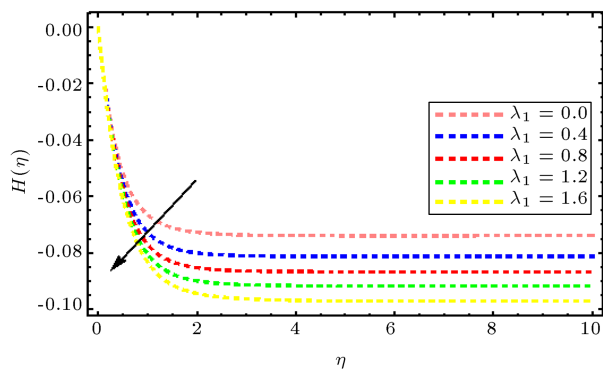
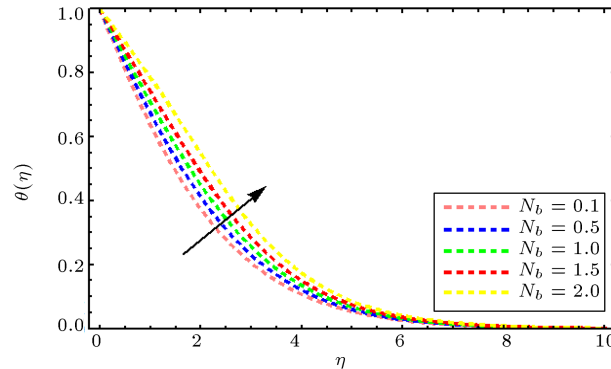


Figure 3. Character of M regarding $f(\eta)$.

Figure 4. Character of λ_1 regarding $f(\eta)$.Figure 8. Character of n regarding $g(\eta)$.Figure 5. Character of n regarding $H(\eta)$.Figure 9. Character of M regarding $g(\eta)$.Figure 6. Character of M regarding $H(\eta)$.Figure 10. Character of λ_1 regarding $g(\eta)$.Figure 7. Character of λ_1 regarding $H(\eta)$.Figure 11. Character of N_b regarding $\theta(\eta)$.

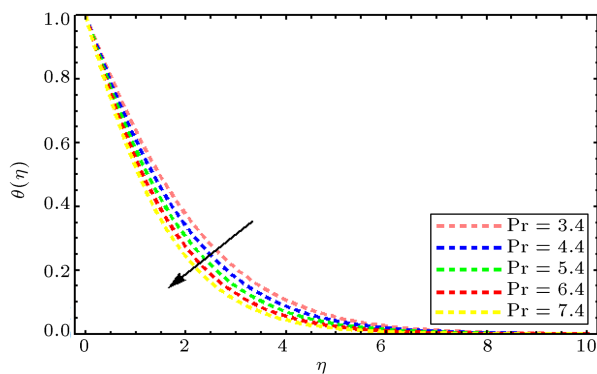


Figure 12. Character of Pr regarding $\theta(\eta)$.

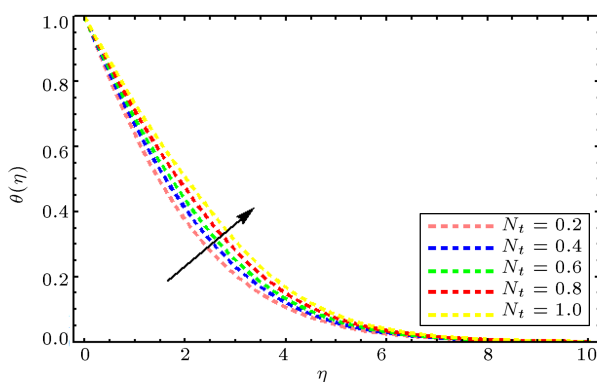


Figure 13. Character of N_t regarding $\theta(\eta)$.

