

An analytic and mathematical synchronization of micropolar nanofluid by Caputo-Fabrizio approach

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Abstract. Nanofluids and enhancement of the heat transfer in real systems have proved to be a widely a research area of nanotechnology, specially, improvement in thermal conductivity, thermophoresis phenomenon, dispersion of nanoparticles volume fraction and few others. Based on the touch of nanotechnology, this research article investigates the analytic and mathematical performance of micropolar nanofluid for the enhancement of heat transfer. The base fluid is taken for the purpose of thermal conductivity subject to two types of nanoparticles so called copper and silver. The mathematical analysis of the micropolar nanofluid has been carried out by invoking non-integer order derivative and transforms methods. By applying mathematical tool on the equations of micropolar nanofluid, the solutions have been explored for temperature, microrotation and velocity. In order to meet the physical aspects of the problem based on micropolar nanofluid, the comparison of velocity field of micropolar nanofluid for the suspensions of ethylene glycol into silver and ethylene glycol into copper is characterize for to enhance the rate of heat transfer.

Key word: Heat transfer of micropolar nanofluids, Fractional derivative of non-singular kernel, Special Fox-**H** function, Suspension of nanoparticles in base fluid.

1. Introduction

Nanoparticles are the small quantities of nanometer-size particles which are composed of metals for instance, copper oxide, alumina, carbides, titania, gold, copper and several others. These nanoparticles have capability to enhance thermophysical properties of base fluids includes ethylene glycol, biofluids, oil, water, lubricants, polymer solutions and many others. In continuation, kerosene, engine oil, water and ethylene glycol are insufficient for heat transfer because they have low thermal conductivity. In order to enhance convective heat transfer performance of such fluids, various techniques have been implemented such as boundary conditions and changing flow geometries. There is no denying fact that fluids have lesser thermal conductivity than metals. The thermal conductivity of base fluids can be enhanced by adding the metals and consequentially

such fluids are characterized as nanofluids [1]. Moreover, the fundamental concepts of micropolar fluid were originated by Eringen [2] that characterize the dynamics of such fluids by considering microscopic impacts rising from micro-motion and local structure of the fluids. A broad assess on micropolar fluid with its engineering applications has investigated by Ariman et al. [3, 4]. Hassanien and et al. [5] explored effects of nonisothermal stretching sheet on micropolar with heat transfer, blowing and suction. Mohammadein and et al. [6] analyzed internal heat generation and dissipation with heat transfer in micropolar. Hussanan and et al. [7] obtained closed form solution for Newtonian heating in micropolar fluid for the problem of free convection flow. Choi [8] studied applications to non-Newtonian fluid flow with nanoparticles to enhance thermal conductivity of fluid. Congedo and et al. [9] perused analysis and modeling of nanofluids with natural convection heat transfer. Ghasemi and et al. [10] traced out impacts of water-CuO nanofluid along with natural convection heat transfer. Hussanan and et al. [11] examined influence of some nanofluids for accelerated plate with magnetic field and porous medium. Water functionalized carbon nanotube flow was analyzed over a moving/static wedge with magnetic field by Khan and et al. [12]. Haq and et al. [13] explored effects of magnetic field with water functionalized metallic nanoparticles for squeezed flow over a sensor surface. In brevity, we include here few very latest references on nanofluids [14-16], modern fractional derivatives [17-22], heat transfer [23-27], nanoparticles [28-31], porosity and magnetic field [32-37] and few different circumstances [38-43]. Motivating by above research work on nanofluids, our purpose is to investigate the analytic and mathematical performance of micropolar nanofluid for the enhancement of heat transfer. The base fluid is taken for the purpose of thermal conductivity subject to two types of nanoparticles so called copper and silver as shown in Fig. 1(a)-1(c). The mathematical analysis of the micropolar nanofluid has been carried out by invoking non-integer order derivative and transforms methods. By applying mathematical tool on the equations of micropolar nanofluid, the solutions have been explored for temperature, microrotation and velocity. In order to meet the physical aspects of the problem based on micropolar nanofluid, the comparison of velocity field of micropolar nanofluid for the suspensions of ethylene glycol into silver and ethylene glycol into copper is characterize for to enhance the rate of heat transfer.

2. Mathematical Equations of Nanofluid

An unsteady flow of ethylene glycol based on micropolar nanofluid occupying the space lying over an oscillating plate which is perpendicular to the y-axis and is situated in the (x, z) plane. Initially, due to constant temperature T_w the fluid is considered at rest. At $t=0^+$, the plate starts to oscillate with the velocity $UH(t)\cos\omega t$ or $U\sin\omega t$ in its plane and level of temperature is increased up to T_w . The governing equations for micropolar fluid are

$$\nabla \cdot (\rho_{nf} \mathbf{V}) = \frac{\partial \rho_{nf}}{\partial t}, \quad (1)$$

$$\nabla \cdot p + \nabla \times (\nabla \times \mathbf{V})(K_1 + \mu_{nf}) - \nabla(\nabla \cdot \mathbf{V})(K_1 + 2\mu_{nf}) = \rho_{nf} \mathbf{b} - \rho_{nf} \frac{d\mathbf{V}}{dt} + K_1 (\nabla \times \mathbf{N}), \quad (2)$$

$$2K_1 \mathbf{N} + \gamma_{nf} \nabla \times (\nabla \times \mathbf{N}) - \nabla(\nabla \cdot \mathbf{N})(\gamma_{nf} + \lambda + \alpha) = \rho_{nf} \mathbf{I} - \rho_{nf} j \frac{d\mathbf{N}}{dt} + K_1 (\nabla \times \mathbf{V}), \quad (3)$$

