Research Note

A numerical investigation of synthetic jet effect on dynamic stall control of oscillating airfoil

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KEYWORDS
Dynamic stall; Oscillating airfoil; Flow control; Synthetic jet actuator; Numerical solution; Momentum coefficient.

Abstract. At high angles of attack, the dynamic stall phenomenon can arise from the vortex shedding, particularly in an oscillating airfoil. As a result of this phenomenon, a considerable decrease in the lift and an increase in the drag and pitching moment coefficients are observed. This study aims to investigate the flow control of an NACA 0015 airfoil using a Synthetic Jet (SJ). The flow was assumed to be unsteady and turbulent at the Mach number of 0.2 and Reynolds number of 1 million. This research was conducted at the angle of attack of 15° ± 10°. In order to carry out the numerical analysis of the problem, the 2D compressible turbulent Navier-Stokes equations based on “Roe” scheme with the second-order accuracy were solved. Turbulence modeling was carried out using the three-equation $k − k_1 − \omega$ model. According to the obtained results, this flow control method could significantly control or eliminate the dynamic stall of the airfoil. In addition, the phase difference between the jet and airfoil oscillations was mostly affected by the dynamic stall decrement. In these changes, using SJ with a momentum coefficient of 0.1 brought about the amplitude of maximum lift at $\varphi = -30^\circ$, and the multiplication of the coefficients of drag and moment amplitudes at $\varphi = -30^\circ$ ensured the best performance.

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1. Introduction

Dynamic stall has many severe consequences and it should be anticipated as quickly as possible. At high angles of attack, flow separation may occur. In addition, when flow is not controlled, dynamic stall occurs at a specific angle. As a result of uncontrolled flow, the lift and drag forces would suddenly decrease and increase, respectively, and the wing loses stability. Thus, flow control is of significance when it comes to preventing such difficulties. In recent years, many experimental and numerical investigations have been conducted to propose an acceptable method that can inhibit the separation flow and control the stall [1-4]. Duvigneau and Visonneau [5] numerically studied the effects of Synthetic Jet (SJ) control on the NACA 0015 airfoil at $Re = 8.96 \times 10^5$. They considered the jet frequency of 0.748, non-dimensional jet velocity of 1.72, and the inclined angle of 25° in their study. The effect of a tangential SJ on aerodynamic characteristics of a NACA 23012 airfoil was also investigated by Esmaeili et al. [6]. They concluded that at the chord Reynolds number of $Re = 2.19 \times 10^6$, two jet oscillating frequencies with different blowing ratios could be obtained and it can be stated that the activation of the SJ could