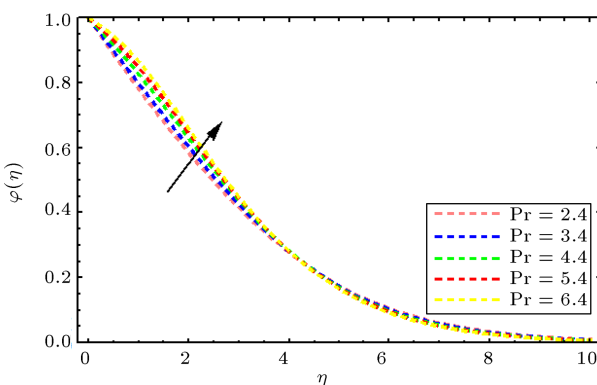


Figure 14. Character of Pr regarding $\varphi(\eta)$.

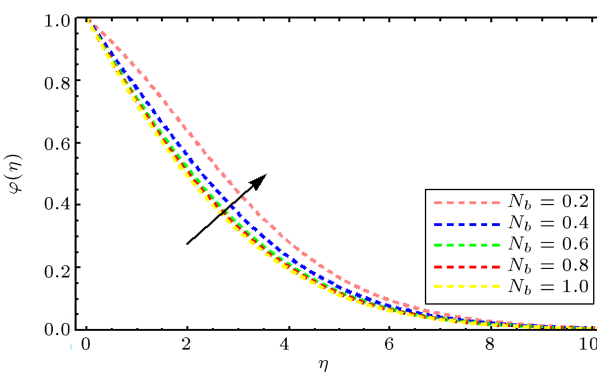


Figure 15. Character of N_b regarding $\varphi(\eta)$.

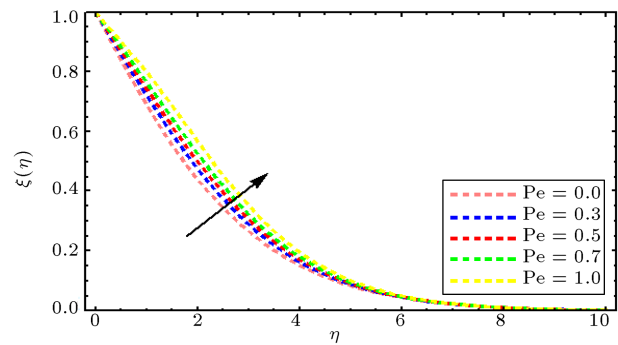


Figure 16. Character of Pe regarding $\xi(\eta)$.

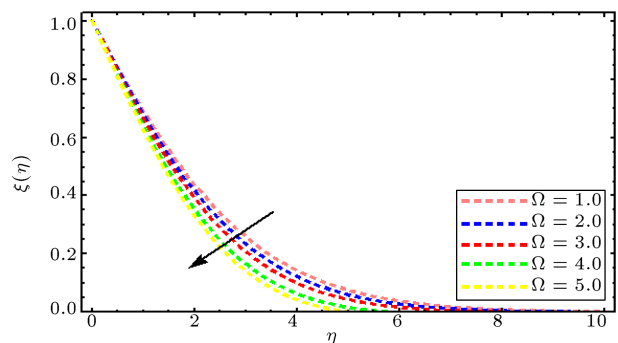


Figure 17. Character of Ω regarding $\xi(\eta)$.

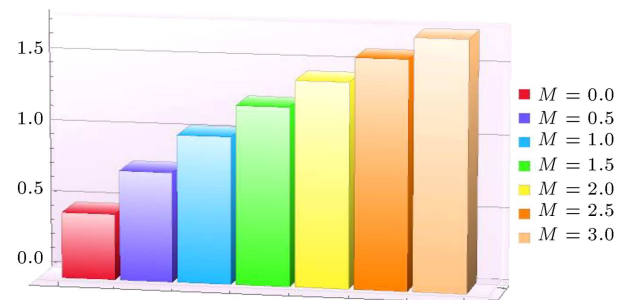


Figure 18. Bar chart of $C_f(\eta)$.