Where, ρ_{nf} , \mathbf{V} , p , K_1 , μ_{nf} , \mathbf{b} , \mathbf{N} , \mathbf{I} , λ , α , j , γ_{nf} are nanofluid density, velocity field, pressure, vortex viscosity, nanofluid dynamic viscosity, body force vector, (gyration) microrotation vectors, body couple per unit mass vector, Spin gradient viscosity coefficients, micro-inertia density, spin-gradient viscosity respectively. Under the constant viscosity, equations (1-3) can be converted in terms of Navier-Stokes equations for micropolar fluid. The Brinkman [44] has stated relationship between base fluid and dynamic viscosity of nanofluid written below

$$\mu_f = \mu_{nf} (1 - \phi)^{2.5}. \quad (4)$$

Aminossadati and et al. [45] and Matin and et al. [46] have described the viscosity of nanofluid in terms as

$$\phi \rho_s + \rho_f (1 - \phi) = \rho_{nf}. \quad (5)$$

The expression of γ_{nf} is based according to published work by Bourantas and et al. [47]

$$j \left(\frac{K_1}{2} + \mu_{nf} \right) = \gamma_{nf}. \quad (6)$$

The continuity equation for incompressible flow is

$$\nabla \cdot \mathbf{V} = 0. \quad (7)$$

Using the vector identity and by neglecting body couple force, for free convection flow equations (2-3) are expressed as

$$\nabla \cdot p - \nabla(\nabla \cdot \mathbf{V})\mu_{nf} - \nabla^2 \cdot \mathbf{V}(K_1 + \mu_{nf}) = \rho_{nf} \mathbf{g} - \rho_{nf} \frac{d\mathbf{V}}{dt} + K_1 (\nabla \times \mathbf{N}), \quad (8)$$

$$2K_1 \mathbf{N} - \gamma_{nf} \nabla^2 \mathbf{N} - \nabla(\nabla \cdot \mathbf{N})(\lambda + \alpha) = K_1 (\nabla \times \mathbf{V}) - \rho_{nf} j \frac{d\mathbf{N}}{dt}, \quad (9)$$

Simplifying equations (8-9) take place by using mass conservation as

$$\nabla \cdot p - \nabla^2 \cdot \mathbf{V}(K_1 + \mu_{nf}) - \rho_{nf} \mathbf{g} = K_1 (\nabla \times \mathbf{N}) - \rho_{nf} \frac{d\mathbf{V}}{dt}, \quad (10)$$

$$2K_1 \mathbf{N} - \gamma_{nf} \nabla^2 \mathbf{N} - \nabla(\nabla \cdot \mathbf{N})(\lambda + \alpha) = K_1 (\nabla \times \mathbf{V}) - \rho_{nf} j \frac{d\mathbf{N}}{dt}, \quad (11)$$

In order to apply the statement of material derivative, equations (10-11) are expressed equivalently as

$$\nabla \cdot p - \nabla^2 \cdot \mathbf{V}(K_1 + \mu_{nf}) - \rho_{nf} \mathbf{g} = K_1 (\nabla \times \mathbf{N}) - \rho_{nf} \left(\mathbf{V}(\nabla \cdot \mathbf{V}) + \frac{d\mathbf{V}}{dt} \right), \quad (12)$$

$$2K_1 \mathbf{N} - \gamma_{nf} \nabla^2 \mathbf{N} - \nabla(\nabla \cdot \mathbf{N})(\lambda + \alpha) = K_1 (\nabla \times \mathbf{V}) - \rho_{nf} j \left(\mathbf{N}(\nabla \cdot \mathbf{N}) + \frac{d\mathbf{N}}{dt} \right), \quad (13)$$

For the problem under consideration of Cartesian coordinates (x, y, z) , we assume the velocity, microrotation and gravitational fields correspondingly

$$\mathbf{V}(w(y, t), 0, 0), \quad \mathbf{N}(0, 0, N(y, t),) \text{ and } \mathbf{g}(g, 0, 0). \quad (14)$$

Simplifying equations (12-14) gives

$$\frac{\partial p}{\partial x} = \rho_{nf} \frac{\partial w}{\partial x} + K_1 \frac{\partial N}{\partial x} + \rho_{nf} g + (K_1 + \mu_{nf}) \frac{\partial^2 w}{\partial y^2}, \quad (15)$$

$$2K_1 N = K_1 \frac{\partial w}{\partial y} - \rho_{nf} j \frac{\partial N}{\partial x} + \gamma_{nf} j \frac{\partial N}{\partial x}, \quad (16)$$

Implementing Boussinesq approximation on equation (15) and making $K_1 = 0$ in equation (16) respectively, we find that

$$(T - T_\infty) g (\beta_T \rho)_{nf} + (K_1 + \mu_{nf}) \frac{\partial^2 w}{\partial y^2} + K_1 \frac{\partial N}{\partial x} - \rho_{nf} \frac{\partial w}{\partial x} = 0, \quad (17)$$

$$\gamma_{nf} j \frac{\partial^2 N}{\partial y^2} - \rho_{nf} j \frac{\partial N}{\partial t} = 0. \quad (18)$$

Energy equation with thermal radiation as previously published papers [48, 49] is defined as

$$K_{nf} \frac{\partial^2 T}{\partial y^2} - (\rho C)_{nf} \frac{\partial T}{\partial t} - \frac{\partial q_r}{\partial y} = 0, \quad (19)$$

Where, K_{nf} is thermal conductivity of nanofluids and $(\rho C)_{nf}$ is heat capacity at a constant pressure described by Khan et al. [50],

$$\phi(\rho C_p)_s + (\rho C_p)_f (1 - \phi) = (\rho C_p)_{nf}, \quad \frac{K_{nf}}{K_f} = \frac{2K_f + 2\phi(K_s - K_f) + K_s}{2K_f + \phi(K_s - K_f) + K_s}, \quad (20)$$

The imposed conditions are set as

$$T(y, 0) = T_\infty, \quad T(0, t) = T_w, \quad T(\infty, t) = T_\infty, \quad (21)$$

$$N(y, 0) = C_\infty, \quad N(0, t) = t, \quad N(\infty, t) = 0, \quad (22)$$

$$w(y, 0) = 0, \quad w(0, t) = UH(t) \cos(\omega t) \text{ or } U \sin(\omega t), \quad w(\infty, t) = 0. \quad (23)$$