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control the stall characteristics of the airfoil. Zhang et al. [7] also investigated the effect of suction control on NACA 0012 through LES methods and concluded that by increasing the suction coefficient, the lift-drag ratio would first increase and then decrease. In addition, they estimated the location of suction and found that the only area with a considerable control effect on flow separation and lift increase was behind the separation point. Tran et al. [8] investigated the ability of dynamic large eddy simulation to predict the flow interactions of a finite-span SJ on NACA 4421 airfoil. Moreover, they compared the results of the large eddy simulation with those of previous experiments and direct numerical simulations. Montazer et al. [9] conducted a numerical study that investigated the effect of the SJ on NACA 0015 at Reynolds number $Re = 896000$. They aimed to optimize the jet implementation to improve the aerodynamic characteristics of the airfoil. The results of their study indicated that the jet implementation was the most useful technique for post-stall angle and could increase the Lift-to-Drag ($L/D$) by 66%. Moreover, Tran et al. [10] numerically investigated the effect of the SJ on S809 airfoil by considering the low-energy input requirements. According to their study, jet control, installed near the leading edge, could reduce the flow separation and, consequently, reduce the hysteresis by 73%. Yousefi et al. [11] studied the effects of blowing and suction flow control on NACA 0012. In fact, they explored the effects of the width jet. It was observed that the lift-to-drag ratio could be improved with an increase in the suction and blowing jet width. Furthermore, Moshtagh and Hur [12] investigated the effect of SJ on S809 airfoil numerically using Detached-Eddy Simulation (DES) turbulence model. At a small angle of attack, the jet was subject to early separation and the lift coefficient was reduced. For the separated flow, the enhancement of the aerodynamic coefficients was observed. Zhao and Zhao [13] numerically investigated flow control around an OA213 rotor with a jet. Furthermore, focusing on a wind tunnel, Tang et al. [14] investigated the effect of the SJ on the low-speed airfoil. They implemented SJ in their proposed model in which the maximum lift coefficient was increased by 27.4% and the drag coefficient was decreased by 19.6%. In addition, Giorgi et al. [15] compared and analyzed the effect of using two different flow control methods, namely SJ and Continuous Jet (CJ), on the boundary layer separation on a NACA 0015 airfoil. They concluded that the Synthetic Jet Actuator (SJA) was more useful in terms of regaining energy. Abe et al. [16] also applied the large-eddy method to conduct a simulation of installing a SJ at the leading edge on NACA 0015 airfoils and investigated the effects of actuation frequency. Neve et al. [17] also carried out parametric analysis to investigate the effect of frequency, jet angle, and jet velocity on the NACA 0015 at Reynolds number of $896000$. The jet angle (30°–40°), jet frequency (100 Hz), and non-dimensional jet velocity (1.8–2.0) could significantly affect the performance. Parthasarathy and Das [18] analyzed the physics of the flow and controlled the separated flow at 20° angle of attack on the NACA 0015 airfoil at the Reynolds number of 896000 using SJ.

Both pitching airfoil and SJ mechanisms can be regarded as periodic functions. When these two mechanisms are simultaneously used, the phase difference between these two oscillations can affect the flow field. According to the previous research studies, the effect of phase difference has not been carefully investigated yet. In this study, an active flow control based on the SJ was applied to the NACA 0015 oscillating airfoil. The present study aims to investigate the effects of SJ on the dynamic stall control and characteristics of aerodynamic amelioration. Furthermore, the effect of phase difference between the airfoil and SJ oscillations on the aerodynamic characteristics was investigated.

To this end, an in-house code based on the Reynolds-averaged Navier-Stokes equations for the unsteady and turbulent flow was developed.

2. Numerical methods

2.1. Governing equations

The integral form of the two-dimensional compressible Navier-Stokes equations is described in the following [19]:

$$\frac{\partial}{\partial t} \int_{O} \mathbf{W} d\Omega + \oint_{O} \left( \mathbf{F}_{e} - V_{S} \mathbf{W} - \mathbf{F}_{v} \right) dS = 0, \quad (1)$$

where $\Omega$, $\partial \Omega$, and $V_{S}$ are the moving control volume, control surface, and control volume speed, respectively [20]. The following conservative variables, convection, and viscous flux are given below:

$$\mathbf{W} = \begin{bmatrix} \rho \\ \rho u \\ \rho v \\ \rho E \end{bmatrix}, \quad \mathbf{F}_{e} = \begin{bmatrix} \rho V_{r} \\ \rho u V_{r} + \rho n x p \\ \rho v V_{r} + \rho n y p \\ \rho H V_{r} + V_{S} p \end{bmatrix},$$

$$\mathbf{F}_{v} = \begin{bmatrix} m_{x} x_{x} + m_{y} x_{y} \\ m_{x} y_{x} + m_{y} y_{y} \\ n_{x} x_{x} + n_{y} x_{y} \end{bmatrix}, \quad (2)$$

where $V_{r}$ is the relative velocity of the motion of flow and system [21]. The static pressure ($p$) is also written as follows:

$$p = \rho (\gamma - 1) \left[ E - \left( \frac{u^{2} + v^{2}}{2} \right) \right],$$

$$p = \rho (\gamma - 1) \left[ E - \left( \frac{u^{2} + v^{2}}{2} \right) \right], \quad (3)$$
where $E$ is the total energy per unit mass [22]. The shear stress components as well as $\Theta_x$ and $\Theta_y$ are expressed in [19]. This study investigates unsteady and turbulent flow of the airfoil at the chord Reynolds number (Reynolds number based on the airfoil chord length) of $10^6$. The chord Reynolds number is defined as follows:

$$\text{Re} = \frac{U_\infty c}{v},$$

where $U_\infty$ is the free-stream velocity, $c$ is the chord length, and $v$ is the kinematic viscosity of the fluid [23]. In order to simulate the turbulent flow of interest, the widely known $k - \varepsilon$ model [24] was employed. The abovementioned model consists of three transport equations and three transport equations including the turbulent kinetic energy $k_T$, laminar kinetic energy $k_L$, and specific dissipation rate $\omega$ [24].

For the time discretization of Eq. (1), an explicit scheme, as expressed in Eq. (5), was utilized. The fourth-order Runge-Kutta method was used to solve Eq. (5):

$$\tilde{W}^{n+1} - \tilde{W}^{n} = -\frac{\Delta t}{\Omega} \tilde{R}^{n},$$

where $\Delta t$, $n$, and $\tilde{R}$ are the global physical time step, time level, and residual vector, respectively [20]. In order to discretize the residual vector, a finite-volume scheme based on the second-order approximate Roe was employed [25]. On a structured grid, $\Delta t$ for a control volume $\Omega$ could be obtained through the following approximate relation [19]:

$$\Delta t = CFL \frac{\Omega_t}{(\xi_t + \xi_t)},$$

where the spectral radii of the convective flux Jacobians are written as follows:

$$\xi_t = (|V|^t + \alpha |S|^t),$$

$$\xi_t = (|V|^t + \alpha |S|^t).$$

After conducting the time-step independency study, a minimum $\Delta t$ over all control volumes was selected to achieve time accuracy.

The following equation is used to simulate the motion of oscillating airfoil [26]:

$$a(t) = a_m + a_0 \sin(\omega t),$$

where $a_m$, $a_0$, and $\omega$ are the main angle of attack, angular amplitude, and angular frequency, respectively. The angular frequency depends on the reduced frequency and is defined as follows:

$$k = \omega c/2U_\infty.$$

### 2.2. Grid generation and boundary conditions

A C-type grid around the NACA 0015 airfoil is generated. This grid is archived with a combination of an o-type grid in the upstream zone and l-type grid in the downstream zone. By utilizing a C-type structure, a proper orthogonal grid, particularly near the leading and trailing edges, can be generated, as shown in Figure 1 [27].

In order to simulate the grid motion, the coordinate system origin was fixed on the one-quarter of the airfoil chord from the leading edge and the airfoil along with the computational domain oscillates around this point. Figure 2 shows the computational domain and applied boundary conditions. The grid domain includes a velocity inlet boundary, a pressure outlet boundary, and a solid wall (airfoil surface). The velocity of the SJ is described as follows [5]:

$$u_j = U_j \sin(ft + \phi),$$

where $U_j$, $\phi$, and $f$ are the jet velocity amplitude, phase difference between the airfoil and jet, and non-dimensional frequency, respectively. The non-dimensional frequency was expressed using the following equation [13]:

$$f = \frac{\omega_j c}{2U_\infty},$$

where $\omega_j$ is the oscillation frequency of the SJ. The

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**Figure 1.** A part of the grid used in flow computations.

**Figure 2.** Computational domain and applied boundary conditions around NACA 0015 airfoil.
ratio between the momentum of the free stream and the momentum of the jet is defined as the momentum coefficient \( c_\mu \), as shown in the following [13]:

\[
c_\mu = C_\mu \sin (f \cdot t + \varphi);
\]

\[
C_\mu = \frac{\rho j U_j^2 h}{1/2 \rho \infty U_\infty^2 c},
\]  \( \text{(13)} \)

where \( C_\mu \) is the amplitude of the SJ and \( h \) is the SJ throat width.