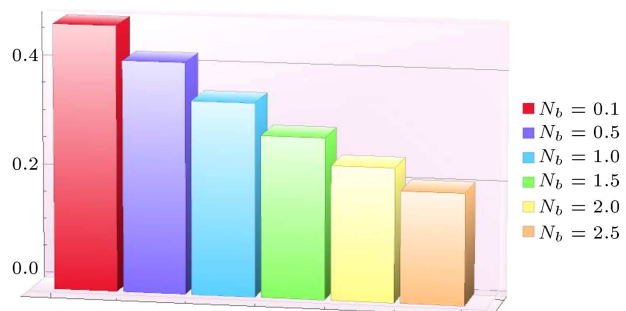
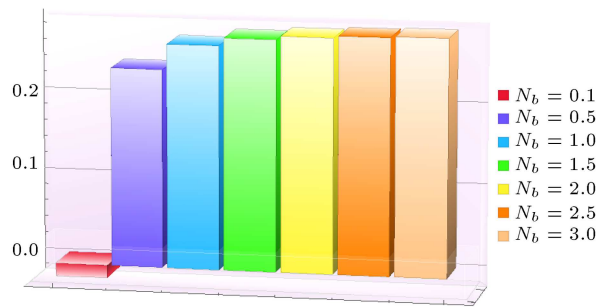
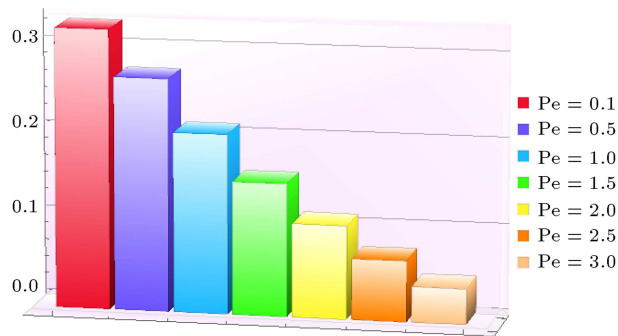


Figure 19. Bar chart of $Nu_r(\eta)$.

The impact of n on fluid flow is discussed and given in Figure 2. According to $n > 1$, the fluid acts as a shear thickening behavior since $f(\eta)$ decreases. In Figure 3, at large values of (M) flow declines due to the enhancement of Lorentz force and generation of greater resistance in fluid flow particles. Similarly, Figure 4 shows the decreasing behavior for $f(\eta)$ following the

**Figure 20.** Bar chart of $Nn_r(\eta)$.**Figure 21.** Bar chart of $Nn_r(\eta)$.**Table 1.** Numerical values of skin friction coefficient

$\text{Re}^{\frac{1}{2}} C_f(0)$ when $St = 0.09$, $\lambda_1 = 0.9$, $\text{Re} = 1.2$, $\text{Pr} = 6.7$, $N_t = 0.1$, $\text{Sc} = 1.2$, $\Omega = 0.2$, $\text{Pe} = 0.7$, $N_b = 0.3$.

M	n	λ_1	Re	$\text{Re}^{\frac{1}{2}} C_f$
$\frac{1}{10}$	01	$\frac{5}{10}$	$\frac{9}{10}$	0.423187
$\frac{2}{10}$	—	—	—	0.500069
$\frac{3}{10}$	—	—	—	0.572402
$\frac{4}{10}$	—	—	—	0.640583
$\frac{1}{10}$	02	$\frac{5}{10}$	$\frac{9}{10}$	0.429413
$\frac{2}{10}$	—	—	—	0.504486
$\frac{3}{10}$	—	—	—	0.576577
$\frac{4}{10}$	—	—	—	0.645889
1	01	$\frac{5}{10}$	$\frac{9}{10}$	0.981269
$\frac{15}{10}$	—	—	—	1.20359
2	—	—	—	1.39237
$\frac{25}{10}$	—	—	—	1.55897
1	02	$\frac{5}{10}$	$\frac{9}{10}$	1.01159
$\frac{15}{10}$	—	—	—	1.26546
2	—	—	—	1.49019
$\frac{25}{10}$	—	—	—	1.69825