The energy equation (19) can be taken place using Rosseland approximation [51, 52]

$$K_{nf} \left(\frac{16T^3 \sigma^*}{3K_{nf} k^*} \right) \frac{\partial^2 T}{\partial y^2} - \frac{\partial T}{\partial t} (\rho C)_{nf} = 0, \quad (24)$$

Implementing below dimensionless quantities in equations (17-19),

$$T = \frac{T - T_\infty}{T_w - T_\infty}, N^* = \frac{v_f N}{U^2}, w^* = \frac{w}{U}, t^* = \frac{U^2 t}{v_f}, y^* = \frac{U y}{v_f}.$$

$$\text{and } K = \frac{k_1}{\mu_f}, R = \frac{16T_\infty^3 \sigma^*}{3K_f k^*}, G_r = \frac{(T_w - T)_\infty g(\beta_T)_f}{U^3}, P_r = \frac{(C_p)_f \mu_f}{K_f}.$$

(*symbol is dropped for simplicity), we obtain governing partial differential equations for temperature distribution, microrotation field and velocity field by employing appendix (A1-A6) respectively

$$\frac{\partial^2 T}{\partial y^2} = \aleph_0 p_r \aleph_1^{-1} \frac{\partial T}{\partial t}, \quad (25)$$

$$\frac{\partial^2 N}{\partial y^2} = \aleph_2 \aleph_3^{-1} \frac{\partial N}{\partial t}, \quad (26)$$

$$\frac{\partial^2 w}{\partial y^2} = \aleph_2 \aleph_5^{-1} \frac{\partial w}{\partial t} - \aleph_5^{-1} G_r \aleph_4 T - \aleph_5^{-1} k \frac{\partial N}{\partial y}, \quad (27)$$

Here, the assumptions for equations (25-27) are

$$T(y, 0) = 0, \quad T(0, t) = t, \quad T(\infty, t) = 0, \quad (21)$$

$$N(y, 0) = 0, \quad N(0, t) = t, \quad N(\infty, t) = 0, \quad (22)$$

$$w(y, 0) = 0, \quad w(0, t) = UH(t) \cos(\omega t) \text{ or } U \sin(\omega t), \quad w(\infty, t) = 0. \quad (23)$$

Finally, expressing governing equations (25-27) in terms of Caputo-Fabrizio fractional derivative, we have

$$\frac{\partial^2 T}{\partial y^2} = \aleph_1^{-1} p_r \aleph_0 \frac{\partial^\alpha T}{\partial t^\alpha}, \quad (31)$$

$$\frac{\partial^2 N}{\partial y^2} = \aleph_2 \aleph_3^{-1} \frac{\partial^\alpha N}{\partial t^\alpha}, \quad (32)$$

$$\frac{\partial^2 w}{\partial y^2} = \aleph_2 \aleph_5^{-1} \frac{\partial^\alpha w}{\partial t^\alpha} - \aleph_5^{-1} G_r \aleph_4 T - \aleph_5^{-1} k \frac{\partial N}{\partial y}. \quad (33)$$

Where, D_t^ζ or $\frac{\partial^\zeta}{\partial t^\zeta}$ represents fractional operator of Caputo-Fabrizio having order $0 \leq \zeta \leq 1$ [53-55]

defined at the normalization functions are $M(1) = M(0) = 1$ as

$${}^{CF} \left(\frac{\partial^\zeta}{\partial t^\zeta} \right) = {}^{CF} (D_t^\zeta) = \int_0^t \frac{G'(\eta)}{1-\zeta} \text{Exp} \left(-\frac{\zeta(t-\eta)}{1-\zeta} \right) d\eta, \quad \text{for } 0 \leq \zeta \leq 1. \quad (34)$$

3. Investigation of Temperature Distribution and Microrotation field

Using Laplace Transform on Caputo-Fabrizio fractionalized differential equation (31-32) and utilizing the

fact $\lambda = \frac{1}{(1-\alpha)}$ we obtain,

$$\frac{\partial^2 \bar{T}}{\partial y^2} = \aleph_1^{-1} p_r \aleph_0 \frac{\alpha s}{(s + \alpha \lambda)} \bar{T}, \quad (35)$$

$$\frac{\partial^2 \bar{N}}{\partial y^2} = \aleph_3^{-1} \aleph_2 \frac{\alpha s}{(s + \alpha \lambda)} \bar{N}, \quad (36)$$

Expressing equations (35-36) in more suitable format equivalently, we have

$$\bar{T} = s^{-2} \exp \left[-y \sqrt{\frac{p_r \aleph_0 \alpha s}{\aleph_1 s + \aleph_1 \alpha \lambda}} \right], \quad (37)$$

$$\bar{N} = s^{-2} \exp \left[-y \sqrt{\frac{\aleph_2 \alpha s}{\aleph_3 s + \aleph_3 \alpha \lambda}} \right], \quad (38)$$

Reworking on equations (37-38), we obtain summation form as

$$\bar{T} = \frac{1}{s^2} + \sum_{l=1}^{\infty} \frac{1}{l!} \left(-y \sqrt{\frac{\alpha p_r \aleph_0}{\aleph_1}} \right)^l \sum_{m=0}^{\infty} \frac{(-\alpha \lambda)^m \Gamma \left(m + \frac{l}{2} \right)}{m! \Gamma \left(\frac{l}{2} \right) s^{m+2}}, \quad (39)$$

$$\bar{N} = \frac{1}{s^2} + \sum_{l=1}^{\infty} \frac{1}{l!} \left(-y \sqrt{\frac{\alpha \aleph_2}{\aleph_3}} \right)^l \sum_{m=0}^{\infty} \frac{(-\alpha \lambda)^m \Gamma \left(m + \frac{l}{2} \right)}{m! \Gamma \left(\frac{l}{2} \right) s^{m+2}}, \quad (40)$$