3. Results

3.1. Grid independence and validation

In order to evaluate the effects of the grid size on the results obtained from the numerical solutions, three grids including \( 285 \times 51 \), \( 316 \times 61 \), and \( 351 \times 66 \) cells were tested. Figure 3 shows the lift coefficients of these three grids. As observed, the results of grids 2 and 3 are very close to each other. Due to a large amount of the computational cost of an unsteady solution, grid 2 is used for the subsequent computations. The first cell size of the selected grid is set so that \( y^+ < 1 \).

In order to evaluate the precision of the developed computer program, the present results were compared with those of previous studies on the lift and drag coefficients. First, the ability of the flow solver was tested for an oscillating airfoil case without SJ. To this end, the flow parameters \( C_l \) and \( C_d \) of the present study and the experimental results obtained by Piziali [28] were compared. A comparison was made in terms of a NACA 0015 airfoil at the Reynolds number of \( 1.935 \times 10^6 \) with \( M = 0.289 \) and \( k = 0.134 \) in the of angle of attack range of \( 17^\circ \pm 4^\circ \). The results of these comparisons are shown in Figures 4–6. The anticipations are in close agreement with the experimental data. After evaluating the flow solver for the baseline airfoil without SJ, the lift coefficients of the stationary airfoil with the SJ of the present flow solver were compared with numerical results of Duvigneau and Visonneau [6]. This comparison was made concerning a NACA 0015 airfoil at the Reynolds number of \( 2 \times 10^6 \) with \( M = 0.289 \) and \( k = 0.134 \) in the of angle of attack range of \( 17^\circ \pm 4^\circ \). The results of these comparisons are shown in Figures 4–6. The anticipations are in close agreement with the experimental data.
airfoil at the Reynolds number of 890000 with $Ma = 0.1$. Figure 7 compares the results of the present solver with those of Duvigneau and Viscomante [5] for the lift coefficients with respect to the angle of attack.

3.2. Simulating the SJ on the airfoil
In this study, the slot of the tangential SJ was placed on the upper surface of the NACA 0015 airfoil centering at 0.5% chord with a height of 0.25% chord. Figure 8 shows the situation of the SJ on the airfoil leading edge. The investigation was conducted at $Ma = 0.2$ and $Re = 10^6$. In the following section, the effects of the investigated parameters of the SJ on the aerodynamic coefficients are studied. These parameters include the magnitude of momentum coefficient, reduced frequency, and phase difference.

3.3. The effect of momentum coefficient
The current study aims to explore the changes in the aerodynamic characteristics with variations in the jet momentum coefficient. To this end, a comparison was made to examine the three jet momentum coefficients including $C_p = 0.07$, 0.1, and 0.13. The results of the comparison of these three SJ control cases are shown in Figures 9–11. These investigations were performed at the reduced frequency of $k = 0.25$ and $f = k$. The results also indicated that SJ with the higher momentum coefficient ($C_p = 0.13$) had a better ability in lift enhancement. In addition, by increasing

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**Figure 7.** Comparison of the present solver and experimental data [5] regarding the lift coefficients and angle of attack.

**Figure 8.** The situation of the Synthetic Jet (SJ) on the leading edge of airfoil.

**Figure 9.** Comparison of the lift coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with different momentum coefficients.

**Figure 10.** Comparison of the drag coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with different momentum coefficients.

**Figure 11.** Comparison of the drag coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with different pitching moment coefficients.
the jet momentum coefficient, the hysteresis loops of aerodynamic coefficients became thinner. Figure 10 reveals that the SJ control case with \( C_\mu = 0.07 \) has the maximum drag coefficient at 25\(^\circ\) and the aerodynamic coefficient loops indicate larger hysteresis in the lift coefficient curve than in other cases. According to the findings, by increasing the jet momentum, more energy can be transferred to the boundary layer. Consequently, further improvements were carried out to increase the lift and reduce the drag coefficients. For the case with \( C_\mu = 0.1, 0.13 \), the lift coefficient curves do not show significant stall, indicating that a stronger jet is able to completely control the dynamic stall. In order to quantify the improvement of the results of SJ control for enhancing the lift coefficient and the decrease in the drag and pitching moment coefficients, the differences in the area under the \( C_l, C_d, \) and \( C_m \) curves among the control cases and the baseline airfoils were calculated [29] as follows (where \( q \) is either drag or moment):

\[
\Delta A C_q = \frac{\int_0^{2\pi} (C_{q,\text{baseline}} - C_{q,\text{control}}) d\alpha}{\int_0^{2\pi} C_{q,\text{baseline}} d\alpha}.
\]