escalation of Carreau parameter (λ_1). Figure 5 indicates that $H(\eta)$ velocity slows down versus large values of n due to the shear thickening fluid. Similar types of behavior are given in Figure 6 against the values of M , while fluid particles decay because of resistive force and the motion fluid particles slow down. Figure 7 reveals

Table 2. Numerical study of Nusselt number $\text{Re}^{-\frac{1}{2}} Nu_r$ when $n = 2.9$, $\lambda_1 = 0.9$, $\text{Re} = 1.2$, $\Omega = 0.2$, $M = 2.02$, $\text{Pe} = 0.5$, $St = 0.3$.

N_b	N_t	Sc	Pr	$\text{Re}^{-\frac{1}{2}} Nu_r(0)$
0.1	0.1	1.0	6.4	0.480056
0.2	—	—	—	0.463862
0.3	—	—	—	0.448108
0.4	—	—	—	0.432789
0.1	0.2	1.0	6.4	0.467692
0.2	—	—	—	0.451871
0.3	—	—	—	0.436486
0.4	—	—	—	0.421534

Table 3. Numerical results with respect to Sherwood number $\text{Re}^{-\frac{1}{2}} Sh_r$ when $St = 0.09$, $M = 2.02$, $\lambda_1 = 0.9$, $\text{Re} = 1.2$, $\Omega = 0.2$, $\text{Pe} = 0.7$.

N_b	N_t	Sc	Pr	$\text{Re}^{-\frac{1}{2}} Sh_r(0)$
0.1	0.1	1.0	6.4	0.015659
0.2	—	—	—	-0.145346
0.3	—	—	—	-0.198932
0.4	—	—	—	-0.225441
0.1	0.2	1.0	6.4	-0.298010
0.2	—	—	—	-0.011591
0.3	—	—	—	-0.114436
0.4	—	—	—	-0.165596

Table 4. Numerical study for $\text{Re}^{-\frac{1}{2}} Nn_r(0)$ when $St = 0.09$, $M = 2.02$, $\lambda_1 = 0.9$, $\text{Re} = 1.2$, $N_t = 0.1$, $\text{Pr} = 6.7$, $N_b = 0.3$.

Pe	Sc	Ω	$\text{Re}^{-\frac{1}{2}} Nn_r(0)$
0.1	0.5	0.2	0.32443
0.2	—	—	0.310349
0.3	—	—	0.296544
0.4	—	—	0.283021
0.1	1.0	—	0.340332
0.2	—	—	0.325989
0.3	—	—	0.311915
0.4	—	—	0.298117

the characteristics and effects of λ_1 on flow $H(\eta)$. It can be measured that increase in λ_1 points to the attenuation of flow phenomenon $H(\eta)$. The impact of n on flow $H(\eta)$ is estimated in Figure 8. In this figure, the fluid flow is shown to be increasing for n resulting from $n < 1$. The velocity of fluid $g(\eta)$ decreases due to the increasing values of intensity of M . According

to Figure 9, it is examined that fluid flow accelerates against the large values of Carreau liquid, as can be implied from Figure 10.