Inverting equations (39-40), we expressed the general solutions of temperature and microrotation field in terms of Fox- \mathbf{H} function as

$$T = t + (l!)^{-1} \int_0^t (t - \tau) \sum_{l=1}^{\infty} \left(-y \sqrt{\frac{\alpha p_r \aleph_0}{\aleph_1}} \right)^l \mathbf{H}_{1,3}^{1,1} \left(\alpha \lambda t \right) \left[\begin{matrix} \left(1 - \frac{l}{2}, 1 \right) \\ (0,1), \left(1 - \frac{l}{2}, 0 \right), (0,1) \end{matrix} \right] d\tau, \quad (41)$$

$$N = t + (l!)^{-1} \int_0^t (t - \tau) \sum_{l=1}^{\infty} \left(-y \sqrt{\frac{\alpha \aleph_2}{\aleph_3}} \right)^l \mathbf{H}_{1,3}^{1,1} \left(\alpha \lambda t \right) \left[\begin{matrix} \left(1 - \frac{l}{2}, 1 \right) \\ (0,1), \left(1 - \frac{l}{2}, 0 \right), (0,1) \end{matrix} \right] d\tau, \quad (42)$$

Where, the special function is defined as [56-59]

$$\sum_f \frac{(-\chi)^f \prod_{j=1}^p \Gamma(a_j + A_j f)}{f! \prod_{j=1}^q \Gamma(b_j + B_j f)} = \mathbf{H}_{p,q+1}^{1,p} \left[\chi \left| \begin{matrix} (1-a_1, A_1), (1-a_2, A_2), (1-a_3, A_3), \dots, (1-a_p, A_p) \\ (0,1), (1-b_1, B_1), (1-b_2, B_2), (1-b_3, B_3), \dots, (1-b_q, B_q) \end{matrix} \right. \right]. \quad (43)$$

4. Investigation of Velocity Field

Using Laplace transform on Caputo-Fabrizio fractionalized differential equation (33) and utilizing the fact

$\lambda = \frac{1}{(1-\alpha)}$ we obtain,

$$\frac{\partial^2 \bar{w}}{\partial y^2} = \frac{\aleph_2 \alpha s}{(\aleph_5 s + \aleph_5 \alpha \lambda)} \bar{w} - \aleph_5^{-1} G_r \aleph_4 \bar{T} - \aleph_5^{-1} k \frac{\partial \bar{N}}{\partial y}, \quad (44)$$

Solving the partial differential equation (44) and using initial and boundary conditions (28-30)₃

$$\begin{aligned} \bar{w} = & \frac{Us}{s^2 + \omega^2} \exp \left\{ -y \sqrt{\frac{\aleph_2 \alpha s}{(\aleph_5 s + \aleph_5 \alpha \lambda)}} \right\} + \frac{k \sqrt{\frac{\aleph_2 \alpha}{\aleph_3}}}{\left[\frac{\aleph_5 \aleph_2 \alpha}{\aleph_3} - \aleph_2 \right]} \frac{1}{s^2} \exp \left\{ -y \sqrt{\frac{\aleph_1 \alpha s}{(\aleph_3 s + \aleph_3 \alpha \lambda)}} \right\} \\ & - \frac{G_r \aleph_4}{\left[\frac{p_r \aleph_0 \alpha \aleph_5}{\aleph_1} - \aleph_2 \right]} \frac{1}{s(s + \alpha \lambda)} \exp \left\{ -y \sqrt{\frac{\aleph_0 \alpha p_r s}{\aleph_1}} \right\}, \end{aligned} \quad (45)$$

Reworking on equations (45), we obtain summation form as

$$\begin{aligned} \bar{w} = & \frac{Us}{s^2 + \omega^2} + U \sum_{l=1}^{\infty} \frac{1}{l!} \left(\frac{-y \sqrt{\aleph_2 \alpha}}{\aleph_5} \right)^l \sum_{m=0}^{\infty} \frac{(-\alpha \lambda) \Gamma\left(m + \frac{l}{2}\right)}{m! \Gamma\left(\frac{l}{2}\right)} \frac{s}{s^m (s^2 + \omega^2)} + \frac{k \sqrt{\frac{\aleph_2 \alpha}{\aleph_3}}}{\left[\frac{\aleph_5 \aleph_2 \alpha}{\aleph_3} - \aleph_2 \right]} \\ & \times \frac{1}{s^2} \exp \left\{ -y \sqrt{\frac{\aleph_2 \alpha s}{\aleph_3 (s + \alpha \lambda)}} \right\} - \frac{G_r \aleph_4}{\left[\frac{p_r \aleph_5 \aleph_0 \alpha}{\aleph_1} - \aleph_2 \right]} \frac{1}{s(s + \alpha \lambda)} \exp \left\{ -y \sqrt{\frac{p_r \aleph_0 \alpha s}{\aleph_1}} \right\}, \end{aligned} \quad (46)$$

Inverting equations (46), we expressed the general solution of velocity as

$$w_c = UH(t) \cos \omega t + UH(t) \sum_{l=1}^{\infty} \frac{1}{l!} \left(\frac{-y \sqrt{\aleph_2 \alpha}}{\aleph_5} \right)^l \int_0^t \cos \omega(t - \tau) \mathbf{H}_{1,3}^{1,1} \left[\left(\frac{\alpha \lambda}{t} \right) \left| \begin{matrix} \left(1 - \frac{l}{2}, 1\right) \\ (0,1), \left(1 - \frac{l}{2}, 0\right), (0,1) \end{matrix} \right. \right] d\tau$$

$$+ \frac{k \sqrt{\frac{\mathcal{N}_2 \alpha}{\mathcal{N}_3}}}{\left[\frac{\mathcal{N}_5 \mathcal{N}_2 \alpha}{\mathcal{N}_3} - \mathcal{N}_2 \right]} \int_0^t \phi \left(y, \tau, \frac{\mathcal{N}_2 \alpha}{\mathcal{N}_3}, \alpha \lambda \right) d\tau - \frac{G_r \mathcal{N}_4}{\left[\frac{\mathcal{N}_5 \mathcal{N}_0 \alpha p_r}{\mathcal{N}_1} - \mathcal{N}_2 \right]} \int_0^t \phi \left(y, \tau, \frac{\mathcal{N}_0 \alpha p_r}{\mathcal{N}_1}, \alpha \lambda \right) \exp(-\alpha \lambda (t - \tau)) d\tau. \quad (47)$$

Where, $L^{-1} \left\{ \frac{1}{s^2} \exp \left(-y \sqrt{\frac{As}{s+B}} \right) \right\} = \int_0^t \phi(y, \tau, A, B) d\tau$ and $L^{-1} \left\{ \frac{1}{s} \exp \left(-y \sqrt{\frac{As}{s+B}} \right) \right\} = \phi(y, \tau, A, B)$.