3.4. The effect of phase difference at \( k = 0.25 \) and \( C_\mu = 0.1 \)

Another control parameter of the jet is responsible for investigating the aerodynamic characteristics of the NACA 0015 airfoil with many phase differences. Given that airfoil and jet oscillations are both sinusoidal, one can consider a phase difference between them to study this effect. In fact, when the oscillating airfoil is at the maximum angle of attack, the oscillating jet will have the highest jet velocity; thus, the phase difference will be zero. Otherwise, there will be a phase difference between these oscillations. Figures 12–23 show the modifications of lift, drag, and pitching moment coefficients compared to the baseline under four phase differences. In all these cases, the same jet

![Figure 12. Comparison of the lift coefficients with respect to \( \alpha \) for Synthetic Jet (SJ) with \( \varphi = 10^\circ \).](image)

| Table 1. Comparison of the effects of the synthetic jet cases and their baseline airfoils \((k = 0.25)\). |
|---|---|---|---|---|
| Coefficient | Baseline | \( C_\mu = 0.07 \) | \( C_\mu = 0.1 \) | \( C_\mu = 0.13 \) |
| \( C_{l,\text{amp}} \) | 1.07 | 0.88 | 1.23 | 1.32 |
| \( C_{d,\text{amp}} \) | 0.2 | 0.21 | 0.155 | 0.149 |
| \( C_{m,\text{amp}} \) | -0.043 | -0.037 | -0.019 | -0.02 |
| \( \Delta A C_l \) | — | 17.76\% | 14.95\% | 23.36\% |
| \( \Delta A C_d \) | — | 5\% | 22.5\% | 25.5\% |
| \( \Delta A C_m \) | — | 13.95\% | 55.81\% | 53.49\% |
Figure 13. Comparison of the drag coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 10^\circ$.

Figure 14. Comparison of the pitching moment coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 10^\circ$.

Figure 15. Comparison of the lift coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 30^\circ$.

Figure 16. Comparison of the drag coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 30^\circ$.

Figure 17. Comparison of the pitching moment coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 30^\circ$.

Figure 18. Comparison of the lift coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = -10^\circ$.

momentum $C_p = 0.1$ with the reduced frequency of $k = 0.25$ and $f = k$ was considered. Similar to the previous section, to quantify the enhancement in the lift and the decrease in the drag and pitching moment during a pitch cycle, the differences of the area under the $C_l$, $C_d$, and $C_m$ curves among the SJ control cases and their baseline airfoils were estimated. Table 2 shows the modifications of the aerodynamic characteristics with many phase differences. As observed, with $\varphi = -30^\circ$, the lift would increase by 20.56% and with $\varphi = -10^\circ$, the drag and pitching moment are reduced by 25.5% and 60.46%, respectively. When the phase differences vary from $-10^\circ$ to $-30^\circ$, the amount of amplitude drag and pitching moment would reduce from 25.5% to
Figure 19. Comparison of the drag coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = -10^\circ$.

Figure 20. Comparison of the pitching moment coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = -10^\circ$.

Figure 21. Comparison of the lift coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = -30^\circ$.

Figure 22. Comparison of the drag coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = -30^\circ$.

Figure 23. Comparison of the pitching moment coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = -30^\circ$.

SJ control case with $\varphi = 30^\circ$ does not outperform the baseline because the multiplication of the amplitudes drag and moment coefficients shows a 5.8 increase with respect to the baseline. Moreover, the multiplication of the amplitudes of drag and moment coefficients in the SJ control case with $\varphi = -10^\circ$ ensures the highest improvement.