Increase in the value of N_b and the subsequent heat energy profile $\theta(\eta)$ outcome are given in Figure 11. It was found that heat energy curves $\theta(\eta)$ increased because of the increment of N_b in fluid particle motion. Temperature profile $\theta(\eta)$ boosts up gradually due to collision between particles. Thermal Boundary Layer (TBL) is reduced upon increasing the values of Pr. Consequently, fluid temperature $\theta(\eta)$ diminishes, as can be shown in Figure 12. The temperature is enhanced because high thermophoresis parameter is N_t . In thermophoresis, the heat from the fluid is reduced and thermal energy increases. The effects of N_t on thermal energy curves $\theta(\eta)$ are given in Figure 13. Concentration profile of nanoparticles $\varphi(\eta)$ increases due to the enhancement of momentum diffusivity. Behavior of these parameters with respect to $\varphi(\eta)$ is shown in Figure 14. Brownian motion parameter (N_b) exhibits the increasing tendency of mass species profile $\varphi(\eta)$ (see Figure 15) and the concentration decreases due to the enhancement of random motion of particles and kinetic energy versus higher N_b . Upon increase in Pe, the speed of cells of swimming microorganisms increased. Through the pattern concentration of motile microorganisms, $\xi(\eta)$ in a moving disk increases. The graphical impact of Pe on the concentration of microorganisms $\xi(\eta)$ is depicted in Figure 16. The effect of difference in microorganism concentration on the concentration of motile microorganisms is depicted in Figure 17. To enhance the concentration difference (Ω), the concentration of motile microorganisms increases for ambient fluid and, yet, declines on the surface of $\xi(\eta)$. Figure 18 shows the bar chart for gradient velocity with different values of Hartmann number (M). The figure points to the increase in values of M due to high resistance between the fluid particles near the surface. Bar charts for both Nu_r and Sh_r show the opposite effect of different active values of N_b according to Figures 19 and 20, respectively. To enhance the values of N_b and facilitate the transport of conductive thermal energy, heat is generated due to the diffusion of nanoparticles which reduces Nu_r while increasing Sh_r . Similarly, the bar chart for Nn_r is plotted related to active parameter Pe in Figure 21. The numerical results regarding gradient velocity are evaluated in Table 1. According to this table, the gradient velocity reduces due to large magnetic values and enhancement of resistance to friction. Further, the rate of heat transport is reduced versus values of (N_b) and (N_t), whereas the rate of transport of species is also reduced because of (N_b) and (N_t) as can be observed in Table 2. Table 3 discusses the comportment of numerous influential parameters with respect to the rate of mass transportation. Table 4

exhibits the decreasing behavior of density of gyrotactic microorganisms at large values of Pe.

4. Key findings of performed analysis

This study investigated the transfer of thermal energy in a different light and conducted a solubility analysis for the magnetohydrodynamic flow of Carreau rheology with nanoparticles and motile microorganisms via heated disk. OHAM was applied to capture the analytic solution of fluid flow phenomenon. The characteristics of the impacts of parameters on flow, heat energy, concentration of motile microorganisms, and nanomaterial observations were conducted graphically. The valuable observation is summarized as follows:

- Increasing the values of (n) had an inverse role in fluid flow due to shear thickening and shear thinning;
- Escalating values of Hartmann and Carreau fluid numbers reduced fluid flow;
- Prandtl and Brownian motion numbers had the opposite effect on fluid temperature;
- Large values of Prandtl number reduced thermal energy field and connected layer;
- The maximum heat was generated upon increasing values of thermophoresis and Brownian motion numbers;
- Enhancement of numerical values of N_b results slowed down the transport of mass species;
- The transport rate for mass species also increased compared to the large values of Prandtl number;
- $\xi(\eta)$ decayed due to the difference in the concentration of microorganisms (Ω) and enhancement of Peclet number (Pe);
- The gradient velocity increased as opposed to the increased values of Hartmann number;
- Large values (N_b) led to a reduction in heat energy (Nu_r) and transport of species accelerated due to large values of N_b ;
- Augmenting values of Peclet number reduced the density number.

Acknowledgement

The research was supported by the National Natural Science Foundation of China (Grant Nos. 11971142, 11871202, 61673169, 11701176, 11626101, 11601485).

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