The case of sine oscillations has been established by applying similar algorithm

$$w_s = U \sin \omega t + U \sum_{l=1}^{\infty} \frac{1}{l!} \left(\frac{-y \sqrt{\mathcal{N}_2 \alpha}}{\mathcal{N}_5} \right)^l \int_0^t \sin \omega (t - \tau) \times \mathbf{H}_{1,3}^{1,1} \left[\left(\frac{\alpha \lambda}{t} \right) \middle| \begin{matrix} \left(1 - \frac{l}{2}, 1 \right) \\ (0,1), \left(1 - \frac{l}{2}, 0 \right), (0,1) \end{matrix} \right] d\tau$$

$$+ \frac{k \sqrt{\frac{\mathcal{N}_2 \alpha}{\mathcal{N}_3}}}{\left[\frac{\mathcal{N}_2 \mathcal{N}_5 \alpha}{\mathcal{N}_3} - \mathcal{N}_2 \right]} \int_0^t \phi \left(y, \tau, \frac{\mathcal{N}_2 \alpha}{\mathcal{N}_3}, \alpha \lambda \right) d\tau - \frac{G_r \mathcal{N}_4}{\left[\frac{\mathcal{N}_5 \mathcal{N}_0 p_r \alpha}{\mathcal{N}_1} - \mathcal{N}_2 \right]} \int_0^t \phi \left(y, \tau, \frac{\mathcal{N}_0 p_r \alpha}{\mathcal{N}_1}, \alpha \lambda \right) \exp(-\alpha \lambda (t - \tau)) d\tau. \quad (48)$$

5. Limiting Cases

5.1 Investigation of Regular or Conventional Nanofluid $K_1 = 0$.

The solutions for velocity field to the regular or conventional Nanofluid are established by letting $K_1 = 0$ (in the absence of microrotation parameter) in the equations (47-48), we obtain

$$w_c = UH(t) \cos \omega t + UH(t) \sum_{l=1}^{\infty} \frac{1}{l!} \left(\frac{-y \sqrt{\mathcal{N}_2 \alpha}}{\mathcal{N}_5} \right)^l \int_0^t \cos \omega (t - \tau) \mathbf{H}_{1,3}^{1,1} \left[\left(\frac{\alpha \lambda}{t} \right) \middle| \begin{matrix} \left(1 - \frac{l}{2}, 1 \right) \\ (0,1), \left(1 - \frac{l}{2}, 0 \right), (0,1) \end{matrix} \right] d\tau$$

$$- \frac{G_r \mathcal{N}_4}{\left[\frac{\mathcal{N}_5 \mathcal{N}_0 p_r \alpha}{\mathcal{N}_1} - \mathcal{N}_2 \right]} \int_0^t \phi \left(y, \tau, \frac{\mathcal{N}_0 p_r \alpha}{\mathcal{N}_1}, \alpha \lambda \right) \exp(-\alpha \lambda (t - \tau)) d\tau. \quad (49)$$

$$\begin{aligned}
w_s = & U \sin \omega t + U \sum_{l=1}^{\infty} \frac{1}{l!} \left(\frac{-y \sqrt{\aleph_2 \alpha}}{\aleph_5} \right)^l \int_0^t \sin \omega(t-\tau) \mathbf{H}_{1,3}^{1,1} \left[\left(\frac{\alpha \lambda}{t} \right) \middle| \begin{matrix} \left(1 - \frac{l}{2}, 1 \right) \\ (0,1), \left(1 - \frac{l}{2}, 0 \right), (0,1) \end{matrix} \right] d\tau \\
& - \frac{G_r \aleph_4}{\left[\frac{\aleph_5 \aleph_0 P_r \alpha}{\aleph_1} - \aleph_2 \right]} \int_0^t \phi \left(y, \tau, \frac{\aleph_0 P_r \alpha}{\aleph_1}, \alpha \lambda \right) \exp(-\alpha \lambda(t-\tau)) d\tau.
\end{aligned} \tag{50}$$

5.2 Investigation of Regular or Conventional Newtonian Fluid $K_1 = \phi = 0$.

It is also pointed out that, the analytic solutions to the regular or conventional Newtonian fluid can be recovered from equations (47-48) by letting $K_1 = \phi = 0$ (in the absence of microrotation parameter). In this continuation, one can also transform the analytic solutions in to ordinary differential operator by substituting $\alpha = 1$.

6. Results and Conclusion

In this research article, micropolar nanofluid has proved to give better thermal performance than conventional fluids based on the mathematical tools of non-integer order fractional derivative and transforms. The analysis has shown vivid effects for the enhancement of high thermal conductivity subject to suspended nanoparticles in to the base fluid. The graphical illustrations have been discussed on the investigated solutions which rectify physical conditions. Various graphs have been depicted by using table 1 for highlighting the effects of nanoparticles and embedded parameters of micropolar nanofluid. However, the key results are enumerated below:

- (i) The analytic solutions are explored for temperature, microrotation and velocity and the similar solutions for velocity field and temperature distribution to the regular or conventional nanofluid $K_1 = 0$. and Newtonian fluid $K_1 = \phi = 0$. have also been recovered as the limiting cases.
- (ii) Fig. 2 is prepared for the influence of nanoparticles based on the two types of solutions i-e fractionalized nanofluids $\alpha = 0.4$ and ordinary nanofluids $\alpha = 1.0$ in which velocity field of copper- ethylene glycol is higher in comparison with pure ethylene glycol and silver- ethylene glycol. It is noted that Velocity field of fractionalized nanofluids $\alpha = 0.4$ has reciprocal behavior with ordinary nanofluids $\alpha = 1.0$. This may be due to the fact of exponential kernel in Caputo-Fabrizio fractional derivative.