3.5. The effect of phase difference at $k = 0.15$ and $C_\mu = 0.1$

This section aims to investigate the effects of SJ control with a reduced frequency of $k = 0.15$ under two phase differences ($\varphi = 0^\circ$ and $\varphi = -30^\circ$). Figures 24–29 show the results obtained from comparing the SJ cases and their baselines regarding the lift, drag, and pitching moment coefficients. Table 3 presents the amplitudes of the aerodynamic coefficients of the many phase differences and their enhancement compared to the baselines. It also indicates that in the SJ control case with $\varphi = -30^\circ$, the amplitude of lift would increase by 27.98%. Furthermore, the amplitude of drag and pitching moment would reduce by 14.98% and 23.08%, respectively. The results indicate that at the
Table 2. Comparison of the effects of the SJ actuation with $C_m = 0.1$ and $k = 0.25$ and the baseline.

<table>
<thead>
<tr>
<th>Coefficient</th>
<th>Baseline</th>
<th>$\varphi = 0^\circ$</th>
<th>$\varphi = 10^\circ$</th>
<th>$\varphi = 30^\circ$</th>
<th>$\varphi = -10^\circ$</th>
<th>$\varphi = -30^\circ$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$C_{l,amp}$</td>
<td>1.07</td>
<td>1.23</td>
<td>1.236</td>
<td>1.25</td>
<td>1.22</td>
<td>1.29</td>
</tr>
<tr>
<td>$C_{d,amp}$</td>
<td>0.2</td>
<td>0.155</td>
<td>0.169</td>
<td>0.205</td>
<td>0.149</td>
<td>0.15</td>
</tr>
<tr>
<td>$C_{m,amp}$</td>
<td>-0.043</td>
<td>-0.019</td>
<td>-0.027</td>
<td>-0.044</td>
<td>-0.017</td>
<td>-0.025</td>
</tr>
<tr>
<td>$\Delta C_l$</td>
<td>—</td>
<td>14.95%↑</td>
<td>15.51%↑</td>
<td>16.82%↑</td>
<td>14.01%↑</td>
<td>20.56%↑</td>
</tr>
<tr>
<td>$\Delta A_{C_d}$</td>
<td>—</td>
<td>22.5%↓</td>
<td>15.5%↓</td>
<td>2.5%↑</td>
<td>25.5%↓</td>
<td>25%↓</td>
</tr>
<tr>
<td>$\Delta A_{C_m}$</td>
<td>—</td>
<td>55.81%↓</td>
<td>37.21%↓</td>
<td>2.32%↑</td>
<td>60.46%↓</td>
<td>41.86%↓</td>
</tr>
<tr>
<td>$\Delta A_{C_l} \times \Delta A_{C_m}$</td>
<td>—</td>
<td>1255.7↓</td>
<td>576.75↓</td>
<td>5.8↓</td>
<td>1541.7↓</td>
<td>1046.5↓</td>
</tr>
</tbody>
</table>

Figure 24. Comparison the lift coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 0^\circ$.

Figure 25. Comparison the drag coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 0^\circ$.

Figure 26. Comparison the pitching moment coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = 0^\circ$.

Figure 27. Comparison the lift coefficients with respect to $\alpha$ for Synthetic Jet (SJ) with $\varphi = -30^\circ$.

$-30^\circ$ phase, i.e., the difference between the phases of the airfoil and jet, the amplitudes of drag and moment coefficient were reduced significantly more than those at the phase $0^\circ$. However, the lift coefficient did not considerably increase compared to that at the phase $0^\circ$.

4. Conclusions

In the present study, the dynamic stall control was numerically investigated using a Synthetic Jet (SJ) on the NACA 0015 at $Re = 10^6$. The solver was validated against the results of the baseline experiment regarding the oscillating airfoil. The comparison of the results of SJ control and those of the numerical investigation of the static airfoil showed good agreement. First, the effects of varying the jet momentum coefficients on the dynamic stall control performance were investigated. The results indicated that using the SJ control with proper momentum coefficient could considerably control the separation. As a result, the dynamic stall was delayed or arrested. Three momentum coefficients of 0.07, 0.1, and 0.13 were investigated. The SJ cases
Table 3. Comparison of the SJ cases with the baseline.