- (iii) Fig. 3 is depicted for temperature distribution with and without Caputo-Fabrizio fractional derivative for the influence of pure ethylene glycol, copper-ethylene glycol and silver-ethylene glycol. In this figure, copper-ethylene glycol has scattering effects for fractionalized temperature distribution and has reversal impacts for ordinary temperature distribution.
- (iv) The influence of nanoparticles on microrotation is underlined in Fig. 4, in which effect of microrotation is observed opposite near the plate. It is also pointed out that copper-ethylene glycol has accelerating behavior in comparison to all others; on the contrary copper-ethylene glycol has decelerating nature near the plate. This may be due to fact of effective fractional parameter α .
- (v) Figs: 5(a), 5(b) and 5(c) are plotted to investigate the effects of different values of volume fraction $\phi = 0.0, 0.01, 0.02$ on velocity field. Here, nanoparticles are suspended as copper-ethylene glycol and silver-ethylene glycol for the velocity field with and without Caputo-Fabrizio fractional derivative. The increase in volume fraction results the scattering behavior of copper-ethylene glycol with Caputo-Fabrizio fractional operator and sequestering behavior of copper-ethylene glycol without Caputo-Fabrizio fractional operator. It is significantly noted that converse phenomenon is observed as the values of volume fraction increases, i-e (silver-ethylene glycol for the velocity field with Caputo-Fabrizio fractional derivative has scattering behavior and copper-ethylene glycol has sequestering behavior). This is due to natural fact that, when temperature is lower than $180^{\circ}C$ then increase in volume fraction generates increase in thermal conductivity. Same phenomenon can also be observed for temperature distribution and microrotation as well.

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Ethical Statement

This material has not been published in whole or in parts elsewhere. The manuscript is not currently being considered for publication in another journal. All authors have been personally and actively involved in substantive work leading to the manuscript, and will hold themselves jointly and individually responsible for its content.

Appendix

$$\mathfrak{N}_0 = \frac{(C_p \rho)_s \varphi}{(C_p \rho)_f} + (1 - \varphi), \quad (\text{A1})$$

$$\mathfrak{N}_1 = R + \frac{k_{nf}}{k_f}, \quad (\text{A2})$$

$$\mathfrak{N}_2 = \frac{\varphi \rho_s}{\rho_f} + (1 - \varphi), \quad (\text{A3})$$

$$\mathfrak{N}_3 = \frac{1}{(1 - \varphi)^{2.5}} + \frac{K}{2}, \quad (\text{A4})$$

$$\mathfrak{N}_4 = \frac{(\beta_T \rho)_s \varphi}{(\beta_T \rho)_f} + (1 - \varphi), \quad (\text{A5})$$

$$\mathfrak{N}_5 = \frac{1}{(1 - \varphi)^{2.5}} + k. \quad (\text{A6})$$

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Biographies

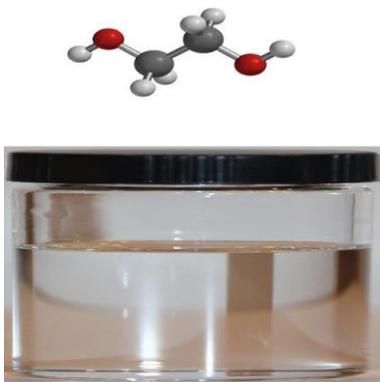
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Table: 1. Fundamental Thermo-physical Properties

Base Fluid/Nanoparticles	ρ (Kg/m ³)	C_p (J/Kg K)	k (W/m.
Cu	8933	385	401
Ag	10500	235	429
Ethylene Glycol	1.115	0.58	0.1490

Fig. 1(a). Ethylene Glycol



1(b). Copper Metal Powder



1(c). Silver Metal Powder



Fig.2. Profile of velocity field with and without Caputo-Fabrizio derivative.

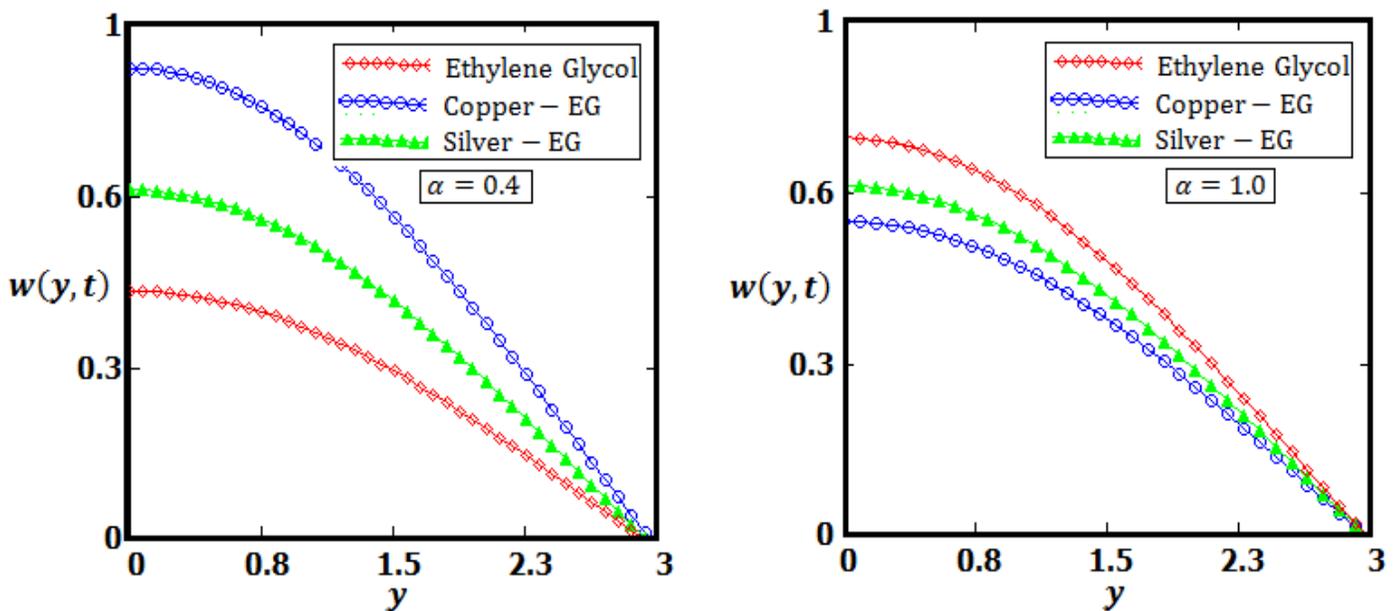


Fig.3. Profile of temperature distribution with and without Caputo-Fabrizio derivative.

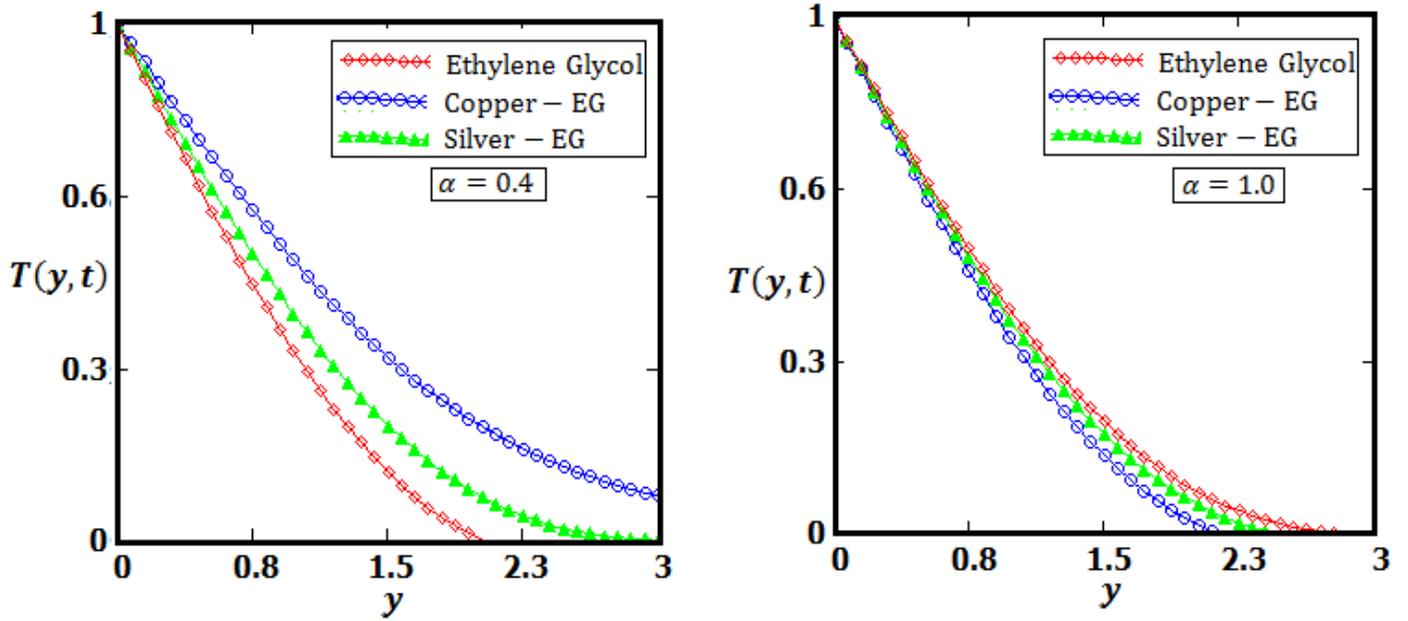


Fig.4. Profile of microrotation with and without Caputo-Fabrizio derivative.