<table>
<thead>
<tr>
<th>Coefficient</th>
<th>Baseline</th>
<th>(\varphi = 0^\circ)</th>
<th>(\varphi = -30^\circ)</th>
</tr>
</thead>
<tbody>
<tr>
<td>(C_{1,amp})</td>
<td>0.965</td>
<td>1.18</td>
<td>1.235</td>
</tr>
<tr>
<td>(C_{d,amp})</td>
<td>0.188</td>
<td>0.186</td>
<td>0.16</td>
</tr>
<tr>
<td>(C_{m,amp})</td>
<td>-0.039</td>
<td>-0.042</td>
<td>-0.03</td>
</tr>
<tr>
<td>(\Delta A_{C_l})</td>
<td>-</td>
<td>22.28% ↓</td>
<td>27.98% ↓</td>
</tr>
<tr>
<td>(\Delta A_{C_d})</td>
<td>-</td>
<td>1.06% ↓</td>
<td>14.89% ↓</td>
</tr>
<tr>
<td>(\Delta A_{C_m})</td>
<td>-</td>
<td>7.69% ↑</td>
<td>23.08% ↓</td>
</tr>
</tbody>
</table>

Figure 28. Comparison the drag coefficients with respect to \(\alpha\) for Synthetic Jet (SJ) with \(\varphi = -30^\circ\).

Figure 29. Comparison the pitching moment coefficients with respect to \(\alpha\) for Synthetic Jet (SJ) with \(\varphi = -30^\circ\).

with the momentum coefficients of \(C_{m} = 0.1\) and 0.13 exhibited an acceptable performance in eliminating the dynamic stall onset. The lower momentum of the jet required lower energy consumption. Therefore, the jet with the momentum coefficient of 0.1 was utilized to continue the investigation. The oscillations of the airfoil and jet were both sinusoidal. Thus, the effect of phase difference among them might be considerable. This effect has not been addressed in previous studies. To this end, in order to evaluate the effect of phase difference between the jet and airfoil, a range of phase differences between \(-30^\circ\) and \(+30^\circ\) were studied. Furthermore, the effect of the SJ at two different reduced frequencies was investigated. The conclusions and improvements of this research versus the baseline airfoils are given below:

1. For the case with the reduced frequency of \(k = 0.25\) and phase difference of \(\varphi = 0^\circ\), the amplitude of lift could increase by 14.95% and the amplitude of drag and pitching moment could decrease by 22.5% and 55.81%, respectively;

2. The results showed that the case with the reduced frequency of \(k = 0.25\) had the highest amplitude of lift at \(\varphi = -30^\circ\), which was improved by 20.56%. Furthermore, at \(\varphi = -10^\circ\), the drag amplitude and pitching moment coefficients were reduced by 25.5% and 60.46%, respectively, which presented the best performance and a considerable decrease;

3. For the case with a lower reduced frequency of \(k = 0.15\), the best aerodynamic performances were achieved at \(\varphi = -30^\circ\). In this phase difference, the amplitude of lift was increased by 27.98%. In addition, the amplitude of drag and the amplitude of pitching moment were reduced by 14.98% and 23.08%, respectively.

Nomenclature

- \(\alpha\): Speed of sound
- \(c\): Airfoil chord
- \(C_l\): Lift coefficient
- \(C_d\): Drag coefficient
- \(C_m\): Pitching moment coefficient
- \(k\): Reduced frequency
- \(f\): Synthetic jet forcing frequency
- \(Re\): Reynolds number
- \(h\): Actuation surface
- \(Ma\): Mach number
- \(C_\mu\): Momentum coefficient

Greek

- \(\alpha\): Angle of attack
- \(\mu\): Dynamic viscosity
- \(\rho\): Density
- \(\eta\): Similarity variable
- \(\varphi\): Phase difference

Subscripts

- \(\infty\): Free stream (far field)
- \(t\): Turbulent

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Biographies

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