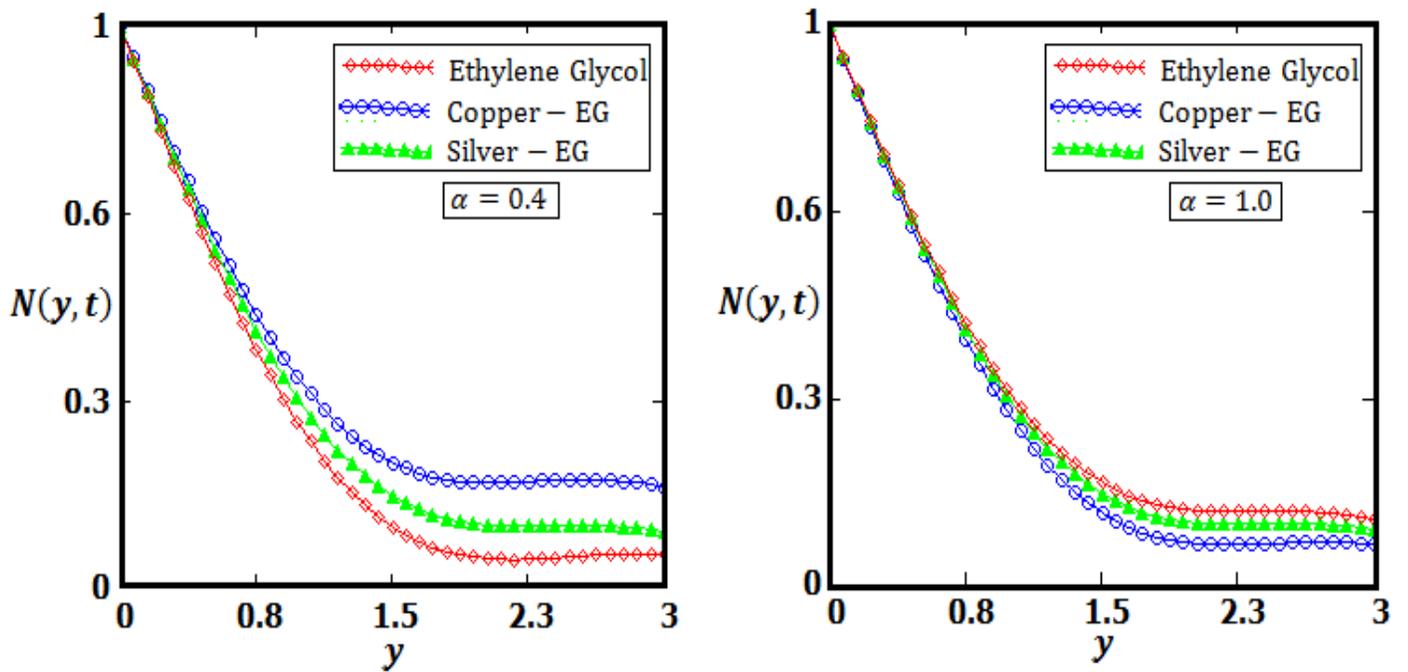


Fig.5(a). Profile of velocity field with and without Caputo-Fabrizio derivative when nanoparticle volume fraction is $\phi = 0.0$.

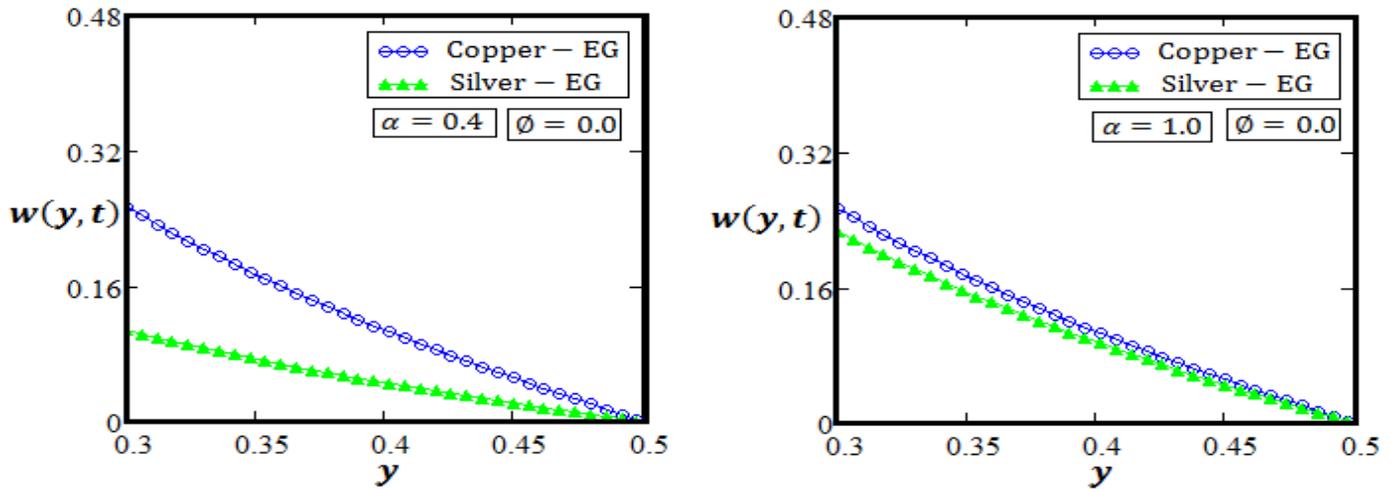


Fig.5(b). Profile of velocity field with and without Caputo-Fabrizio derivative when nanoparticle volume fraction is $\phi = 0.01$.

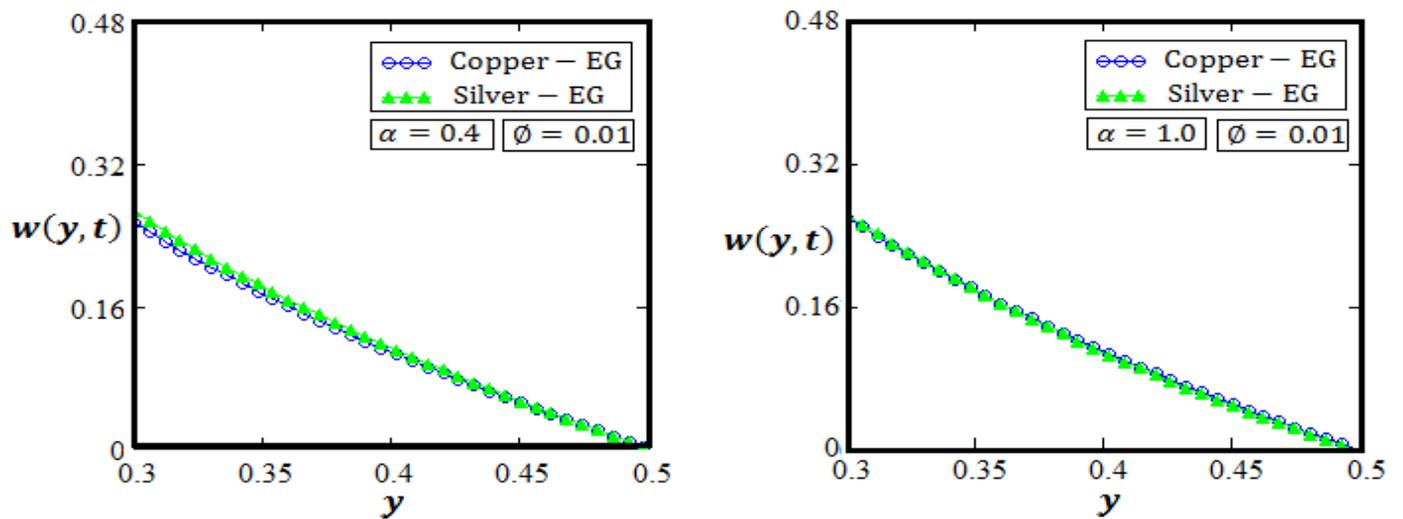


Fig.5(c). Profile of velocity field with and without Caputo-Fabrizio derivative when nanoparticle volume fraction is $\phi = 0.02$.

